Recalibrated waveforms for EMRIs in the EOB frame

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Outline

- Background and motivations;
- Numerical simulation for EMRIs;
- Effective-one-body formalism;
- Recalibrated waveforms in the EOB frame;
- Conclusions and prospect;

Extreme-mass-ratio inspirals: EMRIs



LISA



Chinese space-based projects: Taiji, Tianqin

eLISA pathfinder has gotten good results



Why studying EMRIs

- very long time scale
- SNR 20 up to z=0.7, dozens/year
- studying physics near horizon of SMBH
- testing gravitation theory
- Cosmology

Extreme Mass Ratio Inspirals

- SNR 20 up to $z \approx 0.7$ for $10^5 10^6$ M_{\odot}
- Dozens of events per year
- Mass, spin to 0.1% 0.01 %
- Quadrupole moment to < $0.001 M_{\odot}^{3}G^{2}/c^{4}$
- Do Black Holes have hair?
 - New objects in General Relativity
 - Boson Stars, Gravastars, non-Kerr solutions (e.g. Manko-Novikov)
 - Deviations from General Relativity
 - Chern-Simons, Scalar-Tensor, light scalar fields (axions) and black hole bomb instabilities
- Each has specific GW fingerprint! From Danzmann, 2017 May 25, Beijing



Challenge:

Firstly, we should have huge numbers of waveform templates of EMRIs with high accuracy. The role of waveform templates Matched filtering: find GWs from noise



Recognize the parameters of binaries

EMRIs : Calculation method of waveforms



Teukolsky equation

$$\begin{split} &\left[\frac{\left(r^{2}+a^{2}\right)^{2}}{\Delta}-a^{2}\sin^{2}\theta\right]\frac{\partial^{2}\Psi}{\partial t^{2}}+\frac{4Mar}{\Delta}\frac{\partial^{2}\Psi}{\partial t\partial\phi}+\left[\frac{a^{2}}{\Delta}-\frac{1}{\sin^{2}\theta}\right]\frac{\partial^{2}\Psi}{\partial\phi^{2}}-\\ &-\Delta^{-s}\frac{\partial}{\partial r}\left(\Delta^{s+1}\frac{\partial\Psi}{\partial r}\right)-\frac{1}{\sin\theta}\frac{\partial}{\partial\theta}\left(\sin\theta\frac{\partial\Psi}{\partial\theta}\right)-2s\left[\frac{a\left(r-M\right)}{\Delta}+\frac{i\cos\theta}{\sin^{2}\theta}\right]\frac{\partial\Psi}{\partial\phi}-\\ &-2s\left[\frac{M\left(r^{2}-a^{2}\right)}{\Delta}-r-ia\cos\theta\right]\frac{\partial\Psi}{\partial t}+\left(s^{2}\cot^{2}\theta-s\right)\Psi=0\,. \end{split}$$

energy flux, waveform

$$\begin{split} \frac{dE}{dt} &= \lim_{r \to \infty} \left[\frac{1}{4\pi} \int_{\Omega} \left| \int_{-\infty}^{t} \psi d\tilde{t} \right|^{2} d\Omega \right], \\ \frac{dP_{i}}{dt} &= \lim_{r \to \infty} \left[\frac{1}{4\pi} \int_{\Omega} l_{i} \left| \int_{-\infty}^{t} \psi d\tilde{t} \right|^{2} d\Omega \right], \quad \psi \approx \frac{1}{2} \left(\frac{\partial^{2} h_{+}}{\partial t^{2}} - i \frac{\partial^{2} h_{\times}}{\partial t^{2}} \right) \\ \frac{dL}{dt} &= -\lim_{r \to \infty} \left\{ \frac{1}{4\pi} \operatorname{Re} \left[\int_{\Omega} \left(\partial_{\phi} \int_{-\infty}^{t} \psi d\tilde{t} \right) \left(\int_{-\infty}^{t} \int_{-\infty}^{t} \bar{\psi} d\tilde{t} d\hat{t} \right) d\Omega \right] \right\} \end{split}$$

Describe the orbit of small body

Effective-one-body (EOB) dynamics

Buonanno & Damour, 1999, 2000,

Two-body problem

one-body problem $M = m_1 + m_2$

 $\mu = m_1 m_2 / M$

 $m_1 \quad m_2$



EOB formalism: dynamics

$$\begin{split} \dot{r} &= \frac{\partial H_{\rm EOB}}{\partial p_r}, & H_{\rm EOB} = M\sqrt{1 + 2\nu(H_{\rm eff}/\mu - 1)} \\ \dot{\phi} &= \frac{\partial H_{\rm EOB}}{\partial p_{\phi}}, & H_{\rm eff} = H_{\rm NS} + H_{\rm S} - \frac{\mu}{2Mr^3}S_*^2 \\ H_{\rm NS} &= \beta^i p_i + \alpha\sqrt{\mu^2 + \gamma^{ij} p_i p_j}, \\ \dot{p}_r &= -\frac{\partial H_{\rm EOB}}{\partial r} + \mathcal{F}_{\phi} \frac{p_r}{p_{\phi}}, & H_{\rm S} = H_{\rm SO} + H_{\rm SS} \\ \dot{p}_{\phi} &= \overbrace{\mathcal{F}_{\phi}}^{\bullet} \text{ radiation reaction} \\ \alpha &= \frac{1}{\sqrt{-g^{tt}}}, & \hat{\mathcal{F}} = \frac{-1}{\nu\hat{\Omega}|r \times p|} \frac{dE}{dt}p, \\ \beta^i &= \frac{g^{ti}}{g^{tt}}, & \frac{dE}{dt} = \frac{\hat{\Omega}^2}{8\pi} \sum_{\ell=2}^8 \sum_{m=0}^\ell m^2 \left| \frac{\mathcal{R}}{M} \right| \\ \gamma^{ij} &= g^{ij} - \frac{g^{ti}g^{tj}}{g^{tt}}, \end{split}$$

$$\frac{dE}{dt} = \frac{\hat{\Omega}^2}{8\pi} \sum_{\ell=2}^8 \sum_{m=0}^\ell m^2 \left| \frac{\mathcal{R}}{M} h_{\ell m} \right|^2$$

EOB formalism: waveforms of circular orbits

$$h_{\ell m} = h_{\ell m}^{\text{insp-plunge}} \theta(t_{\text{match}}^{\ell m} - t) + h_{\ell m}^{\text{merger-RD}} \theta(t - t_{\text{match}}^{\ell m})$$

$$h_{\ell m}^{\text{insp-plunge}} = h_{\ell m}^{\text{F}} N_{\ell m}$$
$$h_{\ell m}^{\text{F}} = h_{\ell m}^{(N,\epsilon)} \hat{S}_{\text{eff}}^{(\epsilon)} T_{\ell m} e^{i\delta_{\ell m}} (\rho_{\ell m})^{\ell}$$

factorized waveform Pan et al., PRD 84, 124052 (2011)

$$h_{\ell m}^{\text{merger}-\text{RD}}(t) = \sum_{n=0}^{N-1} A_{\ell m n} e^{-i\sigma_{\ell m n}(t-t_{\text{match}}^{\ell m})}$$

$$\rho_{22} = 1 + \left(\frac{55\nu}{84} - \frac{43}{42}\right)\nu_{\Omega}^{2} + \left(\frac{19583\nu^{2}}{42336} - \frac{33025\nu}{21168} - \frac{20555}{10584}\right)\nu_{\Omega}^{4} + \left(\frac{10620745\nu^{3}}{39118464} - \frac{6292061\nu^{2}}{3259872} + \frac{41\pi^{2}\nu}{192} - \frac{48993925\nu}{9779616} - \frac{428 \operatorname{eulerlog}_{2}(\nu_{\Omega}^{2})}{105} + \frac{1556919113}{122245200}\right)\nu_{\Omega}^{6} + \left(\frac{9202 \operatorname{eulerlog}_{2}(\nu_{\Omega}^{2})}{2205} - \frac{387216563023}{160190110080}\right)\nu_{\Omega}^{8} + \left(\frac{439877 \operatorname{eulerlog}_{2}(\nu_{\Omega}^{2})}{55566} - \frac{16094530514677}{533967033600}\right)\nu_{\Omega}^{10}, \tag{B9a}$$

Accuracy of factorized PN waveform applying to EMRI



Accuracy requirement of EMRIs

- A typical EMRIs have M/μ cycles
- During the inspirals, for dephase less than 2pi, one asks

 $\Delta \dot{E}/\dot{E} < \mu/M$

- So for typical mass-ratio 10^5, the relative precision should be 10^-5.
- F-R PN waveforms break down even for circular orbits

Highly accurate and efficient waveforms

- Numerical simulations are inefficient;
- We try to work on semi-analytical models
- Yunes et al., 2011

$$\begin{split} \rho_{\text{Cal}}^{22} &= \rho^{22} + \left[a_{22}^{(9,1)} + b_{22}^{(9,1)} \text{ eulerlog}_2 v^2 \right] \,\bar{q} \, v^9 \\ &+ \left[a_{22}^{(12,0)} + b_{22}^{(12,0)} \text{ eulerlog}_2 v^2 \right] v^{12} \,, \\ \rho_{\text{Cal}}^{33} &= \rho^{33} + \left[a_{33}^{(8,2)} + b_{33}^{(8,2)} \text{ eulerlog}_3 v^2 \right] \, \bar{q}^2 \, v^8 \\ &+ \left[a_{33}^{(10,0)} + b_{33}^{(10,0)} \text{ eulerlog}_3 v^2 \right] v^{10} \,, \end{split}$$

Yunes' s results:



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$$\dot{E}^{\infty} = \sum_{i=0}^{n} a'_{i} x^{i}, \quad \dot{E}^{\mathrm{H}} = \sum_{i=0}^{n} b'_{i} x^{i},$$
$$\operatorname{Re}[H_{lm}] = \sum_{i=0}^{n} R'_{lm} x^{i}, \quad \operatorname{Im}[H_{lm}] = \sum_{i=0}^{n} I'_{lm} x^{i}.$$

Han, CQG, 2016

Table 1. polynomial parameters for infinity fluxes.

| a_0 | a_1 | a_2 | a_3 | a_4 | a_5 | a_6 | a_7 | a_8 | a_9 | a_{10} |
|---------|--|--|---|---|---|---|---|---|---|---|
| 1.51e-7 | -1.22e-5 | 4.40e-4 | -9.43e-3 | 1.36e-1 | 4.90e0 | -5.27e0 | -1.05e1 | 2.98e1 | -4.34e1 | 3.99e1 |
| 6.61e-7 | -5.98e-5 | 2.41e-3 | -5.69e-2 | 8.76e-1 | -2.90e0 | 5.22e1 | -2.75e2 | 8.81e2 | -1.64e3 | 1.44e3 |
| 1.46e-6 | -1.72e-4 | 9.08e-3 | -2.83e-1 | 5.76e0 | -7.39e1 | 7.63e2 | -5.02e3 | 2.21e4 | -5.77e4 | 7.19e4 |
| 1.26e-6 | -1.85e-4 | 1.22e-2 | -4.78e-1 | 1.22e1 | -2.09e2 | 2.63e3 | -2.21e4 | 1.25e5 | -4.20e5 | 6.78e5 |
| - | $\begin{array}{c} a_0 \\ \hline 1.51e-7 \\ \hline 6.61e-7 \\ \hline 1.46e-6 \\ \hline 1.26e-6 \end{array}$ | $\begin{array}{ccc} a_0 & a_1 \\ \hline 1.51e-7 & -1.22e-5 \\ \hline 6.61e-7 & -5.98e-5 \\ \hline 1.46e-6 & -1.72e-4 \\ \hline 1.26e-6 & -1.85e-4 \end{array}$ | $\begin{array}{c cccc} a_0 & a_1 & a_2 \\ \hline 1.51e-7 & -1.22e-5 & 4.40e-4 \\ \hline 6.61e-7 & -5.98e-5 & 2.41e-3 \\ \hline 1.46e-6 & -1.72e-4 & 9.08e-3 \\ \hline 1.26e-6 & -1.85e-4 & 1.22e-2 \end{array}$ | $\begin{array}{c ccccccccccccccccccccccccccccccccccc$ |

 Table 2. polynomial parameters for horizon fluxes.

| | b_0 | b_1 | b_2 | b_3 | b_4 | b_5 | b_6 | b_7 | b_8 | b_9 | b_{10} |
|----------|----------|----------|----------|----------|----------|-----------------|---------|---------|--------|---------|----------|
| a = 0.9 | -3.44e-9 | 2.93e-7 | -1.09e-5 | 2.25e-4 | -2.66e-3 | 1.34e-2 | 1.07e-1 | -3.19e0 | 8.84e0 | -8.33e0 | 2.37e0 |
| a = 0.7 | 8.88e-8 | -7.97e-6 | 3.17e-4 | -7.39e-3 | 1.11e-1 | - 1.14e0 | 8.00e0 | -3.89e1 | 1.17e2 | -2.04e2 | 1.58e2 |
| a = 0.0 | 4.84e-7 | -5.70e-5 | 3.00e-3 | -9.28e-2 | 1.88e0 | -2.59e1 | 2.48e2 | -1.62e3 | 6.97e3 | -1.79e4 | 2.11e4 |
| a = -0.9 | 5.30e-7 | -7.76e-5 | 5.11e-3 | -1.99e-1 | 5.08e0 | -8.90e1 | 1.08e3 | -9.10e3 | 5.05e4 | -1.69e5 | 2.63e5 |

Table 3. polynomial coefficients for waveform (2,2) mode: real part.

| | R_{22}^{0} | R_{22}^{1} | R_{22}^2 | R_{22}^{3} | R_{22}^4 | R_{22}^{5} | R_{22}^{6} | R_{22}^{7} | R_{22}^{8} | R_{22}^{9} | R_{22}^{10} |
|----------|--------------|--------------|------------|--------------|------------|--------------|--------------|--------------|--------------|--------------|---------------|
| a = 0.9 | 4.00e-5 | 4.96e-1 | -9.20e-1 | -3.43e-1 | -6.23e0 | 1.73e2 | -1.02e3 | 3.20e3 | -6.14e3 | 6.78e3 | -3.27e3 |
| a = 0.7 | 2.88e-5 | 4.96e-1 | -9.20e-1 | 1.03e0 | -1.98e1 | 2.75e2 | -1.57e3 | 5.33e3 | -1.18e4 | 1.54e4 | -9.09e3 |
| a = 0.0 | 1.78e-5 | 4.97e-1 | -8.67e-1 | 5.44e0 | -6.22e1 | 6.73e2 | -4.23e3 | 1.89e4 | -5.58e4 | 9.83e4 | -7.12e4 |
| a = -0.9 | 1.38e-5 | 4.97e-1 | -7.79e-1 | 1.17e1 | -1.17e2 | 1.34e3 | -9.37e3 | 4.84e4 | -1.46e5 | 2.09e5 | 1.13e5 |

Our model: Fully calibrated version1



Our model: Fully calibrated version1



Figure 5. Orbital evolution of the EMRIs (mass ratio 1/1000) for a = 0; the right panel shows the details of the orbit evolution at the final time.

Han CQG 2016



Our Fully calibrated version1: shortcoming

- Each group of polynomials only works in one special case;
- For getting the polynomials, one firstly must generate some numerical reference data;
- Polynomials do not have physics inside.
- Inspiring us to use factorized forms in EOB frame

$$\rho_{lm} = 1 + a_1 v^2 + b_1 q v^3 + (a_2 + b_2 q^2) v^4 + b_3 q v^5 + [b_4 q^2 + a_3 + a_4 \text{eulerlog}(mv)] v^6$$
$$[b_5 q + b_6 q^3] v^7 + [b_7 q^2 + b_8 q^4 + a_5 + a_6 \text{eulerlog}(mv)] v^8 + (b_9 q + b_{10} q^3) v^9 + [a_7 + a_8 \text{eulerlog}(mv)] v^{10} + [a_9 + a_{10} \text{eulerlog}(mv)] v^{12} \quad \text{for } \mathsf{I} = \mathsf{M}$$

$$\rho_{lm} = 1 + b_1 qv + (a_1 + b_2 q^2)v^2 + (b_3 q + b_4 q^3)v^3 + (a_2 + b_5 q^2 + b_6 q^4)v^4 + (b_7 q + b_8 q^3 + b_9 q^5)v^5 + [a_3 + a_4 \text{eulerlog}(mv) + b_{10} q^2 + b_{11} q^4 + b_{12} q^6]v^6 + [b_{13} q + b_{14} \text{eulerlog}(mv)q + b_{15} q^3 + b_{16} q^5 + b_{17} q^7]v^7 + [a_5 + a_6 \text{eulerlog}(mv)]v^8 + [a_7 + a_8 \text{eulerlog}(mv)]v^{10} + [a_9 + a_{10} \text{eulerlog}(mv)]v^{12} \quad \text{for } \mathbf{I} = /\mathbf{M}$$

$$a_{11}^{\text{Hor},\text{S}} = p_1 q + p_2 q^2 + p_3 q^3 + p_4 q^4 + p_5 q^5 + p_6 q^6 \text{ for horizon absorption}$$

$$a_{12}^{\text{Hor},\text{S}} = p_7 q + p_8 q^2 + p_9 q^3 + p_{10} q^4 + p_{11} q^5 + p_{12} q^6 + p_{13} q^7$$

Version2: recalibrated waveforms in EOB frame We use a least square method to find the global coefficients with the highly accurate Teukolskybased data which q=0.9~0,-0.3,-0.5,-0.7,-0.9

| | l = 2, m = 2 | l = 3, m = 3 | l = 4, m = 4 | l = 5, m = 5 | l = 6, m = 6 | |
|----------|---------------------------------|----------------------------------|--------------------|---------------------|--------------------|-----|
| a_1 | -1.023975805956E+00 | -1.166721079966E + 00 | -1.2227931618E+00 | -1.2487811477E+00 | -1.2619660215E+00 | -1. |
| a_2 | -1.773785209399E+00 | -1.632220902128E+00 | -1.5488861713E+00 | -1.5137494102E+00 | -1.4927232619E+00 | -1. |
| a_3 | 3.318171742361E + 01 | $1.485007940230\mathrm{E}{+}01$ | 1.1208215801E + 01 | 8.7424413748E + 00 | 6.5654511682E + 00 | 4. |
| a_4 | $3.953835298796E{+}01$ | $1.245654119994\mathrm{E}{+01}$ | 1.2709894830E + 01 | 1.1883585053E + 01 | 1.1939038028E + 01 | 1.: |
| a_5 | -1.104050524831E + 03 | $-6.678847665787\mathrm{E}{+02}$ | -8.6428660908E+02 | -9.5407467730E + 02 | -1.0672751384E+03 | -1. |
| a_6 | $1.765670955233\mathrm{E}{+03}$ | $6.364340267479\mathrm{E}{+}02$ | 6.4805577773E + 02 | 6.1182473525E + 02 | 6.1235129084E + 02 | 6.: |
| a_7 | -1.584521768582E + 04 | -7.018461656809E + 03 | -8.2330639578E+03 | -8.5105354171E + 03 | -9.1262632809E+03 | -9. |
| a_8 | 1.007746475744E + 04 | $3.520695972919E{+}03$ | 3.6160554839E + 03 | 3.4023945061E + 03 | 3.4000165921E + 03 | 3.4 |
| a_9 | -1.943757060254E + 04 | -7.516377182961E + 03 | -8.4638217440E+03 | -8.4396976216E + 03 | -8.8310465543E+03 | -9. |
| a_{10} | $6.701134943008E{+}03$ | $2.262034338117\mathrm{E}{+03}$ | 2.3482970193E + 03 | 2.2049397140E + 03 | 2.2024481099E + 03 | 2. |
| b_1 | -6.733755278411E-01 | -6.696972546143E-01 | -6.6915298811E-01 | -6.6901906546E-01 | -6.6904049613E-01 | -6 |
| b_2 | 5.843602324003E-01 | 5.360214944791E-01 | 5.2813489804E-01 | 5.2594458075E-01 | 5.2571830732E-01 | 5. |
| b_3 | -1.662151538548E + 00 | -1.308471947237E+00 | -1.2252358713E+00 | -1.2090008750E+00 | -1.2146809603E+00 | -1. |
| b_4 | -1.061146392648E+00 | -4.975131465872E-01 | -4.5878852649E-01 | -4.6545217148E-01 | -4.8516130249E-01 | -5 |
| b_5 | 3.827859278692E + 00 | $2.543651584540\mathrm{E}{+00}$ | 2.4583522998E + 00 | 2.4535277945E + 00 | 2.4730693730E + 00 | 2.4 |
| b_6 | -7.756293634998E-01 | -3.514346900192E-01 | -1.8135025452E-01 | -1.4988716272E-01 | -1.4384902166E-01 | -1 |

TABLE III. Total coefficients for $m \neq l$

| | l = 5, m = 4 | l = 5, m = 3 | l = 6, m = 5 | l = 6, m = 4 |
|----------|----------------------|--------------------|----------------------|----------------------|
| a_1 | -1.2781593391E+00 | -9.6137951691E-01 | -1.2846696264E+00 | -1.0236865411E+00 |
| a_2 | -9.7034478849E-01 | -6.7312218658E-01 | -1.0280853586E+00 | -8.0181417330E-01 |
| a_3 | 2.8215508168E + 00 | 6.6046438604E + 00 | 1.4723471598E + 00 | $1.3317579430E{+}01$ |
| a_4 | $1.2611970259E{+}01$ | -2.3057664622E+01 | $1.2336196550E{+}01$ | -1.9494642667E + 01 |
| a_5 | -5.2596249376E+02 | 8.0535017960E + 02 | -6.2959344624E+02 | 8.2565169087E + 02 |
| a_6 | 3.0038340199E + 02 | -7.0899298915E+02 | 3.3177604825E + 02 | -5.8059965421E+02 |
| a_7 | -1.7661904159E + 03 | 6.8217825869E + 03 | -2.5543377525E+03 | $6.3637960489E{+}03$ |
| a_8 | 6.0774411335E + 02 | -3.2601890463E+03 | 8.8256123155E + 02 | -2.6727578200E+03 |
| a_9 | -2.4280639567E+02 | 5.8657539419E + 03 | -1.0359578359E+03 | $5.2005709295E{+}03$ |
| a_{10} | 3.2316814336E + 01 | -1.7108913454E+03 | 2.3941084370E + 02 | -1.3988354619E+03 |
| b_1 | -2.4005464853E-01 | 8.6374961193E-06 | -1.9451648765E-01 | -9.2761151062E-06 |
| b_2 | -1.1607843889E-01 | 1.2299226986E-04 | -9.5157480501E-02 | 2.0275436115 E-04 |
| b_3 | 3.2207333467E-01 | -1.1811709880E+00 | 1.6169577622 E-01 | -1.0374290314E+00 |
| b_4 | -8.2537773600E-02 | 5.7012441031E-03 | -6.6960649581E-02 | 5.2750953889E-03 |
| b_5 | 3.3019097935E-01 | 5.8255283839E-01 | 3.7135374972E-01 | 5.2694468496E-01 |
| b_6 | -7.1149022817E-02 | 1.6974336944 E-02 | -5.8801917957E-02 | 1.3846828797E-02 |



Cheng & Han, 2017, submitted to PRD, arXiv: 1706.03884

Including fluxes to horizon



Cheng & Han, 2017, submitted to PRD, arXiv: 1706.03884

Dephasing between recalibrated model and Teukolskybased waveforms



Evolution



q=0.9, M=10^5 solar mass, r0=11M

Waveforms



Conclusions

- The recalibrated formalism looks ugly but works
- Only works for EMRIs
- Published all the coefficients and will publish our codes;
- It is an effort to construct the waveform templates for LISA et. al. in the EOB formalisms.
- recalibrating mass-ratios dependent terms;

Thank you!