

# The Role of Dust in the Early Universe

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■ references:

I: Protogalaxy Evolution ( DY+ 2011, ApJ, 735, 44 )

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( DY+ 2011, ApJ, 735, 44 )

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## INTRODUCTION

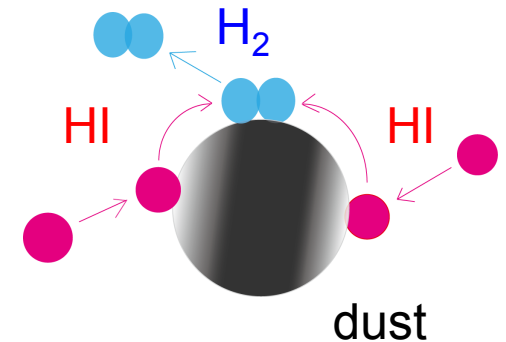
# Role of Dust in Star Formation

■  $H_2$  molecule is dominant coolant in the early Universe :

- Stars form in cool dense gas that is formed by  $H_2$  cooling.

■  $H_2$  formation on dust is very quick than gas phase :

- Dust is catalyst in  $H_2$  formation.
- $H_2$  formation is due to collisions of hydrogen atoms to dust surface.
- The collision rate depends on dust size, dust-to-gas mass ratio,  $T_{\text{gas}}$ , and  $\rho_{\text{gas}}$ .

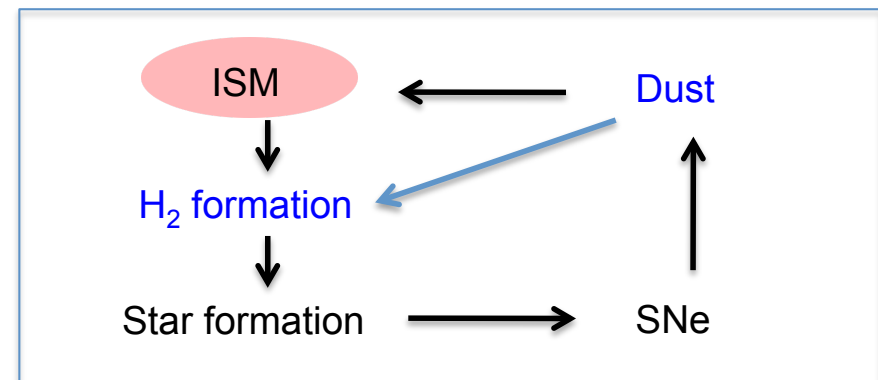


■ Dust size distribution in high redshift ( $z > 5$ ) :

- This is different from Milky Way dust.
- This is because dominant dust origin is SNe.

# Our Study

- We investigate time evolution of dust mass and **dust size distribution** in high  $z$  Universe, considering the dust production by SNe and **dust destruction** by sputtering in the high-velocity shocks driven by SNe.
- We consistently treat following processes in our one-zone model
  - (i) the formation and **size evolution** of dust,
  - (ii) the chemical reaction networks including  $H_2$  formation **both** on the dust grains and in gas phase,
  - (iii) gas cooling and heating
  - (iv) the **SFR based on  $H_2$  mass**.



# MODEL

## Dust Evolution (1/2)

### ■ Dust in the early Universe :

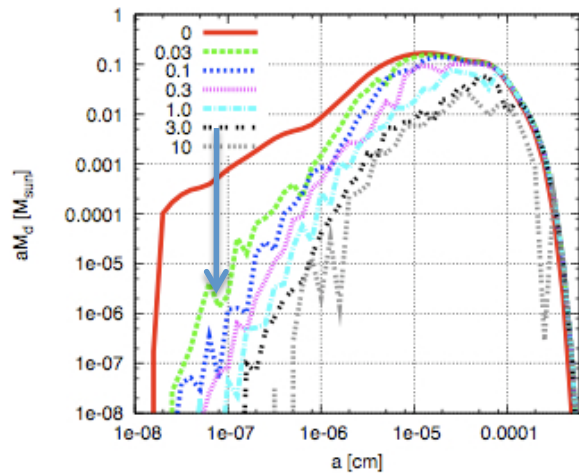
- SNe II are the source of dust in the early Universe (  $5 < z < 10$  )  
( Dwek+ 07; Gall+ 2011a, b, c; but see also Valiante+09; Dwek&Cherchneff 11; Valiante+11 )
- Average size of SN dust is small.

### ■ Dust injected from SN II into ISM through a reverse shock :

A reverse shock destroys dust.

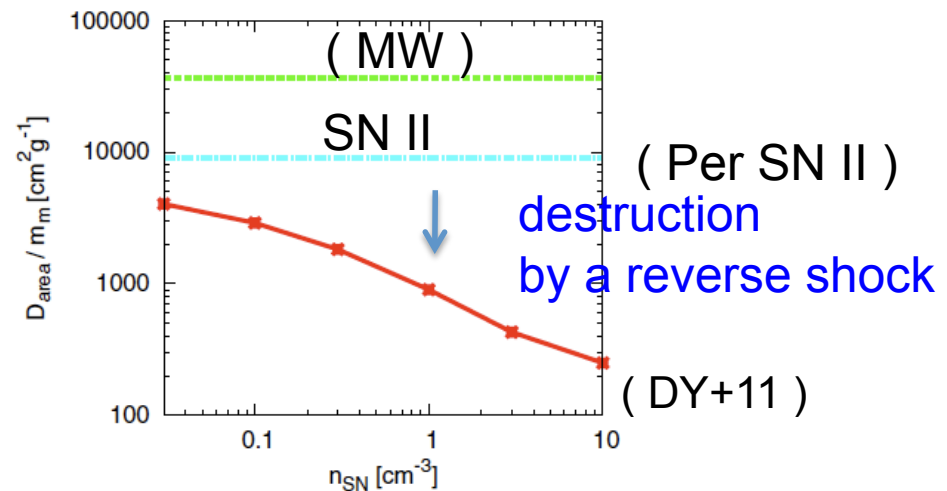
This is more efficient in large  $n_{SN}$ .

$$D_{area} = 4\pi a^2 f(a) da$$



adopt:  
the models by  
Nozawa+ 07

( DY+11 )



$n_{SN}$  : gas density around SN progenitor

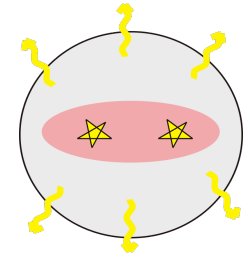


## MODEL

# Galaxy One-zone Model (1/2)

### ■ Dark matter halo and physical state of gas :

- We assume (i) the DM halo as a **singular isothermal sphere** and (ii) the baryonic gas as a **uniform, rotating gas disk**



- Initial gas temperature :  $T = T_{\text{vir}}$   $T_{\text{vir}} \equiv \frac{G\mu m_{\text{H}} M_{\text{vir}}}{3k_{\text{B}} r_{\text{vir}}}$ , (11)

- A rotation timescale,  $t_{\text{cir}}$  :  $t_{\text{cir}} \equiv \frac{2\pi r_{\text{disk}}}{v_{\text{c}}}$ .

$r_{\text{disk}}$  : the radius of disk  
( Mo+ 98,  $\lambda = 0.04$  )

- The number density of H,  $n_{\text{H}}$  :

H : the typical scale height of Tgas  
( Shakura & Sunyaev 88 )

$$n_{\text{H}} = \frac{M_{\text{H}}}{\pi r_{\text{disk}}^2 2H m_{\text{H}}} \quad (16)$$

the typical volume of the galaxy

$$H = \sqrt{2} \frac{v_{\text{s}}}{v_{\text{c}}} r_{\text{disk}} = \left( \frac{2T}{3T_{\text{vir}}} \right)^{\frac{1}{2}} r_{\text{disk}}$$

### ■ A star formation law based on H<sub>2</sub> :

$$\Psi(t) = \frac{f_{\text{H}_2}(t) M_{\text{H}}(t)}{t_{\text{cir}}(z_{\text{vir}})}, \quad (18)$$

$f_{\text{H}_2}$  : the molecular fraction

$f_{\text{H}_2} \times M_{\text{H}}$  : the mass of H<sub>2</sub> in the galaxy

# MODEL

## Galaxy One-zone Model (2/2)

### Chemistry and cooling :

- the time evolution of H<sub>2</sub> fraction, f<sub>H<sub>2</sub></sub> :

$$\frac{df_{H_2}}{dt} = \left[ \frac{df_{H_2}}{dt} \right]_{\text{gas}} + \left[ \frac{df_{H_2}}{dt} \right]_{\text{dust}} + \left[ \frac{df_{H_2}}{dt} \right]_{\text{dest}} + \left[ \frac{df_{H_2}}{dt} \right]_{\text{UV}} + \left[ \frac{df_{H_2}}{dt} \right]_{\text{star}}, \quad (22)$$

### H<sub>2</sub> formation on dust grains

- cooling ( H<sub>2</sub>, H, H<sup>+</sup>, C<sub>I</sub>, C<sub>II</sub>, O<sub>I</sub> ) and heating ( photoheating )

### Formation of molecular hydrogen on dust grains :

$$\left[ \frac{df_{H_2}}{dt} \right]_{\text{dust}} = 2R_{\text{dust}} \mathcal{D} n_H f_0$$

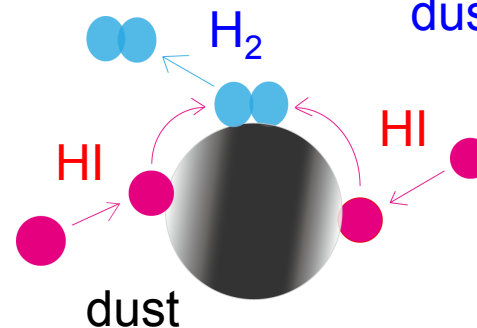
$$= \sum_j \int_0^\infty \underline{f_0 f_j(a) \pi a^2 \bar{v} S da}$$

the collision rate (30)  
between the dust grains  
and HI atoms

$$R_{\text{dust}}(a) \mathcal{D} = \sum_j \int_0^\infty \left( \frac{3m_H \bar{v} S}{8a\rho_j} \right) \left( \frac{4\pi a^3 \rho_j f_j(a)}{3n_H m_H} \right) da.$$

a : radius of dust ↑ (33)

dust-to-gas mass ratio



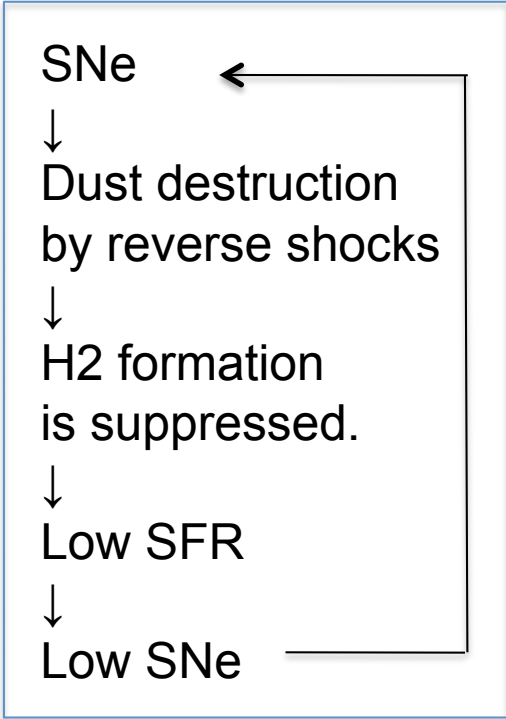
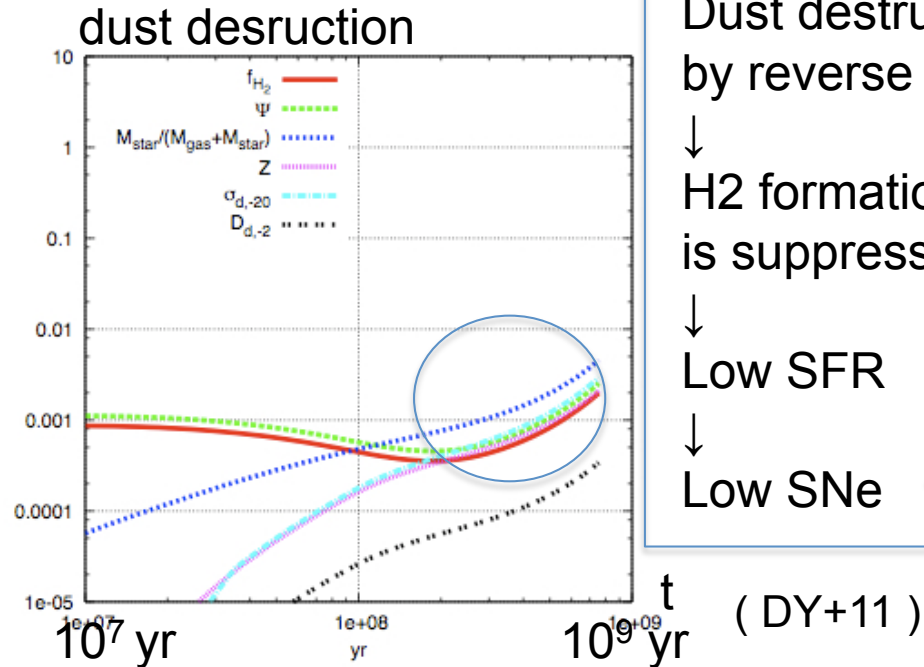
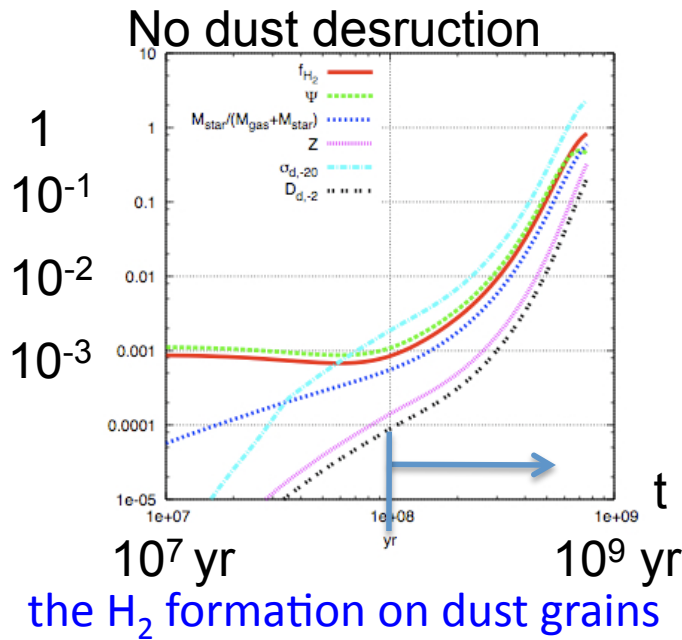


# RESULTS

## Galaxy Evolution

### Time evolution of galaxy on dust models :

- $M_{\text{vir}} = 10^9 M_{\text{sun}}$ ,  $z_{\text{vir}} = 10$ ,  $n_{\text{SN}} = 1 \text{ cm}^{-3}$  :



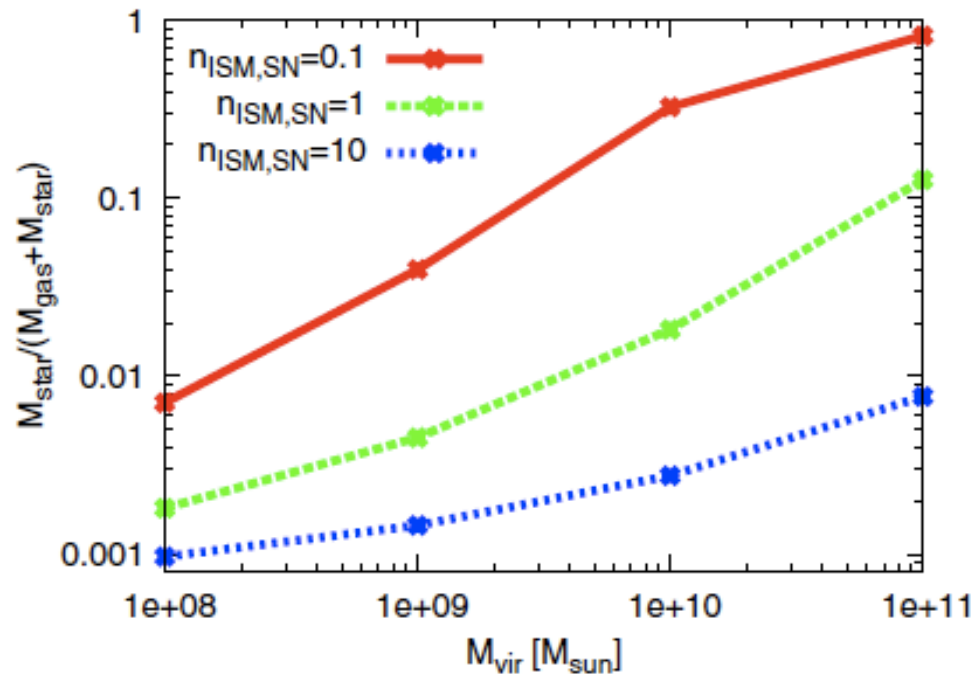
reverse	×	○
forward	×	○

- The dust destruction by the reverse shocks is very effective and suppresses  $\text{H}_2$  formation on dust grains.

## RESULTS

# SFE vs. DM Halo Mass

- SFR is very suppressed in our model with dust destruction, and stellar mass fraction is very low at  $t \sim 0.8$  Gyr.  
note : SFE can be well characterized by stellar mass fraction.



- SFE is large for large  $M_{\text{vir}}$  :  
The gas cools the  $T_{\text{CMB}} \sim 40$  K in all  $M_{\text{vir}}$ .  
 $T_{\text{vir}}$  increase with  $M_{\text{vir}}$ .  
↓  
The final gas density becomes higher because of smaller  $H / r_{\text{disc}}$ .  
↓  
This results in more rapid  $\text{H}_2$  formation.  
↓  
This enhances the SFR.

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(ApJ in this July)

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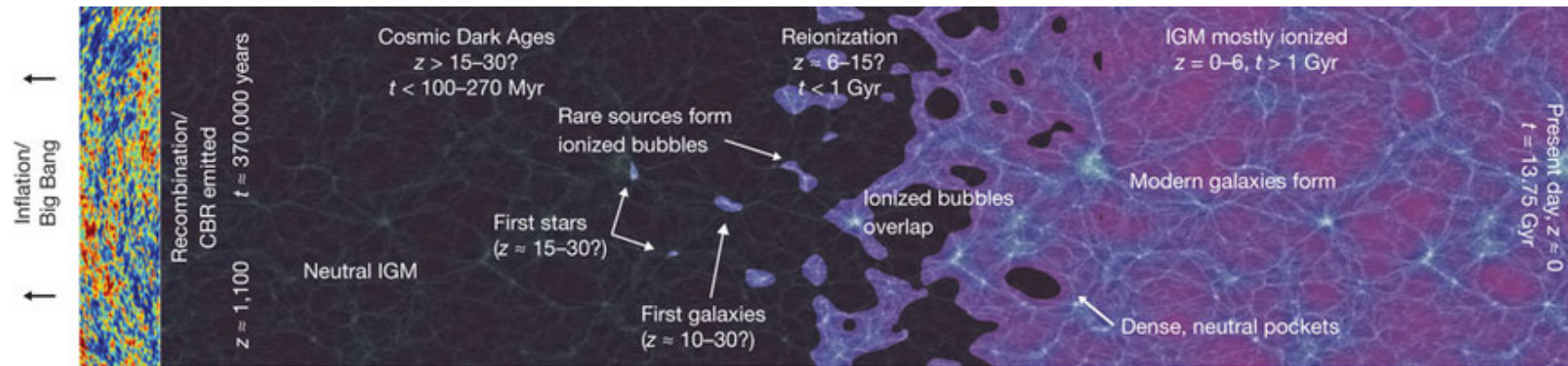
RESULTS

The Effects of  $D_{\text{crit}}$  on Reionization

## III: SUMMARY

# INTRODUCTION

## Reionization in the Universe



( Robertson+ 10 )

### ■ Observational constraints on reionization :

- Gun-Perterson test : reionization for  $z > 6$

### ■ Reionizing photons :

- Most of the reionization radiation is expected to come from galaxies less than  $\sim 10^{9.5}$  Msun.

### ■ Star formation efficiency ( SFE ) :

- The stellar mass fraction,  $M_{\text{star}} / (M_{\text{gas}} + M_{\text{star}})$ , is assumed to be constant independent of  $M_{\text{vir}}$  in previous analytic works.

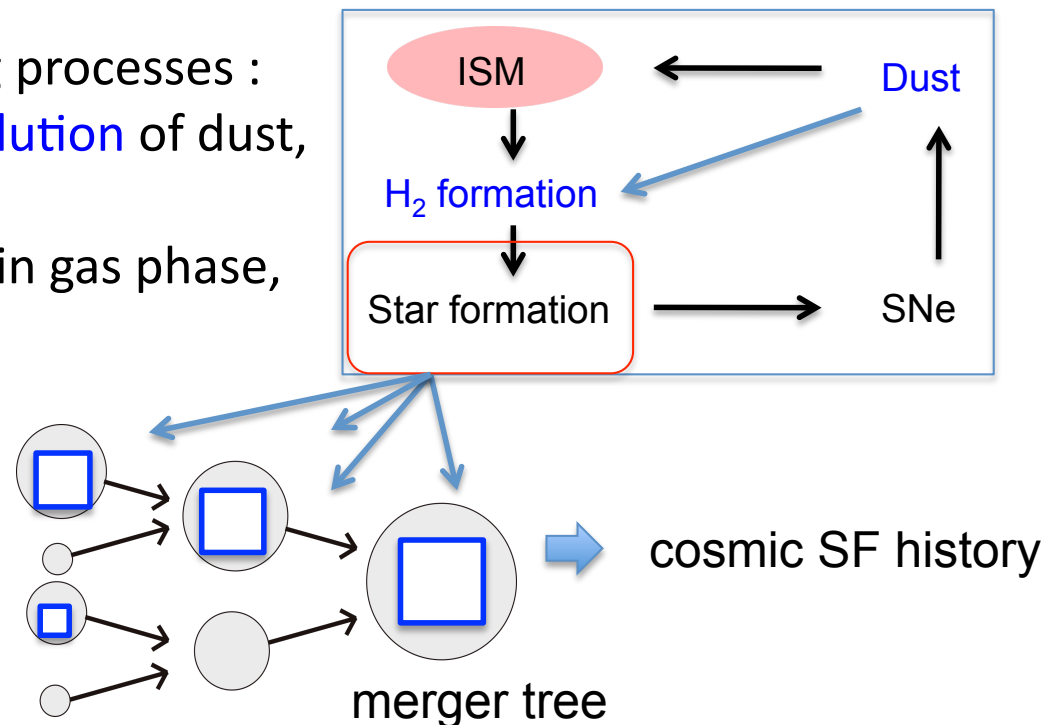
# Our Study

■ We study reionization by our one-zone protogalaxy model which includes dust production and dust size evolution.

■ Because of  $H_2$  formation rate on dust grains decreasing, SFE decreases with small dark matter halo mass.

■ We consistently treat following processes :

- (i) the formation and **size evolution** of dust,
- (ii)  $H_2$  formation
- both** on the dust grains and in gas phase,
- (iii) gas cooling and heating
- (iv) the **SFR based on  $H_2$  mass** and
- (v) DM halo evolution
- (vi) IMF transition from Pop III to Pop II

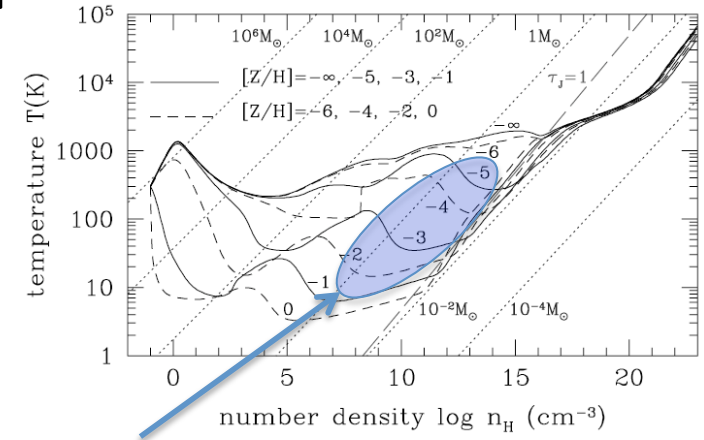


# MODEL

## Pop III.1 and Pop III.2

### ■ Critical dust-to-gas mass ratio, $D_{\text{crit}}$ :

- We assume the IMF transition from Pop III to Pop II due to dust cooling.  
( Schneider+ 03; Omukai+ 05; Schneider+ 06; Schneider and Omukai 10; Omukai+ 10 )



dust cooling

(Omukai+ 05)

### ■ IMF transition :

Dust	DM halo	IMF	SFR
$D < D_{\text{crit}}$	$T_{\text{vir}} < 10000 \text{ K}$	Pop III.1 ( <u>100 - 500</u> $M_{\text{sun}}$ )	( Machacek+ 03 )
		Pop III.2 ( <u>10 - 100</u> $M_{\text{sun}}$ )	$\Psi(t) = \frac{f_{\text{H}_2}(t) M_{\text{H}}(t)}{t_{\text{vir}}(z_{\text{vir}})}$
$D \geq D_{\text{crit}}$	$T_{\text{vir}} \geq 10000 \text{ K}$	Pop II ( <u>0.1 - 60</u> $M_{\text{sun}}$ )	( Our model )

all slope index : -2.35

## MODEL

# Reionization

- Ionizing photons,  $n_{\text{ion}}$ , emitted by massive stars :

$$\frac{1}{n_b} \frac{dn_{\text{ion}}(z)}{dz} = \frac{1}{\rho_m} \frac{\Omega_m}{\Omega_b} f_{\text{esc}} \eta_{\text{ion}} \Psi_*(z) \left| \frac{dt}{dz} \right|$$

cosmic SFR (per Mpc<sup>3</sup>)

$n_{\text{ion}}$  : comoving density of ionizing photons

$f_{\text{esc}}$  : escape fraction

$\eta_{\text{ion}}$  : number of ionizing photons  
emitted per stellar baryon

- ionizing photon number per stellar mass for Pop III.2 stars is **10 times** larger than that for Pop II stars (Schaerer 02).
- Pop III.2 stars are very effective to the reionization.

- Evolution of ionized volume fraction,  $Q_{\text{ion}}$  :

$$\frac{dQ_{\text{ion}}(z)}{dz} = \frac{1}{n_b} \frac{dn_{\text{ion}}(z)}{dz} - \alpha_B n_b C(z) Q_{\text{ion}}^2(z) (1+z)^3 \left| \frac{dt}{dz} \right|$$

ionization

recombination

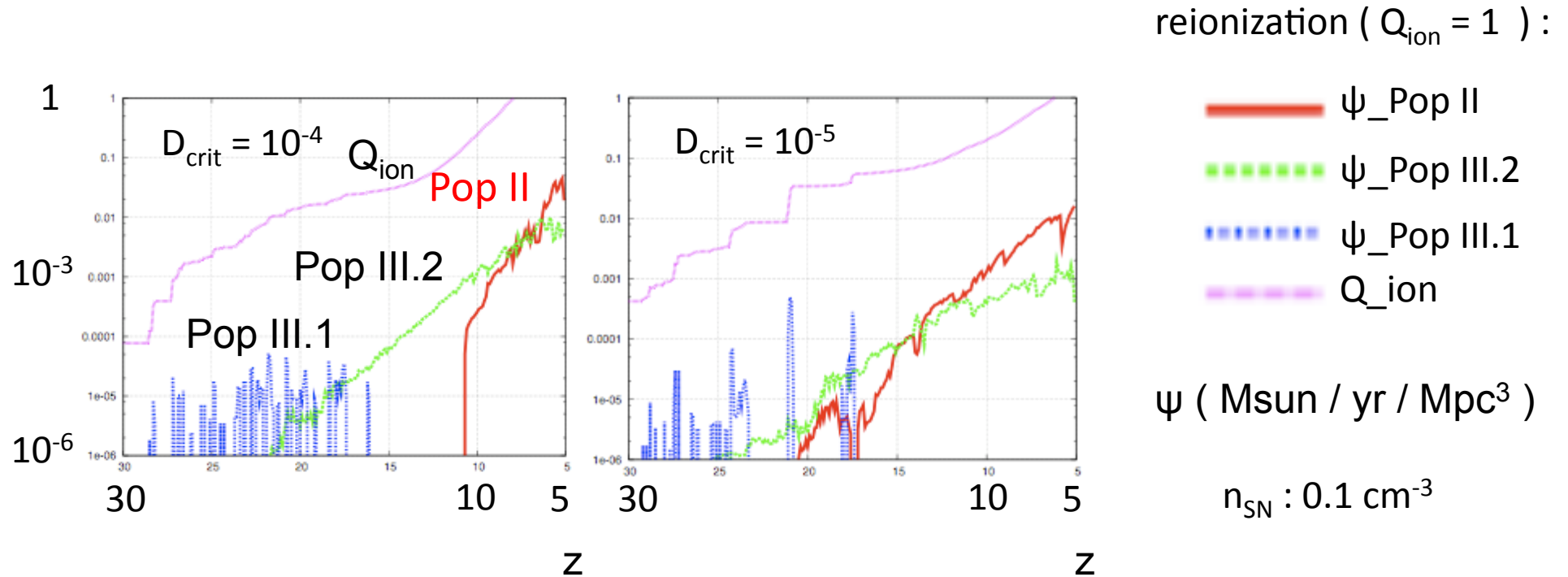
$Q_{\text{ion}}$  : volume fraction of ionizing regions

$\alpha_B$  : recombination coefficient

$C(z)$  : clumping factor

## RESULTS

# Reionization and Dcrit



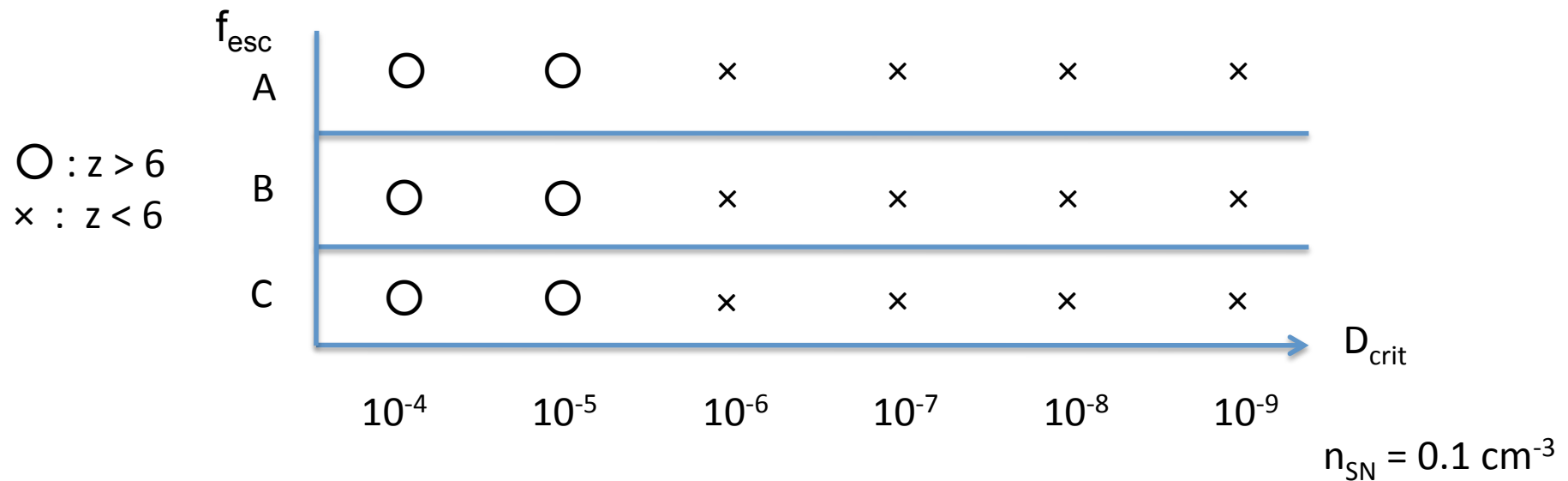
- Small  $D_{\text{crit}}$   $\rightarrow$  early transition from Pop III to Pop II  $\rightarrow$  late reionization epoch  
note: ionizing photon per Pop III mass  $>$  ionizing per Pop II mass



## RESULTS

# Reionization and $D_{\text{crit}}$

■ reionization (  $Q_{\text{ion}} = 1$  ) :



A : Pop III.2 :  $f_{\text{esc}} = 10\%$ , Pop II :  $f_{\text{esc}} = 30\%$  ( Greif and Bromm 06 )

B : Pop III.2 :  $f_{\text{esc}} = 50\%$ , Pop II :  $f_{\text{esc}} = 60\%$  ( Wise and Cen 09, cosmological )

C : Pop III.2 :  $f_{\text{esc}} = 70\%$ , Pop II :  $f_{\text{esc}} = 80\%$  ( Wise and Cen 09, isolated )

- In the cases of  $D_{\text{crit}} > 10^{-5}$  ,  
reionization occurs at  $z > 6$  independent of escape fraction,  $f_{\text{esc}}$  .

# SUMMARY

## I: Protogalaxy Evolution :

We conclude that

the amount and **the size distribution** of dust **strongly affects** the evolution of galaxies in the early Universe, since SF activity depends on H<sub>2</sub> formation on dust grains.

- We show that H<sub>2</sub> formation is suppressed by the dust destruction, especially by **the reverse shocks** in SNRs.

## II. Reionization :

We study reionization by our galaxy model

and find that, in the cases of  $D_{\text{crit}} > 10^{-5}$ ,

reionization occurs at  $z > 6$  independent of other parameters.

- This result show that study of critical dust-to-gas mass ratio is very important for reionization process.