

Post-TOV modelling of relativistic stars and astrophysical applications

$$G_{\mu\nu} = 8\pi T_{\mu\nu} ?$$

(based on PRD 92, 2015 & 94, 2016)

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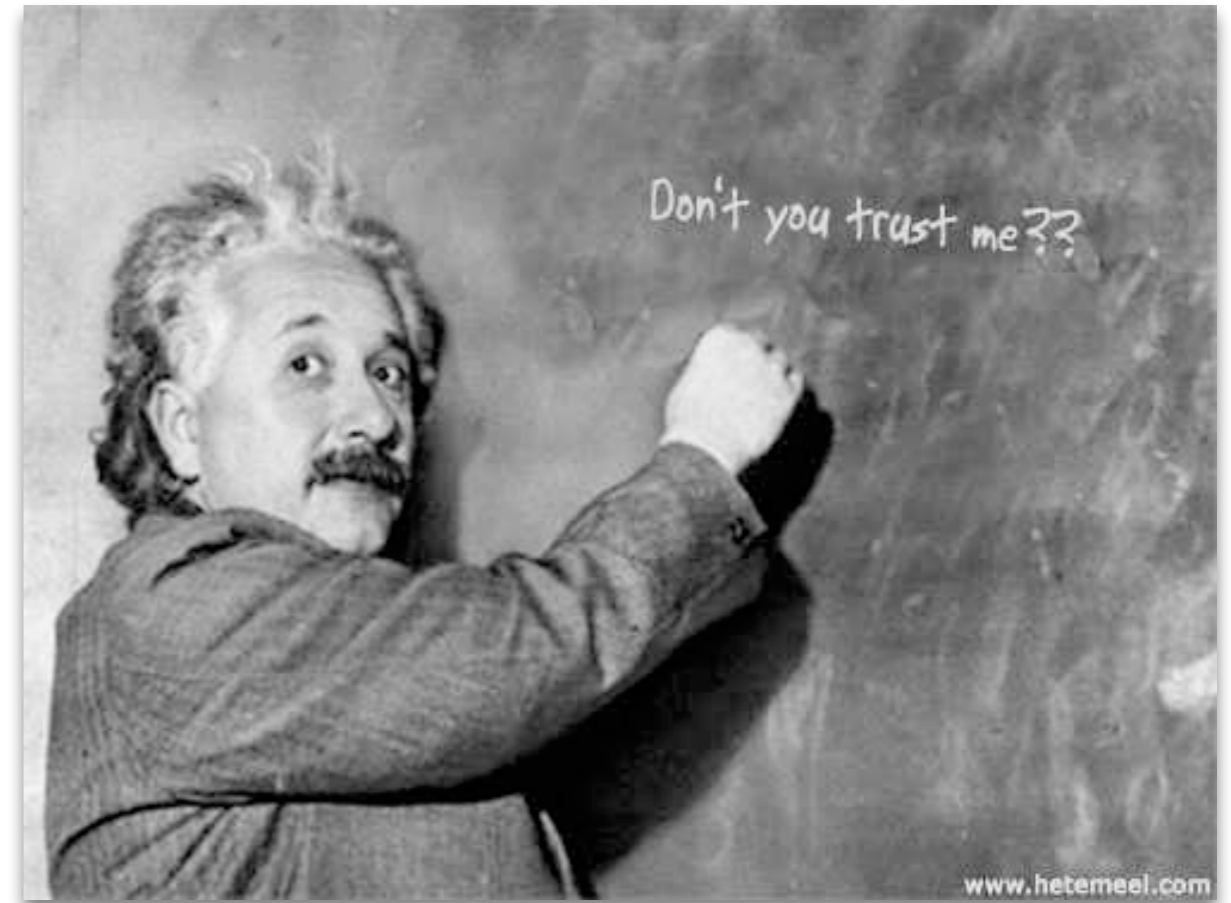
IAP, March 2017

Outline

- Neutron stars as gravity probes.
- Gravity-matter degeneracy in neutron stars.
- The post-TOV formalism (interior and exterior).
- Astrophysical applications: redshift, X-ray bursts and QPOs.
- Further/future extensions of the formalism.

GR-exit?

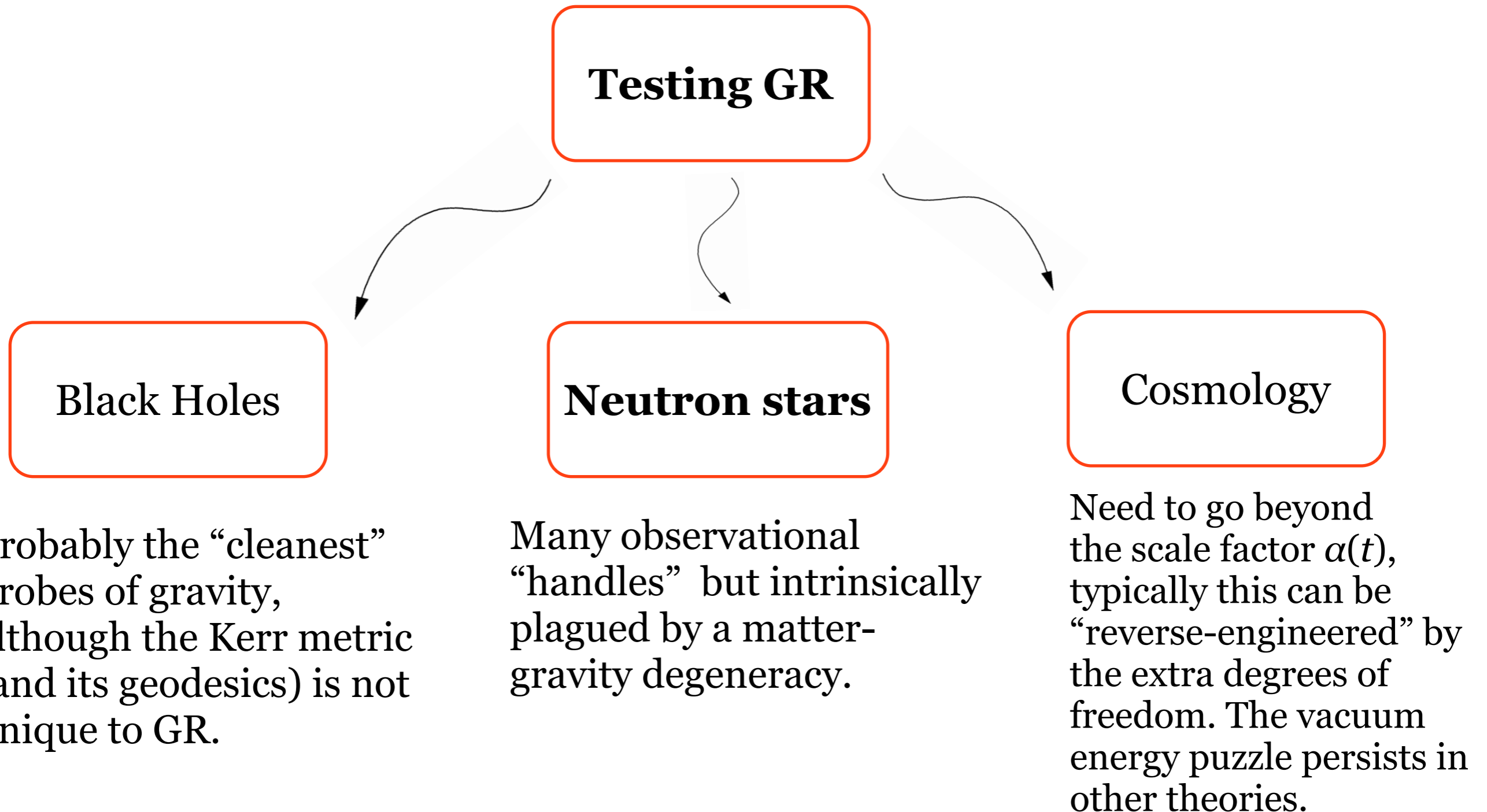
- General Relativity is arguably the most elegant theory invented so far and is, most likely, *The* theory of (classical) gravity.
- Good science: test all theories, no matter how elegant!
- Build phenomenology
—> how “special” is GR?
- A host of modified theories of gravity on the market.



The zoo of gravity theories

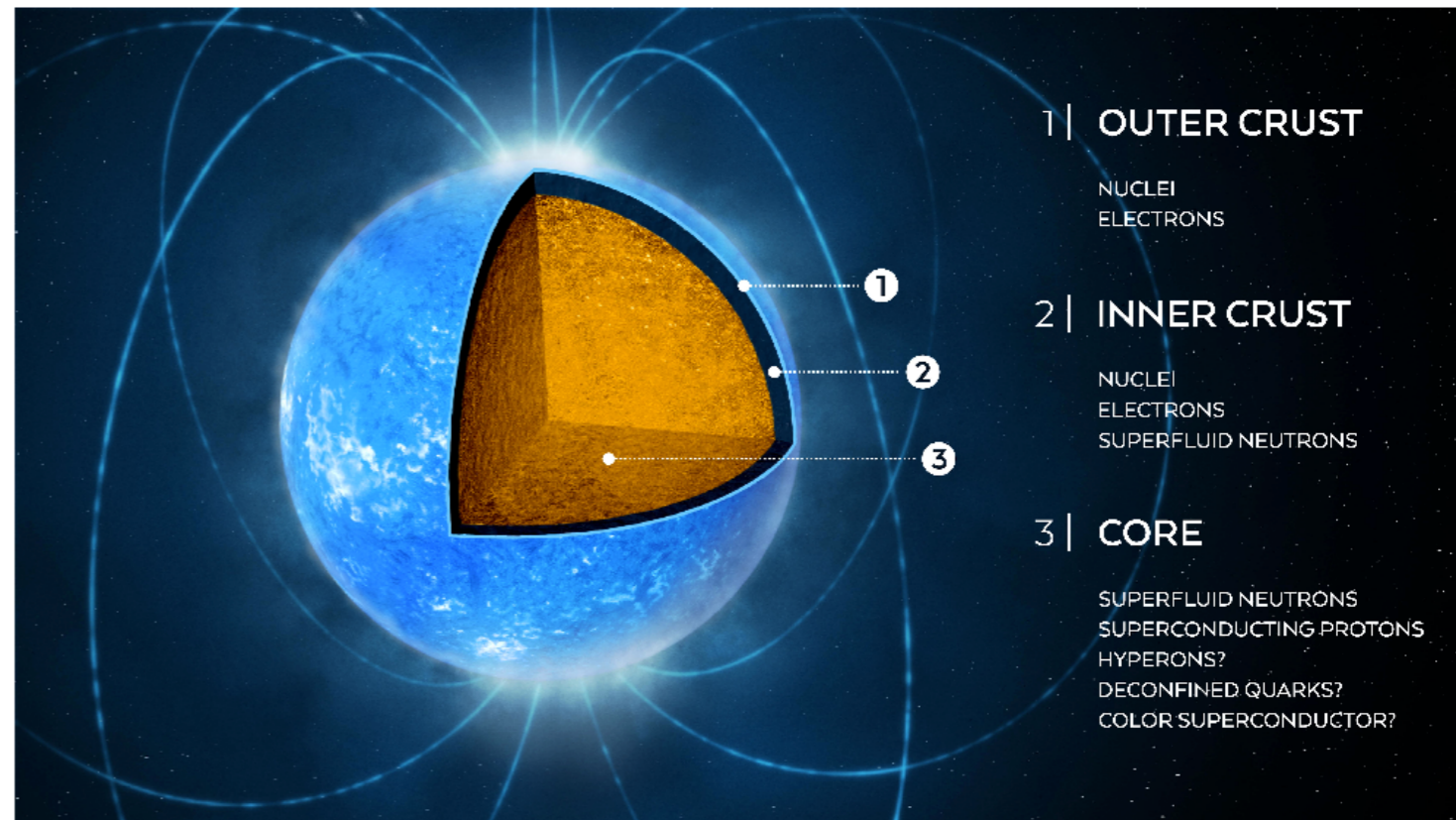
Theory	Field content	Strong EP	Massless graviton	Lorentz symmetry	Linear $T_{\mu\nu}$	Weak EP	Well-posed?	Weak-field constraints
Extra scalar field								
Scalar-tensor	S	\times	\checkmark	\checkmark	\checkmark	\checkmark	\checkmark [30]	[31-33]
Multiscalar	S	\times	\checkmark	\checkmark	\checkmark	\checkmark	\checkmark ?	[34]
Metric $f(R)$	S	\times	\checkmark	\checkmark	\checkmark	\checkmark	\checkmark [35,36]	[37]
Quadratic gravity								
Gauss-Bonnet	S	\times	\checkmark	\checkmark	\checkmark	\checkmark	\checkmark ?	[38]
Chern-Simons	P	\times	\checkmark	\checkmark	\checkmark	\checkmark	$\times\checkmark$? [39]	[40]
Generic	S/P	\times	\checkmark	\checkmark	\checkmark	\checkmark	?	
Horndeski	S	\times	\checkmark	\checkmark	\checkmark	\checkmark	\checkmark ?	
Lorentz-violating								
\mathcal{A} -gravity	SV	\times	\checkmark	\times	\checkmark	\checkmark	\checkmark ?	[41-44]
Khronometric/ Hořava-Lifshitz	S	\times	\checkmark	\times	\checkmark	\checkmark	\checkmark ?	[43-46]
n-DBI	S	\times	\checkmark	\times	\checkmark	\checkmark	?	none ([47])
Massive gravity								
dRGT/Bimetric	SVT	\times	\times	\checkmark	\checkmark	\checkmark	?	[16]
Galileon	S	\times	\checkmark	\checkmark	\checkmark	\checkmark	\checkmark ?	[16,48]
Nondynamical fields								
Palatini $f(R)$	–	\checkmark	\checkmark	\checkmark	\times	\checkmark	\checkmark	none
Eddington-Born-Infeld	–	\checkmark	\checkmark	\checkmark	\times	\checkmark	?	none
Others, not covered here								
TeV S	SVT	\times	\checkmark	\checkmark	\checkmark	\checkmark	?	[33]
$f(R)\mathcal{L}_m$?	?	\checkmark	\checkmark	\checkmark	\times	?	
$f(T)$?	\times	\checkmark	\times	\checkmark	\checkmark	?	[49]

Testing strong gravity



Neutron stars on a napkin

- Relativistic objects $R/M \sim 5$.
- Supra-nuclear density matter.
- Lots of exotic physics (superfluidity, superconductivity, deconfined quarks ...).
- Fast rotation, strong B-fields.
- Future experiments (SKA, NICER ...) should provide accurate NS information (masses, radii).

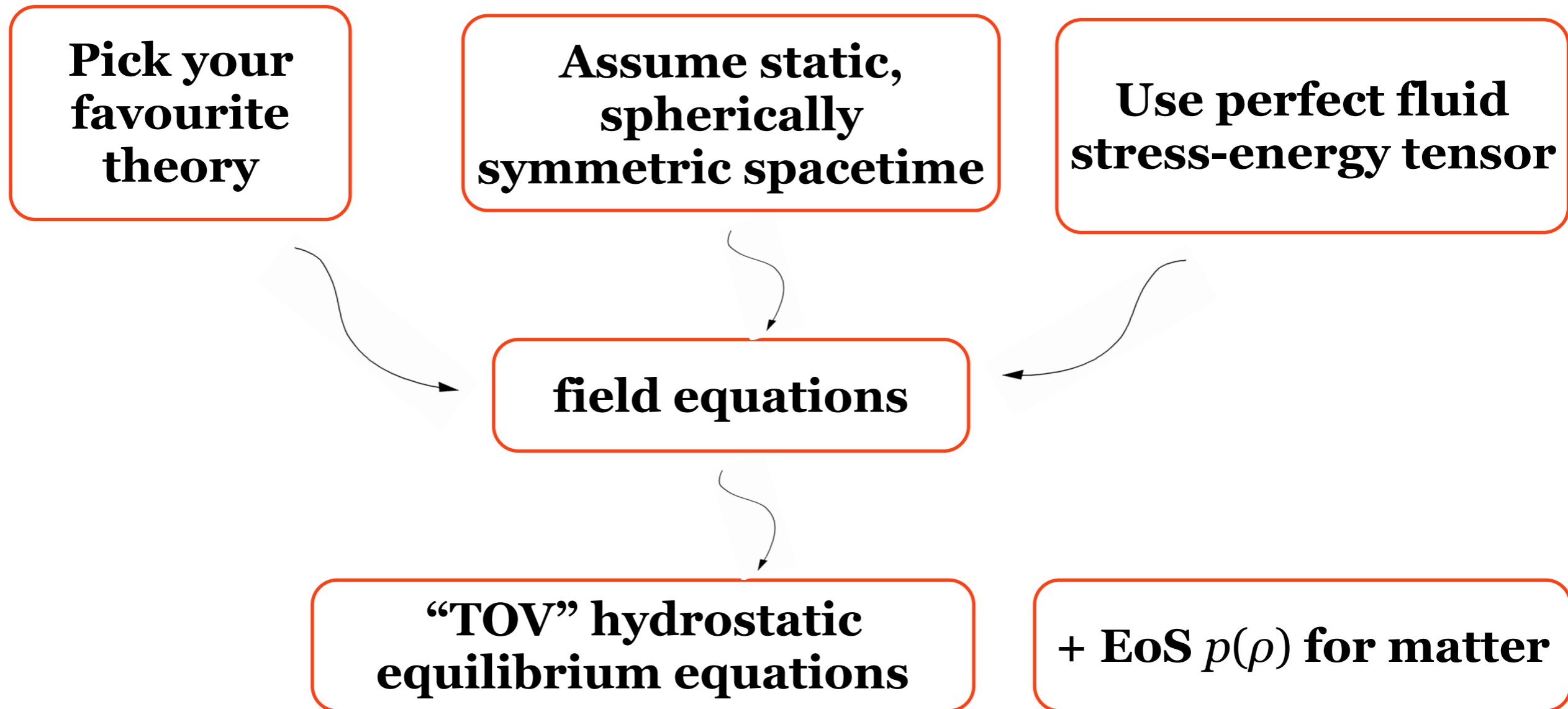


$$M \sim M_{\text{P}} \left(\frac{M_{\text{P}}}{m_n} \right)^2 \approx 1.5 M_{\odot}$$

$$R \sim \lambda_n \frac{M_{\text{P}}}{m_n} \approx 10^6 \text{ cm}$$

$$f_{\text{spin}} \lesssim 1 \text{ kHz}$$

Recipe for building neutron stars



Recipe for building neutron stars

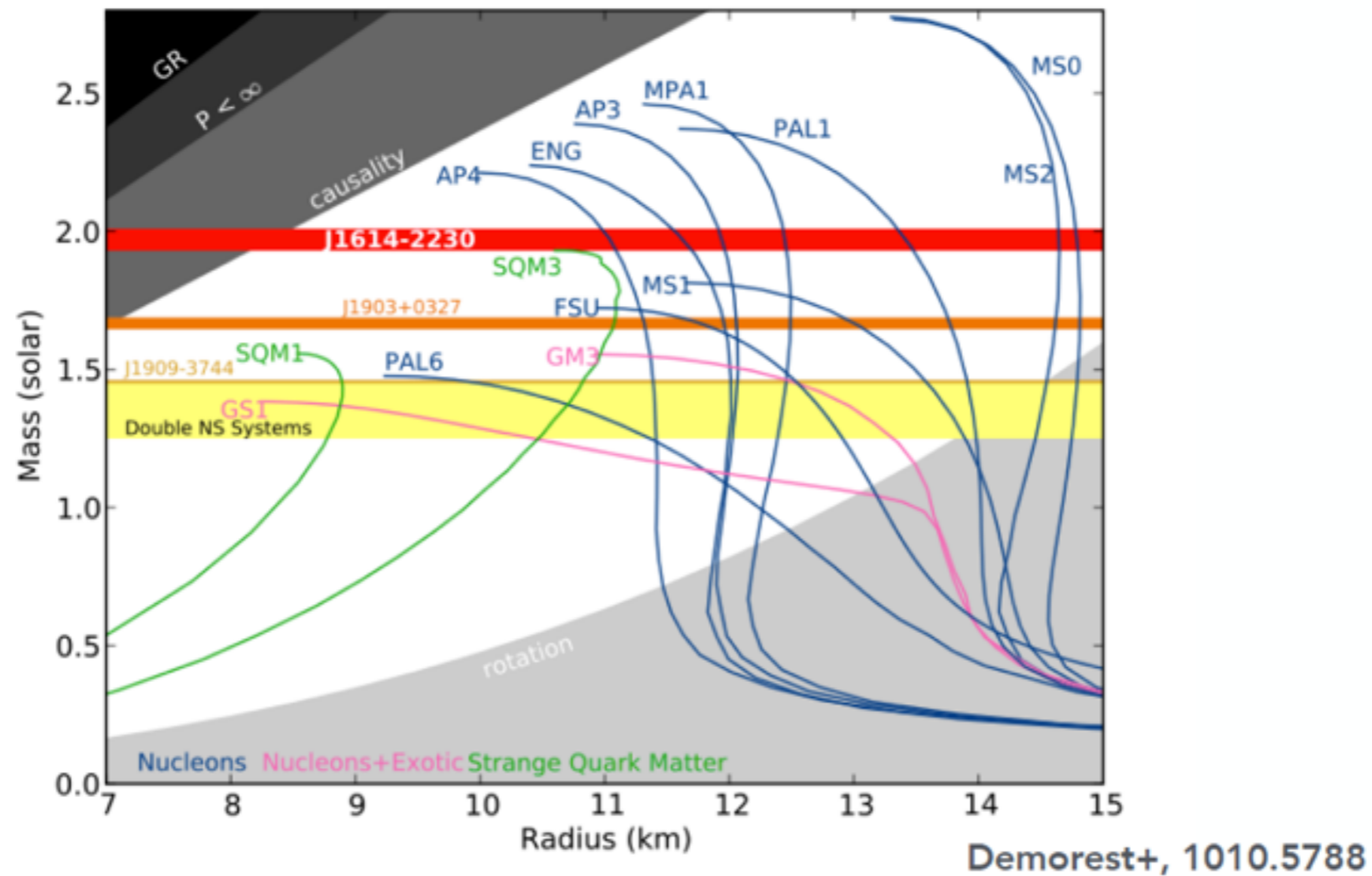
Done many times!

[Berti et al. 2015]

Theory	Structure		Collapse	Sensitivities	Stability	Geodesics	
	NR	SR					FR
Extra scalar field							
Scalar-Tensor	[109–114]	[112, 115, 116]	[117–119]	[120–127]	[128]	[129–139]	[118, 140]
Multiscalar	?	?	?	?	?	?	?
Metric $f(R)$	[141–153]	[154]	[155]	[156, 157]	?	[158, 159]	?
Quadratic gravity							
Gauss-Bonnet	[160]	[160]	[77]	?	?	?	?
Chern-Simons	\equiv GR	[25, 40, 161–163]	?	?	[162]	?	?
Horndeski	?	?	?	?	?	?	?
Lorentz-violating							
\mathcal{A} -gravity	[164, 165]	?	?	[166]	[43, 44]	[158]	?
Khronometric/ Hořava-Lifshitz	[167]	?	?	?	[43, 44]	?	?
n-DBI	?	?	?	?	?	?	?
Massive gravity							
dRGT/Bimetric	[168, 169]	?	?	?	?	?	?
Galileon	[170]	[170]	?	[171, 172]	?	?	?
Nondynamical fields							
Palatini $f(R)$	[173–177]	?	?	?	–	?	?
Eddington-Born-Infeld	[178–184]	[178, 179]	?	[179]	–	[185, 186]	?

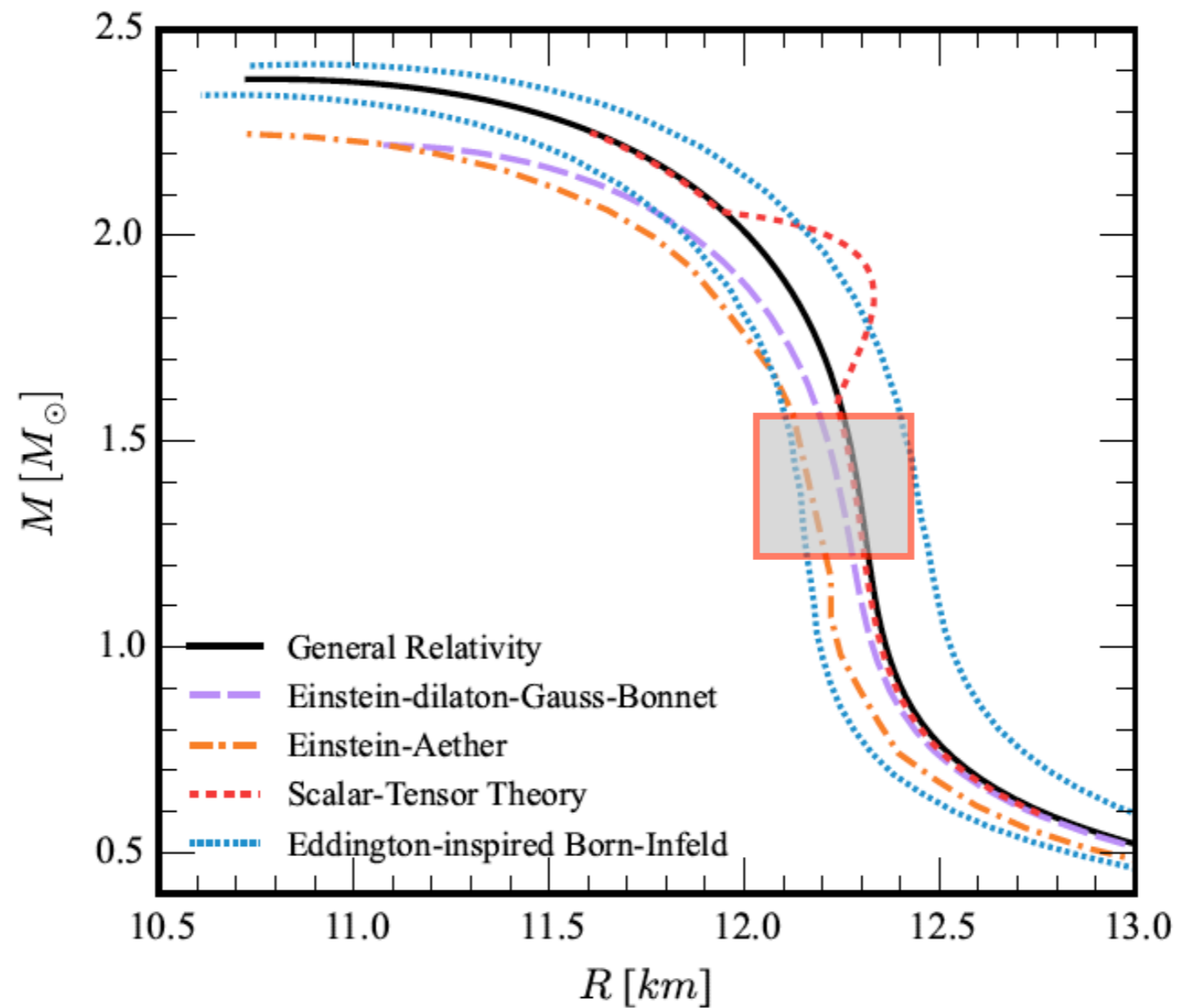
NR = non-rotating, SR = slow rotation, FR = fast rotation

“Matter-gravity” degeneracy



Fixed gravity theory (GR), varying EoS

“Matter-*gravity*” degeneracy



Fixed EoS (APR), varying gravity

(different theories can cause similar deviations from GR)

A different approach

- *Do not commit to any particular “favourite” theory of gravity.*
- Parametrize deviations from GR in a PPN theory-manner.
- At the same time allow for strong gravity, by going beyond “simple” PN expansions.
- Examples from the literature:
 - “Bumpy” or “quasi-Kerr” BH metrics,
 - “post-Friedmannian” cosmology,
 - “post-Einsteinian” GW waveforms.



Strategy for a “post-TOV” formalism

- This is a formalism for *relativistic stars in spherical symmetry* (i.e. no rotation), *not* a fully-fledged theory of gravity.

- **Main idea:**

augment the General Relativistic TOV stellar structure equations by adding 1PN and 2PN corrections with *arbitrary coefficients*.

These terms are to be built out of the available parameters:

$$p, \rho, \Pi, m, r \quad \text{matter energy density: } \epsilon = \rho(1 + \Pi)$$

- The hydrostatic equations for the pressure $p(r)$ and mass function $m(r)$ take the symbolic form:

$$\frac{dp}{dr} = \left(\frac{dp}{dr} \right)_{\text{GR}} + \{ \text{PN corrections} \}, \quad \frac{dm}{dr} = \left(\frac{dm}{dr} \right)_{\text{GR}} + \{ \text{PN corrections} \}$$

Looking for post-TOV corrections

- **1PN order:** these can be extracted from the existing PPN theory:

$$\Lambda_1 \sim \Pi, \frac{m}{r}, \frac{r^3 p}{m}$$

- **2PN :** use dimensional analysis (excluding the presence of dimensional coupling constants):

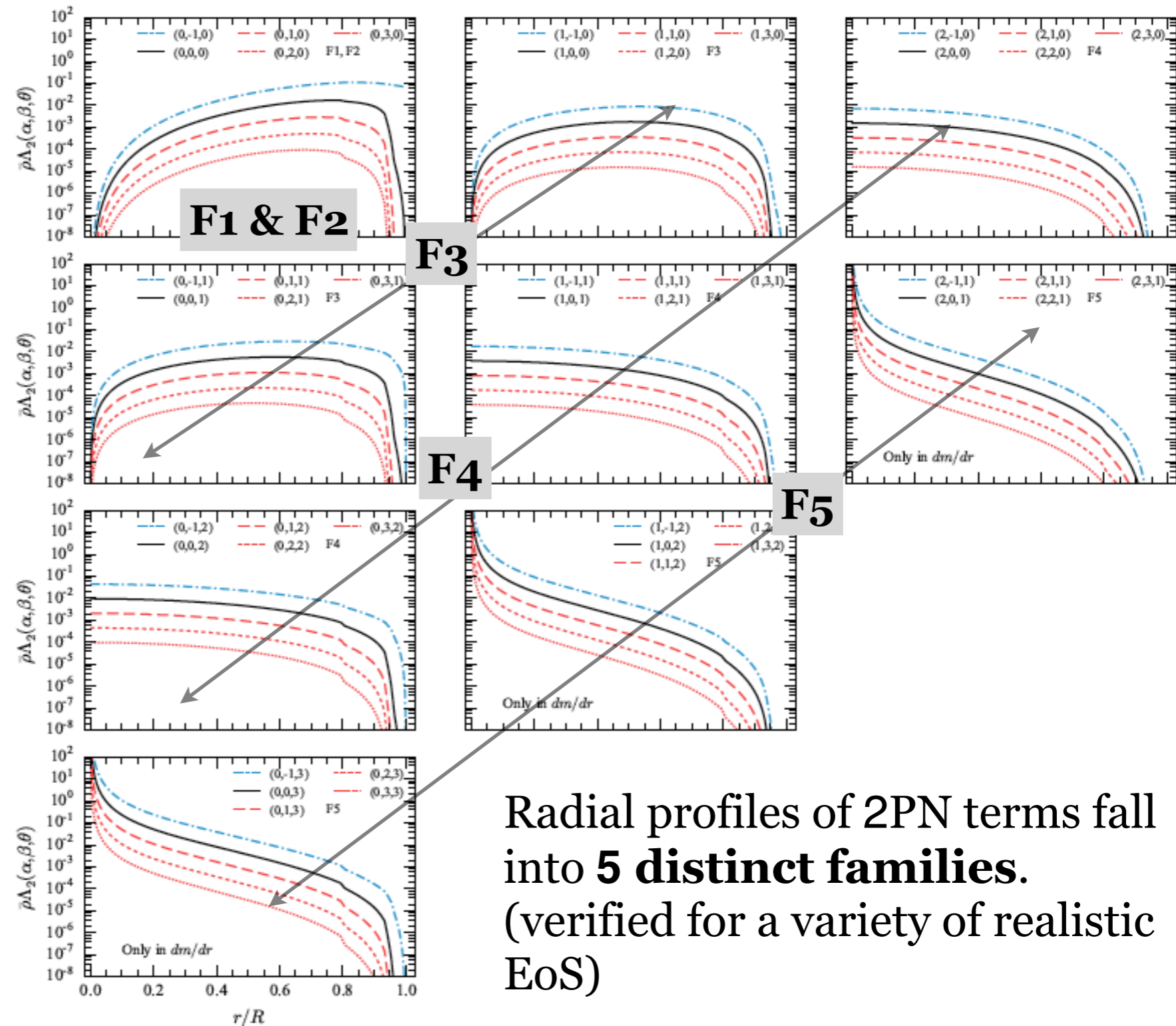
General 2PN term (dimensionless): $\Lambda_2 \sim \Pi^\theta (r^2 p)^\alpha (r^2 \rho)^\beta \left(\frac{m}{r}\right)^{2-2\alpha-\beta-\theta}$

- Limits on $\{\alpha, \beta, \theta\}$: (i) avoid divergence at the stellar center/surface and (ii) assume field equations *linear* in the stress-energy tensor:

$$\left\{ \begin{array}{l} \{\text{geometry}\} \sim 8\pi T^{\mu\nu} \\ \{\text{geometry}\} \sim (\epsilon + \tau p)^n \\ \tau, n = \mathcal{O}(1) \end{array} \right. \left\{ \begin{array}{l} \beta \geq -1 \\ 0 \leq \theta \leq 2 \text{ or } 3 \\ 0 \leq \alpha \leq 2 - \theta \text{ or } 3 - \theta \end{array} \right. \quad \begin{array}{l} \text{In principle,} \\ \text{there are infinite} \\ \text{2PN terms ...} \end{array}$$

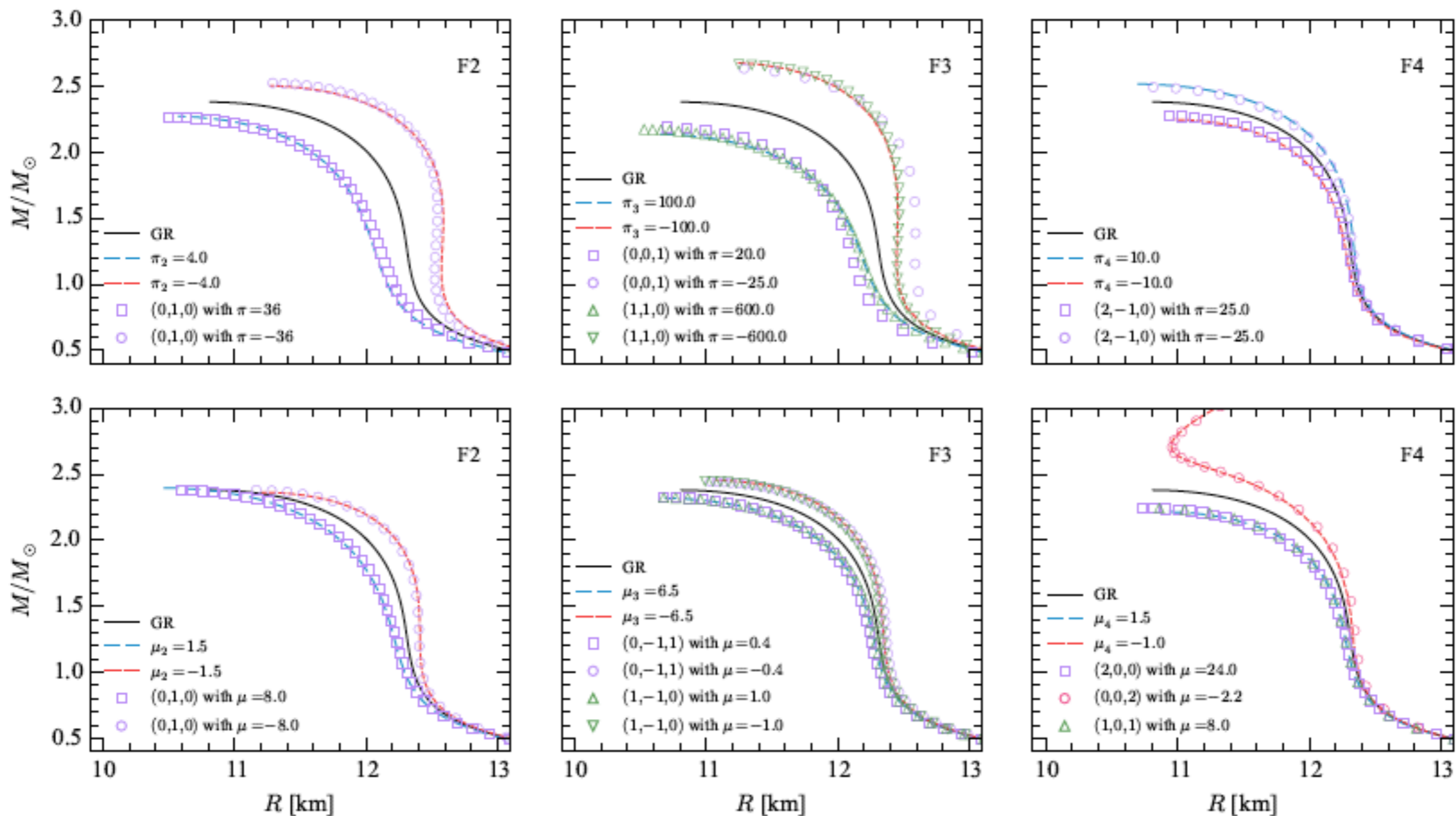
It's all about Families!

Family	2PN term	(α, β, θ)
F1	$m^3/(r^5 \rho)$	$(0, -1, 0)$
F2	$(m/r)^2$	$(0, 0, 0)$
F2	$rm\rho$	$(0, 1, 0)$
F3	$mp/(r\rho)$	$(1, -1, 0)$
F3	$r^2 p$	$(1, 0, 0)$
F3	$\Pi m^2/(r^4 \rho)$	$(0, -1, 1)$
F3	$\Pi m/r$	$(0, 0, 1)$
F3	$r^2 \Pi \rho$	$(0, 1, 1)$
F4	$r^3 p^2/(\rho m)$	$(2, -1, 0)$
F4	$r^6 p^2/(m^2)$	$(2, 0, 0)$
F4	$\Pi p/\rho$	$(1, -1, 1)$
F4	$\Pi r^3 p/m$	$(1, 0, 1)$
F4	$\Pi^2 m/(r^3 \rho)$	$(0, -1, 2)$
F4	Π^2	$(0, 0, 2)$
F5	$\Pi r^4 p^2/(\rho m^2)$	$(2, -1, 1)$
F5	$\Pi r^7 p^2/m^3$	$(2, 0, 1)$
F5	$\Pi^2 r p/m\rho$	$(1, -1, 2)$
F5	$\Pi^2 r^4 p/m^2$	$(1, 0, 2)$
F5	$\Pi^3/(r^2 \rho)$	$(0, -1, 3)$
F5	$\Pi^3 r/m$	$(0, 0, 3)$



Families: M - R self-similarity

Remarkably, the 2PN terms of the same family lead to self-similar M - R curves when added as post-TOV corrections (results shown assume APR but have been verified for other EoS too).



Post-TOV structure equations (I)

$$\frac{dp}{dr} = \left(\frac{dp}{dr} \right)_{\text{GR}} - \frac{\rho m}{r^2} (\mathcal{P}_1 + \mathcal{P}_2)$$

$$\frac{dm}{dr} = \left(\frac{dm}{dr} \right)_{\text{GR}} + 4\pi r^2 \rho (\mathcal{M}_1 + \mathcal{M}_2)$$

- 1PN-order corrections:

$$\mathcal{P}_1 = \delta_1 \frac{m}{r} + 4\pi \delta_2 \frac{r^3 p}{m}$$

$$\mathcal{M}_1 = \delta_3 \frac{m}{r} + \delta_4 \Pi$$

- Current PPN limits: $|\delta_i| \ll 1 \rightarrow |\mathcal{P}_1|, |\mathcal{M}_1| \ll 1$

1PN terms can be ignored

Post-TOV structure equations (II)

$$\frac{dp}{dr} \approx \left(\frac{dp}{dr} \right)_{\text{GR}} - \frac{\rho m}{r^2} \mathcal{P}_2 \qquad \frac{dm}{dr} \approx \left(\frac{dm}{dr} \right)_{\text{GR}} + 4\pi r^2 \rho \mathcal{M}_2$$

- 2PN-order corrections:

use *just one representative term per family*

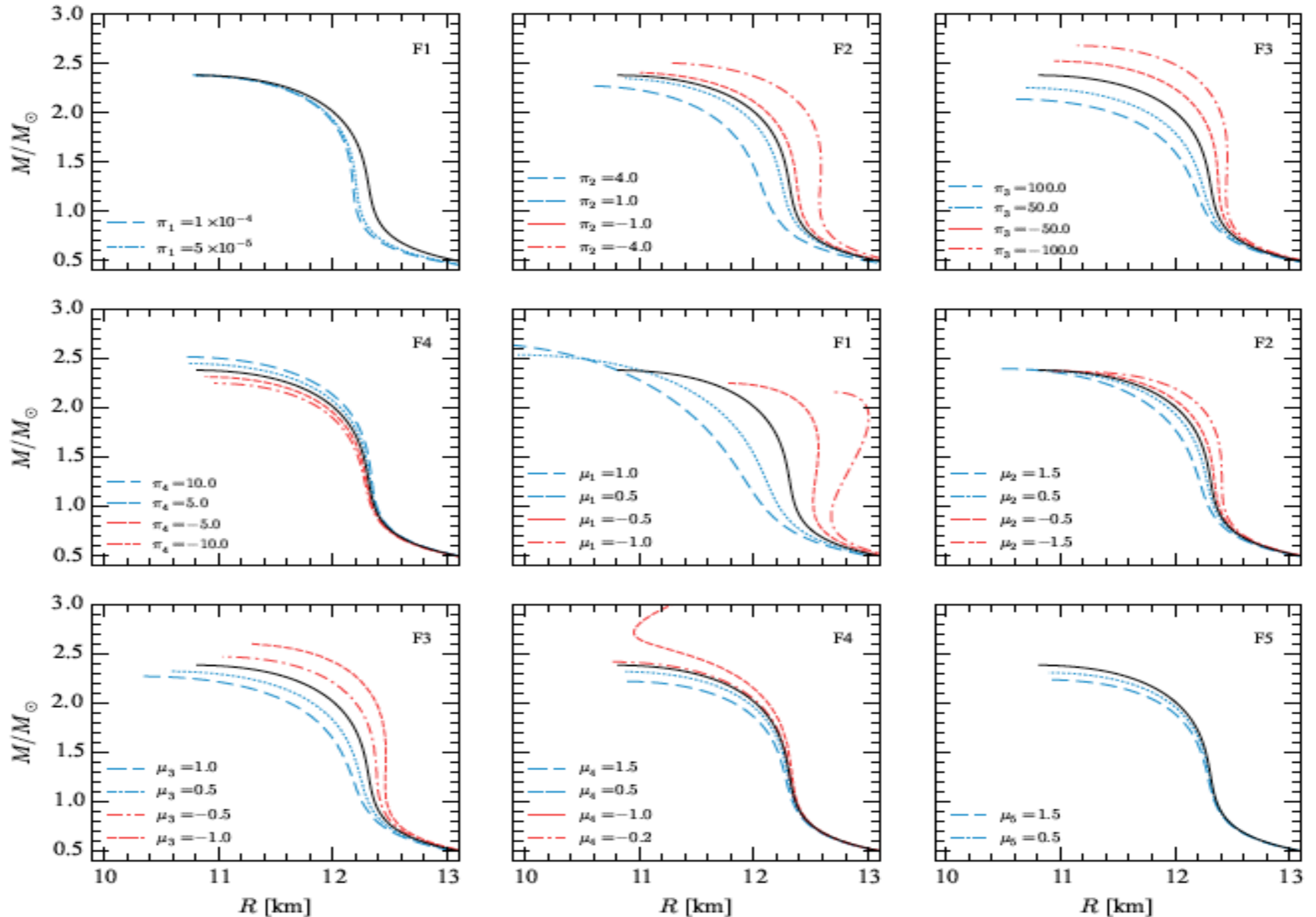
self-similarity: all other terms are accounted for by varying the corresponding coefficient

$\{\pi_i, \mu_i\} \leftrightarrow F_i$
family

$$\mathcal{P}_2 = \pi_1 \frac{m^3}{r^5 \rho} + \pi_2 \frac{m^2}{r^2} + \pi_3 r^2 p + \pi_4 \frac{\Pi p}{\rho}$$

$$\mathcal{M}_2 = \mu_1 \frac{m^3}{r^5 \rho} + \mu_2 \frac{m^2}{r^2} + \mu_3 r^2 p + \mu_4 \frac{\Pi p}{\rho} + \mu_5 \Pi^3 \frac{r}{m}$$


post-TOV: sample of M - R curves



Post-TOV as “effective GR”

- The post-TOV equations can be mapped onto an *effective* GR formulation:

$$\nabla_{\nu} T_{\text{eff}}^{\mu\nu} = 0, \quad T_{\text{eff}}^{\mu\nu} = (\epsilon_{\text{eff}} + p)u^{\mu}u^{\nu} + pg^{\mu\nu}$$


$$\frac{dp}{dr} = -\frac{1}{2} (\epsilon_{\text{eff}} + p) \frac{d\nu}{dr} \quad \frac{dm}{dr} = 4\pi r^2 \epsilon_{\text{eff}}$$

- *Gravity-shifted effective* EoS: $p = p(\epsilon_{\text{eff}}), \quad \epsilon_{\text{eff}} = \epsilon + \rho\mathcal{M}_2$

- *Effective interior* metric:

$$g_{\mu\nu} = \text{diag}[-e^{\nu(r)}, (1 - 2m(r)/r)^{-1}, r^2, r^2 \sin^2 \theta]$$

Exterior metric (I)

- Our scheme also allows the construction of an exterior metric.

- Set all fluid parameters to zero: $p = \epsilon = \rho = \Pi = 0$

- The resulting post-TOV equations are:

$$\left\{ \begin{array}{l} \frac{d\nu}{dr} = \left(\frac{d\nu}{dr} \right)_{\text{GR}} + 2(\pi_2 - \mu_2) \frac{m^3}{r^4} \\ \frac{dm}{dr} = 4\pi\mu_1 \frac{m^3}{r^3} \end{array} \right.$$

- Integrating and keeping the leading post-TOV terms:

$$m(r) \approx M_\infty \left(1 - 2\pi\mu_1 \frac{M_\infty^2}{r^2} \right) \quad \Rightarrow \quad M \approx \underbrace{M_\infty}_{\text{ADM mass}} \left(1 - 2\pi\mu_1 \frac{M_\infty^2}{R^2} \right)$$

$$\nu(r) \approx \log \left(1 - \frac{2M_\infty}{r} \right) - \frac{2\chi}{3} \frac{M_\infty^3}{r^3} \quad \text{Schwarzschild mass } M \equiv m(R)$$

Exterior metric (II)

- We assume the *same* effective metric form as in the interior:

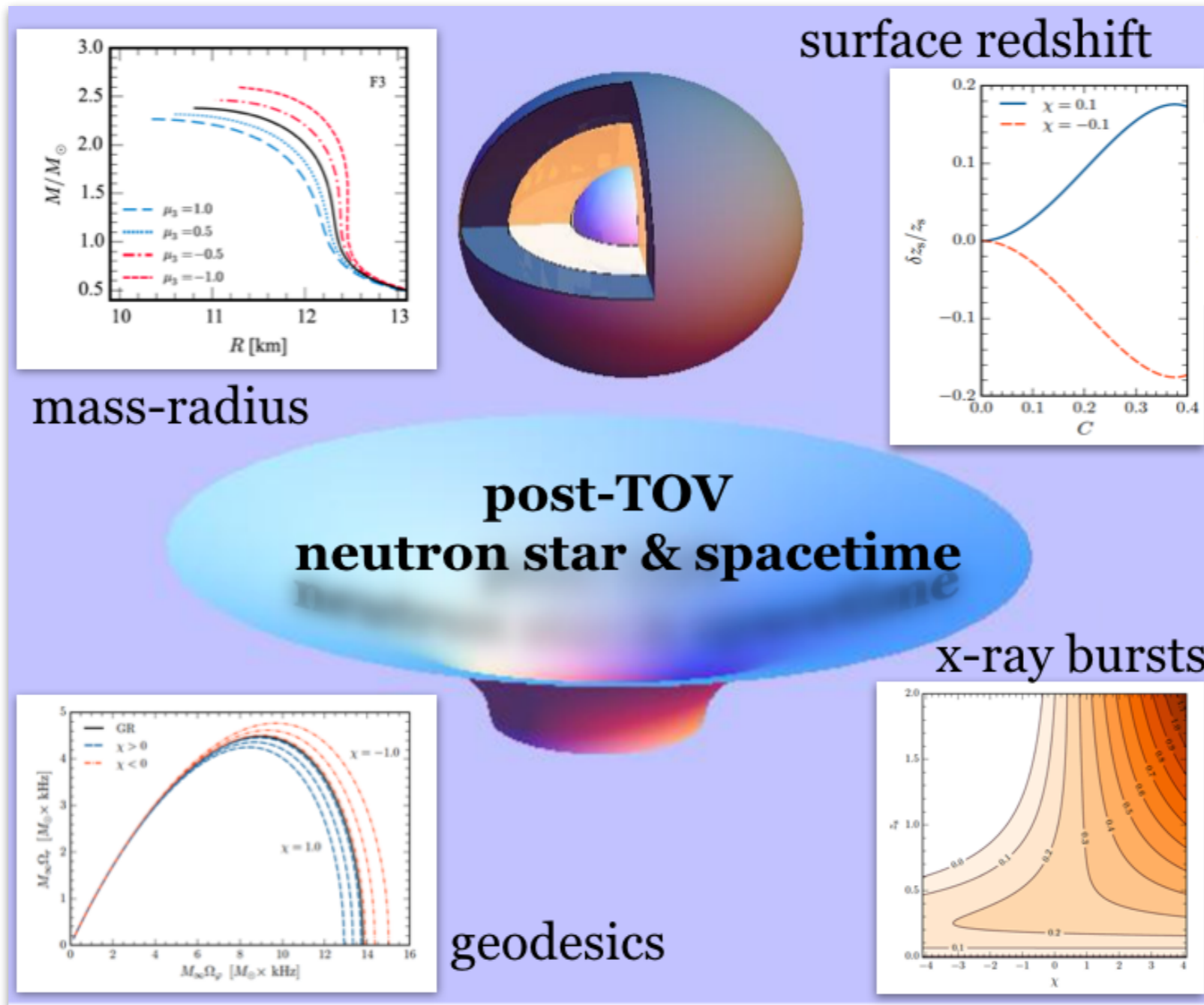
$$g_{\mu\nu} = \text{diag}[-e^{\nu(r)}, (1 - 2m(r)/r)^{-1}, r^2, r^2 \sin^2 \theta]$$

- The exterior metric takes a *post-Schwarzschild* form:

$$g_{tt}(r) \approx - \left(1 - \frac{2M_\infty}{r} \right) + \frac{2\chi}{3} \frac{M_\infty^3}{r^3}$$
$$g_{rr}(r) \approx \left(1 - \frac{2M_\infty}{r} \right)^{-1} - 4\pi\mu_1 \frac{M_\infty^3}{r^3}$$

where $\chi = \pi_2 - \mu_2 - 2\pi\mu_1$

post-TOV: astrophysics




Surface redshift (I)

- The first “observable” we can construct is surface redshift (of absorption & emission lines). This is defined in the usual way:

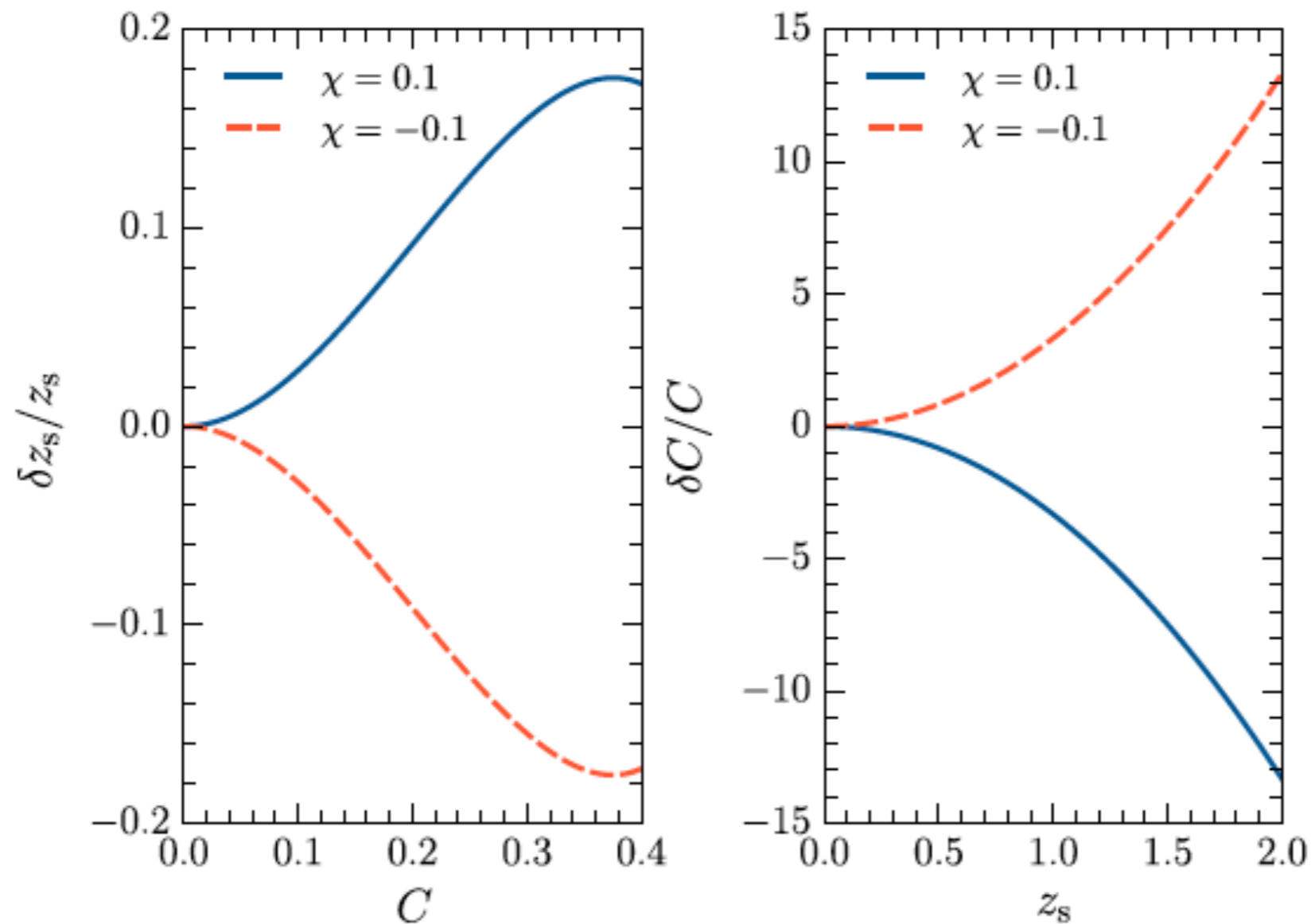
$$z_s = \frac{f_s}{f_\infty} - 1 \quad \text{compactness } C = M_\infty/R$$

- For any static spacetime:

$$\frac{f_\infty}{f_s} = \left[\frac{g_{tt}(R)}{g_{tt}(\infty)} \right]^{1/2} \Rightarrow \left\{ \begin{array}{l} z_s(C) = \underbrace{(1 - 2C)^{-1/2} - 1}_{\text{GR part}} + \frac{\chi}{3} C^3 \\ C(z_s) = \frac{1}{2} \underbrace{[1 - (1 + z_s)^{-2}]}_{\text{GR part}} \left(1 - \frac{\chi}{3} z_s^2 \right) \end{array} \right.$$


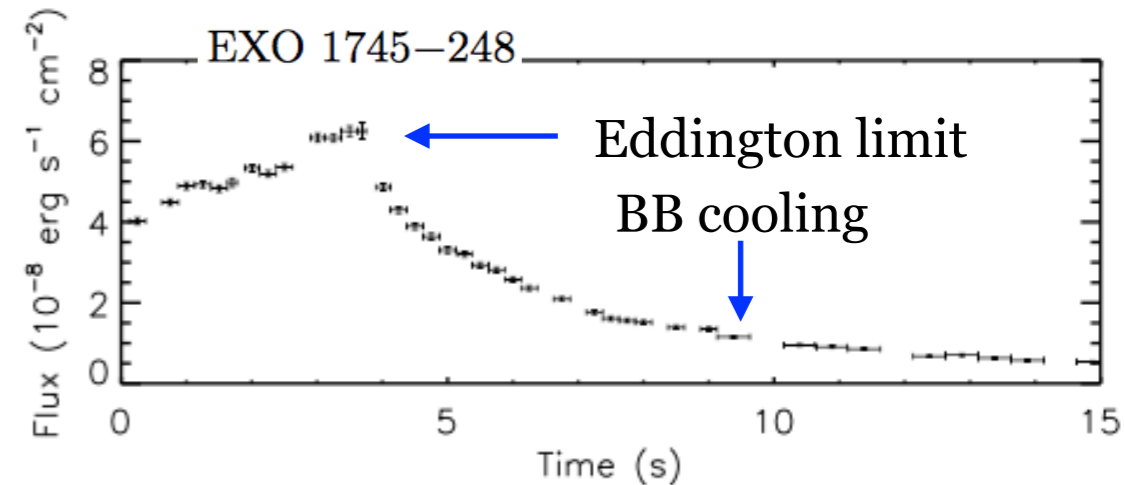
Surface redshift (II)

- The redshift depends *only* on $\chi = \pi_2 - \mu_2 - 2\pi\mu_1 \Rightarrow$ degeneracy!



Thermonuclear bursts (I)

- These are X-ray flashes produced by nuclear detonation of accreted matter on the surface of a neutron star.
- We follow the recipe of Psaltis (2008).



all relations as in GR

- The *observed* flux & apparent radius:

$$\left\{ \begin{array}{l} L_{\infty} = 4\pi D^2 F_{\infty} = \sigma_{\text{SB}} S_{\text{app}} \bar{T}_{\infty}^4 \\ R_{\text{app}} \equiv \left(\frac{S_{\text{app}}}{4\pi} \right)^{1/2} = D \left(\frac{F_{\infty}}{\sigma_{\text{SB}} \bar{T}_{\infty}^4} \right)^{1/2} \end{array} \right.$$

- Relation between the “colour” and effective BB temperature:

$$\underbrace{\bar{T}_{\infty}}_{\text{extracted from BB spectrum}} = f_c \underbrace{\sqrt{-g_{tt}(R)}}_{\text{“colour correction”}} T_{\text{eff}}$$

extracted from
BB spectrum

“colour
correction”

$$L_s = 4\pi R^2 \sigma_{\text{SB}} T_{\text{eff}}^4$$

conserved number of emitted photons: $L_{\infty} = -g_{tt}(R) L_s$

Thermonuclear bursts (II)

- The second key observable is the Eddington flux (“touchdown luminosity”):

$$L_E^\infty = 4\pi D^2 F_E^\infty = \frac{4\pi R^2}{\underbrace{\kappa (1+z_s)^2}_{\substack{\text{opacity of matter} \\ \text{(Thomson scattering)}}}} g_{\text{eff}} \Rightarrow g_{\text{eff}} = \kappa \sigma_{\text{SB}} \frac{F_E^\infty}{F_\infty} \left(\frac{\bar{T}_\infty}{f_c} \right)^4 (1+z_s)^4$$

- *Only the effective surface “g” takes a non-GR form:*

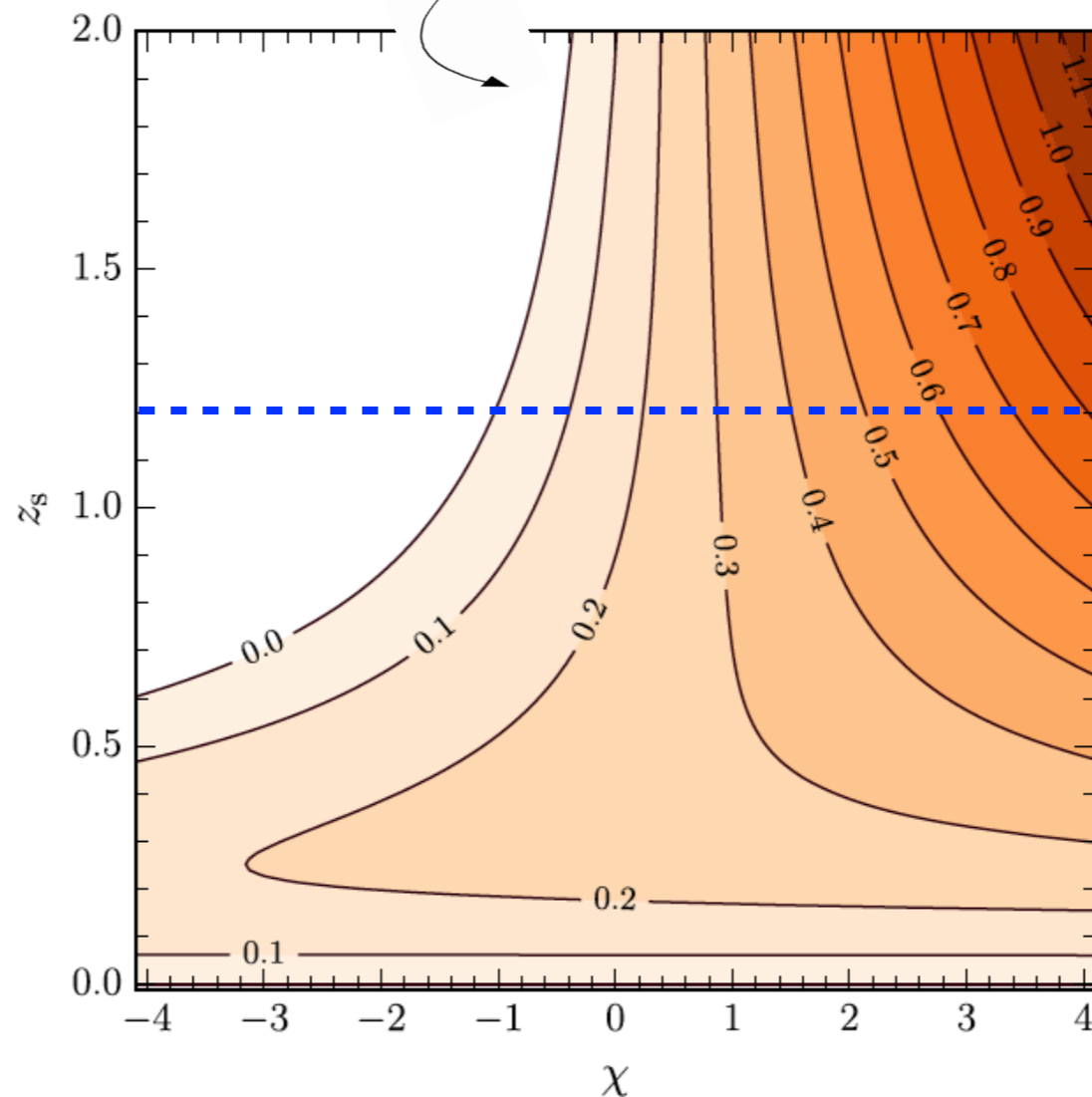
$$g_{\text{eff}} = \frac{1}{2\sqrt{g_{rr}(R)}} \frac{g'_{tt}(R)}{g_{tt}(R)} \Rightarrow g_{\text{eff}} = \frac{z_s (2+z_s)}{2R (1+z_s)} \left(1 + \frac{2}{3} \chi z_s^2 \right)$$

- Combine the above equations to produce a relation:

$$f(\chi, z_s) = \text{observables}$$

X-ray bursts (III)

$$\underbrace{\frac{z_s(2+z_s)}{(1+z_s)^4} \left(1 + \frac{2}{3}\chi z_s^2\right)}_{\text{contour plot}} = 2\kappa D \frac{F_E^\infty}{f_c^2} \left(\frac{\sigma_{\text{SB}} \bar{T}_\infty^4}{F_\infty}\right)^{1/2}$$



measure redshift
to get χ !

Geodesics & QPOs (I)

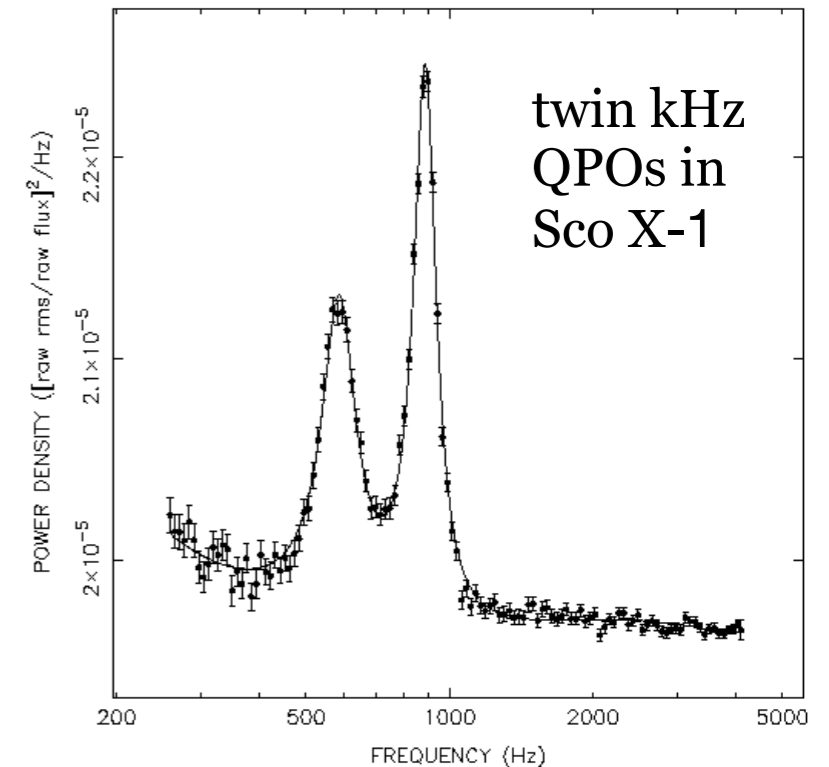
- In the most popular models, QPOs from accreting systems are associated with geodesic frequencies (the reality, the situation may not be that simple!)
- Azimuthal frequency :

$$\Omega_{\varphi}^2 = -\frac{g'_{tt}}{g'_{\varphi\varphi}} \approx \frac{M_{\infty}}{r^3} \left(1 + \chi \frac{M_{\infty}^2}{r^2} \right)$$

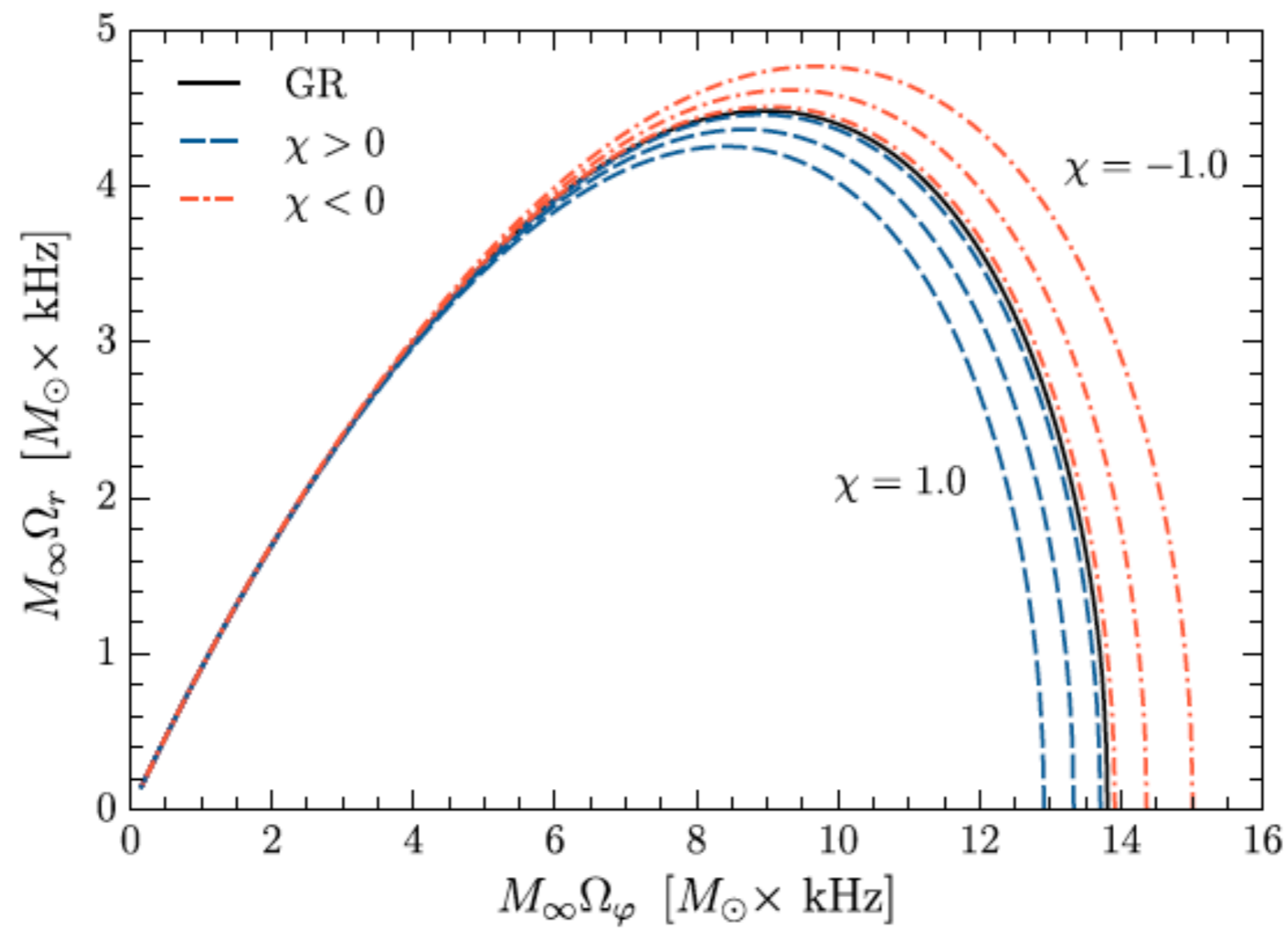
- Radial (epicyclic) frequency:

$$\Omega_r^2 = -\frac{g^{rr}}{2\dot{t}^2} V_{\text{eff}}''(r) \approx \frac{M_{\infty}}{r^3} \left(1 - \frac{6M_{\infty}}{r} - \chi \frac{M_{\infty}^2}{r^2} \right)$$

- ISCO radius: $r_{\text{isco}} \approx 6M_{\infty} \left(1 + \frac{19}{324}\chi \right)$



Geodesics & QPOs (II)



post-TOV: how general is it?

- Modified theories of gravity typically include one (or, less frequently, more) additional dynamical degrees of freedom (e.g. a scalar field)
- Our post-TOV formalism, featuring only two equations for dp/dr , dm/dr , implies that the extra d.o.f. ψ can be expressed in terms of the matter variables, i.e. $\psi = \psi(p, \rho, \Pi)$, *when building a NS*.
- This is the case for scalar-tensor theory and the same is likely true for other theories with a single extra d.o.f.
- By construction, the post-TOV scheme assumes “small” departures from GR, so it may *not* capture non-perturbative effects like “spontaneous scalarization”.

Outlook

- The post-TOV formalism is a toolkit for building relativistic stellar models with small/moderate departures from GR.
- Its parametrised form should (eventually) encompass a large class of modified theories of gravity.
- **Plenty of extensions:**
 - Include *dimensional* coupling constants (in progress): more post-TOV terms but families still exist.
 - Map formalism onto various alternative theories. Capture non-linear effects (e.g. scalarisation).
 - Add *slow rotation*: necessary step for a realistic framework.
 - Study I-Q “universal” relations (and perhaps break matter-gravity degeneracy?)