

Critical Phenomena in Gravitational Collapse

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GRECO Seminar, IAP, Paris

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Summary

- 1 Historical introduction
 - Christodoulou
 - Choptuik
 - Other results
- 2 The current model
 - Self-similarity
 - GR as a dynamical system
 - Mass scaling
- 3 Interesting results
 - Non-spherical systems
 - Global structure of the critical solution
 - Chaos
- 4 Conclusions and open questions

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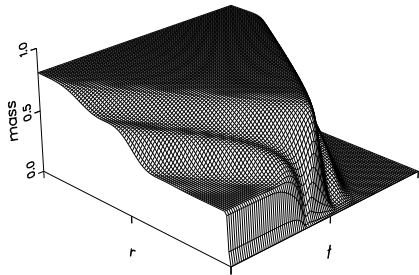
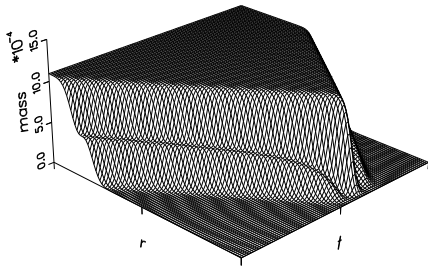
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- Goldwirth and Piran, PRD'87:

We present a numerical study of the gravitational collapse of a massless scalar field. We calculate the future evolution of new initial data, suggested by Christodoulou, and we show that in spite of the original expectations these data lead only to singularities engulfed by an event horizon.

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- And then in 1993...



1.2. Choptuik's setup

- The system:

$$ds^2 = -\alpha^2(t, r)dt^2 + a^2(t, r)dr^2 + r^2d\Omega^2, \quad \Phi \equiv \phi', \quad \Pi \equiv a\dot{\phi}/\alpha$$

$$\dot{\Phi} = \left(\frac{\alpha}{a}\Pi\right)', \quad \dot{\Pi} = \frac{1}{r^2} \left(r^2\frac{\alpha}{a}\Phi\right)', \quad \frac{\alpha'}{\alpha} = \frac{a'}{a} + \frac{a^2 - 1}{r} = 2\pi r(\Pi^2 + \Phi^2).$$

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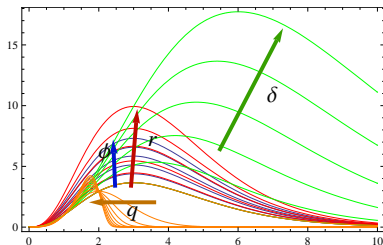
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- One-parameter (ρ) families of initial conditions with the property:
 - Small ρ leads to no BH formation (small finite data).
 - Large ρ produces a BH (large data).

Example (pure ingoing):

$$\phi(0, r) = \phi_0 r^3 \exp(-[(r - r_0)/\delta]^q)$$

$$\rho = \phi_0, r_0, \delta, q$$



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Comment: Self-similarity is dynamically found, but in a more general form!

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- Abrahams & Evans (PRL'93): axisymmetric vacuum (DSS).
- Evans & Coleman (PRL'94): perfect fluid, $\rho = \rho/3$ (CSS).
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- Proca, Dirac, sigma fields, ..., Vlasov(?)
- With/without mass, charge, conformal couplings, ...
- Different equations of state for fluids.
- Other dimensions.

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- New bet!

Whereas Stephen W. Hawking (having lost a previous bet on this subject by not demanding genericity) still firmly believes that naked singularities are an anathema and should be prohibited by the laws of classical physics,

And whereas John Preskill and Kip Thorne (having won the previous bet) still regard naked singularities as quantum gravitational objects that might exist, unclothed by horizons, for all the Universe to see,

Therefore Hawking offers, and Preskill/Thorne accept, a wager that

When any form of classical matter or field that is incapable of becoming singular in flat spacetime is coupled to general relativity via the classical Einstein equations, then

A dynamical evolution from generic initial conditions (i.e., from an open set of initial data) can never produce a naked singularity (a past-incomplete null geodesic from \mathcal{I}_+).

The loser will reward the winner with clothing to cover the winner's nakedness. The clothing is to be embroidered with a suitable, truly concessionary message.

Stephen W. Hawking

John P. Preskill Kip S. Thorne
John P. Preskill & Kip S. Thorne

Pasadena, California, 5 February 1997

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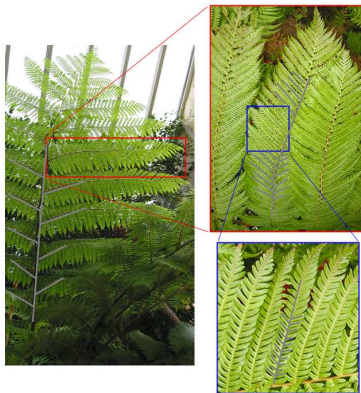
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- Any continuous symmetry has a discrete version.



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- CSS: Homothetic Killing vector ξ^a :

$$\mathcal{L}_\xi g_{ab} = -2g_{ab}$$

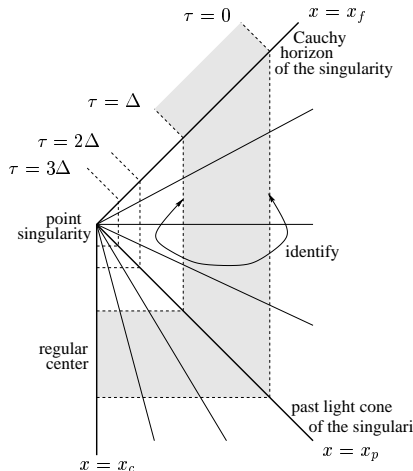
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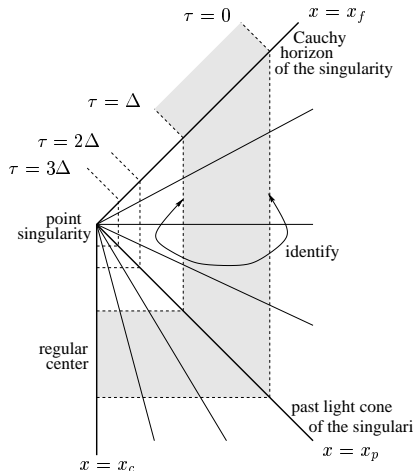
$$x \equiv \frac{r}{-t}, \quad \tau \equiv -\log \frac{-t}{t_0}$$

- Then any metric is:

$$e^{-2\tau} (Ad\tau^2 + 2Bd\tau dx + Cdx^2 + Fd\Omega^2)$$

with $\xi = \partial_\tau$.

- CSS: A, B, C, F functions of x only.
DSS: also periodic in τ , period Δ .



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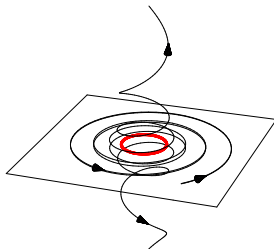
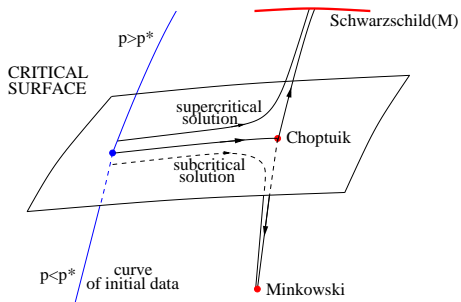
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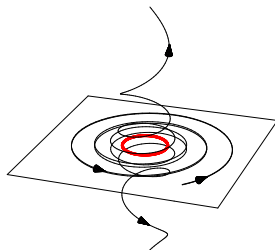
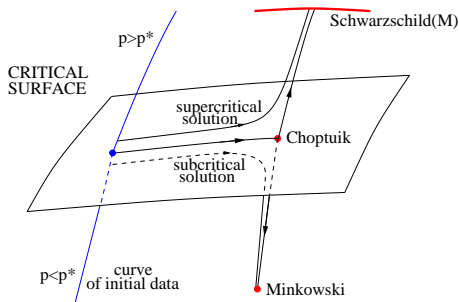
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- Warnings:
 - Which functional space? Asymptotic properties of the spacetimes.
 - Which foliations? Which coordinates?
 - Meaning of "attraction"?

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- Attraction \Rightarrow Forget initial details \Rightarrow Highly symmetric solutions:
 - Spherical or axisymmetric
 - Static (“type I”) or self-similar (“type II”). Both continuous or discrete.
- Two approaches to criticality:
 - The nonlinear way:
 - Evolution code and fine tune IC families to critical surface.
 - Search for critical phenomena (mainly universality).
 - The linear way:
 - Construct candidate critical solution.
 - Check there is a unique unstable linear mode.

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$$K(p - p^*) e^{\lambda_0 \tau_p} s_0(x) \approx \epsilon$$

- At that time τ_p the system has forgotten everything except for the scale

$$(-t_p) = t_0 e^{-\tau_p} \propto (p - p^*)^{1/\lambda_0}, \quad M_{BH} \propto (p - p^*)^{1/\lambda_0}, \quad \max R \propto (p - p^*)^{-2/\lambda_0}$$

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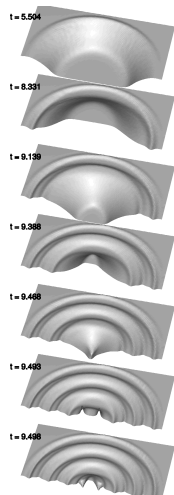
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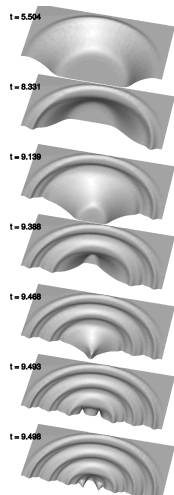
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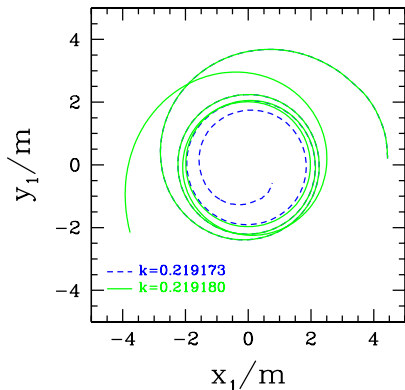
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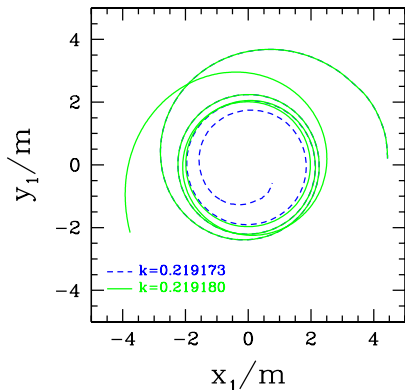
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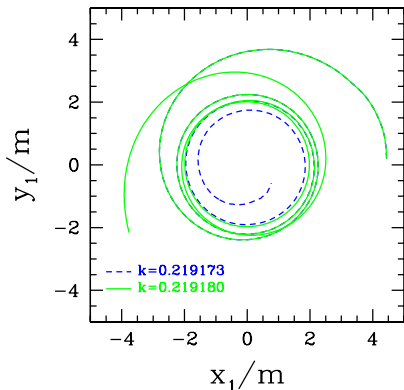
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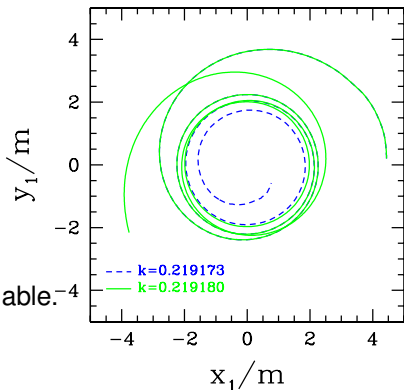
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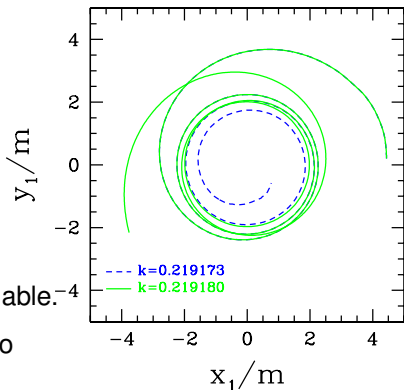
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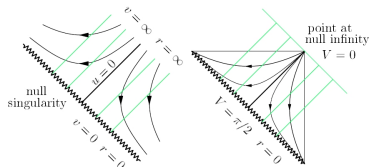
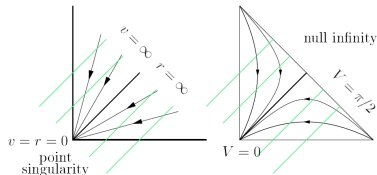
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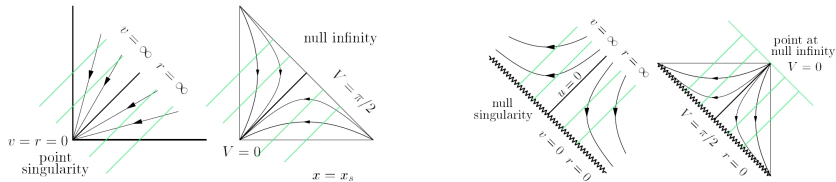
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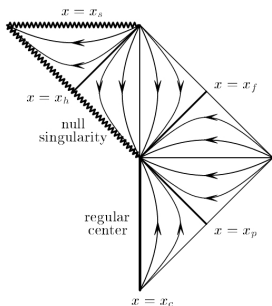


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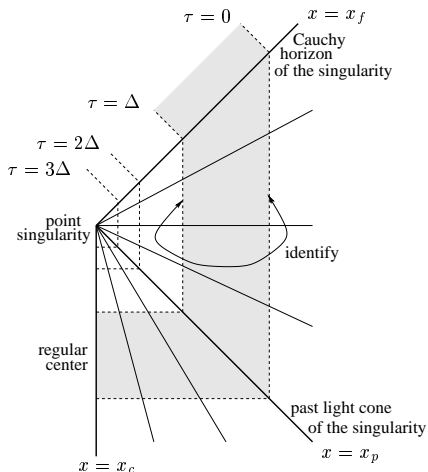
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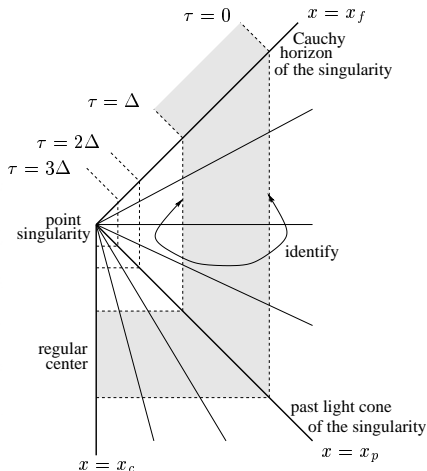
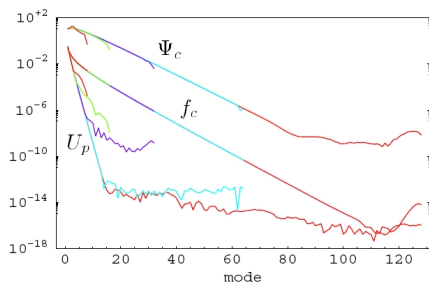
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- Three regions
- Pseudospectral code. Fourier in τ ;
4th order FD in x .



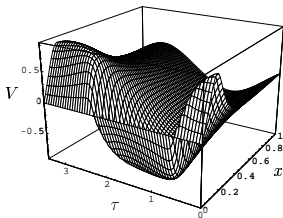
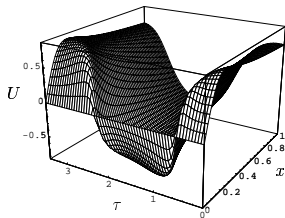
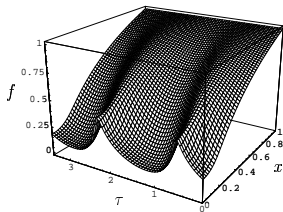
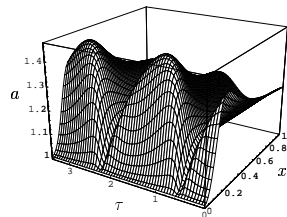
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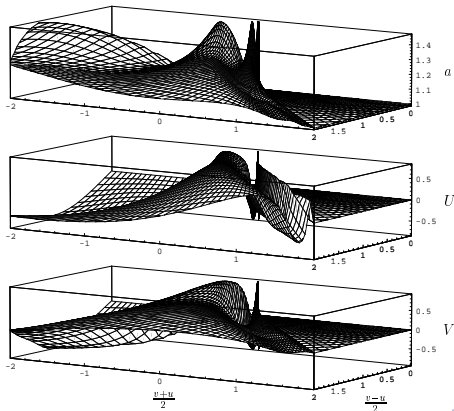
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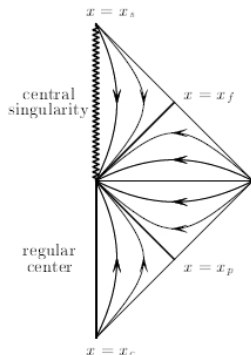
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- Unique DSS continuation with regular center (nearly flat):



3.2. Global structure of the Choptuik spacetime

All other continuations produce a negative mass singularity at the centre, with no new self-similarity horizon:



4+1 vacuum collapse

Bizoń et al PRD'05, PRL'05, PRL'06

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Szybka & Chmaj PRL'08

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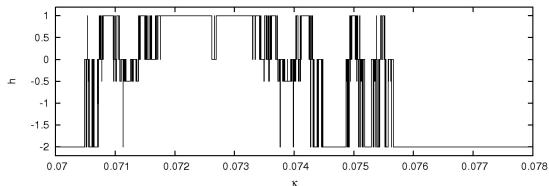
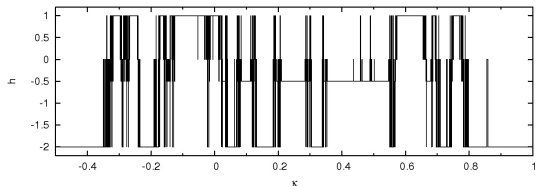
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