



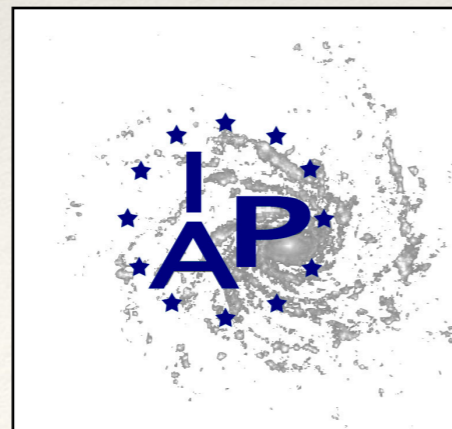
TÉCNICO
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The Era of Gravitational Wave Astronomy, IAP

Paris, June 2017

Gravitational waves from compact binaries in scalar-tensor gravity to 2PN order

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Outline

- ❖ Scalar-Tensor Gravity
- ❖ Motivation
- ❖ To-date
- ❖ Landau-Lifshitz Formalism
- ❖ Direct Integration of Relaxed Einstein Equations (DIRE)
- ❖ To-do

Scalar-Tensor Gravity

- ❖ Alternate theory of gravity (ATG)

- ❖ Alternative to Newtonian:

Nordström: $g_{\mu\nu} = \Phi\eta_{\mu\nu}$, General Relativity: $8\pi GT_{\mu\nu} = R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu}$.

- ❖ Alternative to GR:

- ❖ Supported by same experimental evidence: Whitehead, Scalar-Tensor.

- ❖ Seeks to explain recent observations / issues: (super)String Theory, f(R), etc.

- ❖ Variable gravitational “constant” → Scalar field(s),

$$S_{GR} = \frac{1}{16\pi G} \int R\sqrt{-g}d^4x + \int \mathcal{L}(m, g_{\alpha\beta}) \sqrt{g}d^4x,$$

$$T^{\alpha\beta}_{,\beta} = 0$$

$$S_{ST} = \frac{1}{16\pi} \int \left[\phi R - \frac{1}{\phi} \omega(\phi) g^{\alpha\beta} \partial_\alpha \phi \partial_\beta \phi \right] \sqrt{-g}d^4x + \int \mathcal{L}(m, g_{\alpha\beta}) \sqrt{g}d^4x.$$

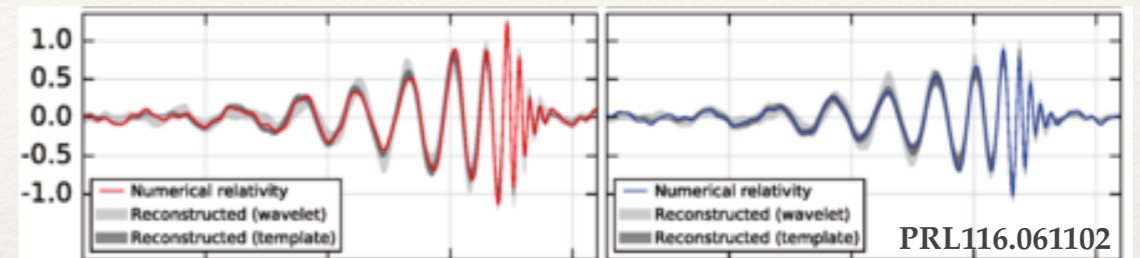
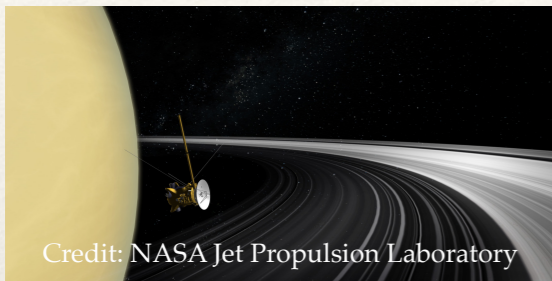
- ❖ Obeys Einstein’s Equivalence principle

- ❖ Violates the Strong Equivalence Principle

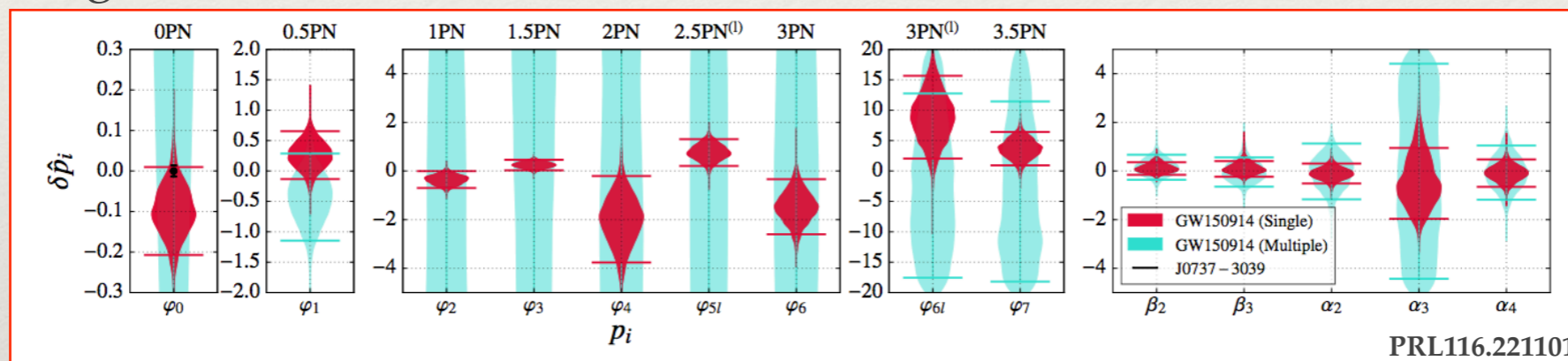
Internal Structure
effects motion and
GW emission

Motivation

- ❖ Seek to verify / constrain / discard ATG's
- ❖ New tool → Gravitational Wave Astronomy



- ❖ Current Testing



- ❖ Why Scalar Tensor?

- ❖ One of the simplest variations of GR
- ❖ Encapsulates some of F(R) and Superstring theories

Warning

- ❖ Hawking Theorem - GR and ST black hole binaries indistinguishable
- ❖ General argument for small differences expected at every PN order :(

To date ...

❖ Necessary Ingredients:

3PN
EOM

❖ EOM (Mirshekeri & Will 2013: 2.5PN)

❖ Tensor gravitational waves and tensorial energy flux 2PN
(Lang 2014: 2PN)

2.5PN
 Ψ

❖ Scalar gravitational waves and scalar energy flux (Lang 2015:
1.5PN and 1PN respectively)

❖ Ready to use waveforms (Sennett, Marsat, Buonanno 2016:
incomplete 2PN)

❖ Culprit: Non-vanishing scalar dipole moment:

$\Psi = \Psi_{-1/2} + \Psi_0 + \Psi_{1/2} + \Psi_1 + \Psi_{3/2}, \Rightarrow$ We require $\Psi_{n+1/2}$ for $\dot{E}_n,$

$\dot{E}_S \propto \dot{\Psi}^2 \Rightarrow \dot{E} = \dot{E}_{-1} + \dot{E}_0 + \dot{E}_{1/2} + \dot{E}_1. \Rightarrow$ We require EOM_(n+1) for $\Psi_{n+1/2}.$

Landau-Lifshitz Formalism

❖ Field equations:

$$G_{\alpha\beta} = \frac{8\pi}{\phi} T_{\mu\nu} + \frac{\omega(\phi)}{\phi^2} \left(\phi_{,\alpha} \phi_{,\beta} - \frac{1}{2} g_{\alpha\beta} \phi_{,\gamma} \phi^{,\gamma} \right) + \frac{1}{\phi} (\phi_{,\alpha\beta} - g_{\alpha\beta} \square_g \phi),$$

$$\square_g \phi = \frac{1}{3 + 2\omega(\phi)} \left(8\pi T - 16\pi \phi \frac{\partial T}{\partial \phi} - \frac{d\omega}{d\phi} \phi_{,\gamma} \phi^{,\gamma} \right).$$

$$\psi = \frac{\phi}{\phi_0}.$$

❖ Wave equations:

$$\square_\eta \tilde{h}^{\alpha\beta} = -16\pi \tau^{\alpha\beta}, \quad \square_\eta \psi = -8\pi \tau_s.$$

$$\tau^{\alpha\beta}_{,\beta} = 0$$

$$16\pi \tau^{\alpha\beta} = 16\pi (-g) \frac{\psi}{\phi_0} T^{\alpha\beta} + \Lambda^{\alpha\beta} + \Lambda_S^{\alpha\beta},$$

$$\tau_s = -\frac{1}{3 + 2\omega} \sqrt{-g} \frac{\psi}{\phi_0} \left(T - 2\psi \frac{\partial T}{\partial \psi} \right) - \frac{1}{8\pi} \tilde{h}^{\mu\nu} \psi_{,\mu\nu} + \frac{1}{16\pi} \frac{d}{d\psi} \left[\ln \left(\frac{3 + 2\omega}{\psi^2} \right) \right] \psi_{,\mu} \psi_{,\nu} \tilde{g}^{\mu\nu}$$

❖ Integration of wave equations via flat Green function:

$$\tilde{h}^{\alpha\beta}(t, x) = 4 \int \frac{\tau^{\alpha\beta}(t', x') \delta(t' - t + |x - x'|)}{|x - x'|} d^4 x', \quad \psi(t, x) = 2 \int \frac{\tau_s(t', x') \delta(t' - t + |x - x'|)}{|x - x'|} d^4 x'.$$

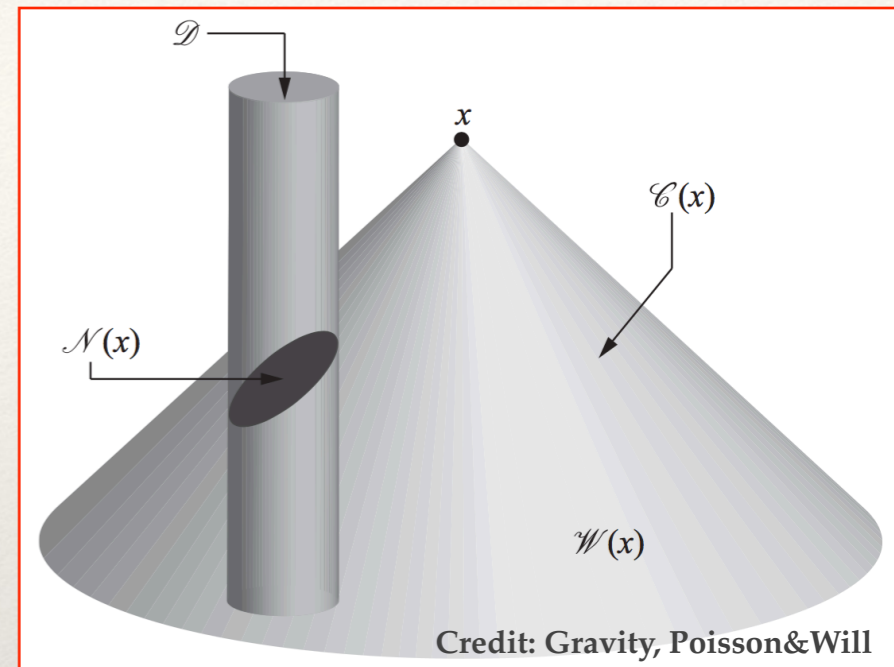
Direct Integration of Relaxed Einstein Equations

- ❖ Zones: Near zone, wave zone

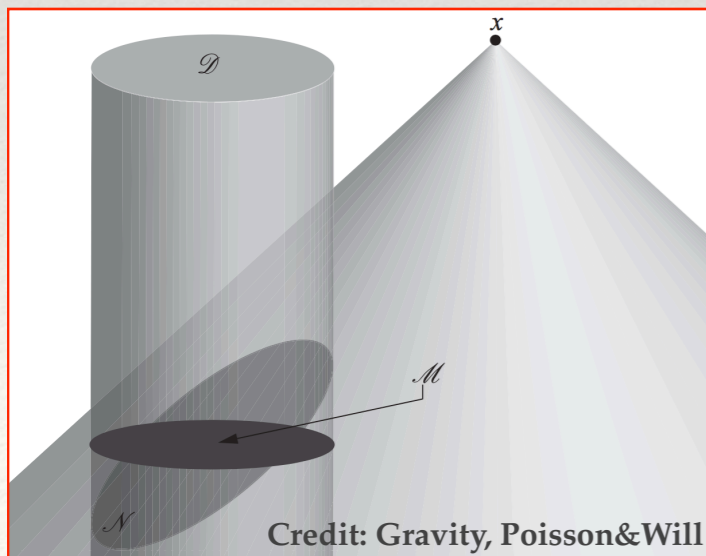
$$\Psi = \Psi_{\mathcal{N}} + \Psi_{\mathcal{W}}, \quad h^{\alpha\beta} = h_{\mathcal{N}}^{\alpha\beta} + h_{\mathcal{W}}^{\alpha\beta}$$

$$\tilde{h}_{(\mathcal{N}/\mathcal{W})}^{\alpha\beta}(t, x) = 4 \int_{\mathcal{N}/\mathcal{W}} \frac{\tau^{\alpha\beta}(t', x') \delta(t' - t + |x - x'|)}{|x - x'|} d^4 x',$$

$$\psi_{(\mathcal{N}/\mathcal{W})}(t, x) = 2 \int_{\mathcal{N}/\mathcal{W}} \frac{\tau_s(t', x') \delta(t' - t + |x - x'|)}{|x - x'|} d^4 x'.$$



- ❖ Near zone field, wave zone field point

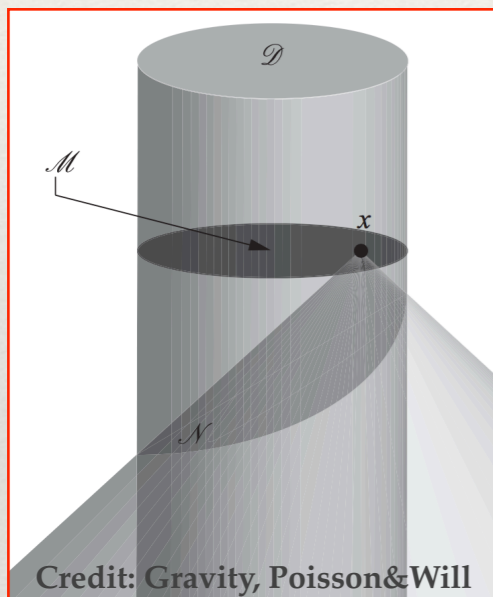


$$\begin{aligned} \psi_{\mathcal{N}}/h_{\mathcal{N}}^{\alpha\beta} &= \int_{\mathcal{N}} \int \frac{\tau^A(t - |x - x'|, y) \delta(y - x') d^3 y d^3 x'}{|x - x'|} \Rightarrow x' \text{ small} \\ &= \sum_{l=0}^{\infty} \frac{(-1)^l}{l!} \partial_L \left[\frac{1}{r} \int_{\mathcal{M}} \tau^A(\tau, x') x'^L d^3 x' \right], \\ &= \frac{1}{r} \sum_{l=0}^{\infty} \frac{1}{l!} n_L \left(\frac{d}{d\tau} \right)^l \int_{\mathcal{M}} \tau^A(\tau, x') x'^L d^3 x' + \mathcal{O}(r^{-2}). \end{aligned}$$

Far away
wave zone

Direct Integration of Relaxed Einstein Equations

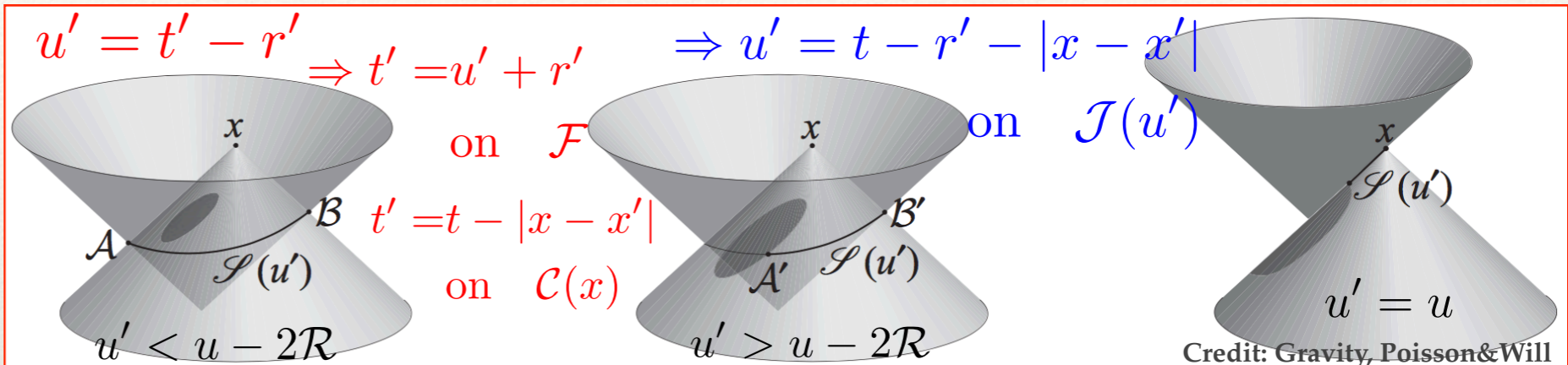
- ❖ Near zone field, near zone field point



$\Rightarrow |x - x'| \text{ small}$

$$\psi_{\mathcal{N}}/h_{\mathcal{N}}^{\alpha\beta} = \sum_{l=0}^{\infty} \frac{(-1)^l}{l!} \left(\frac{\partial}{\partial t} \right)^l \times \int_{\mathcal{M}} \tau^A(t, x') |x - x'|^{l-1} d^3 x'$$

- ❖ Far zone field

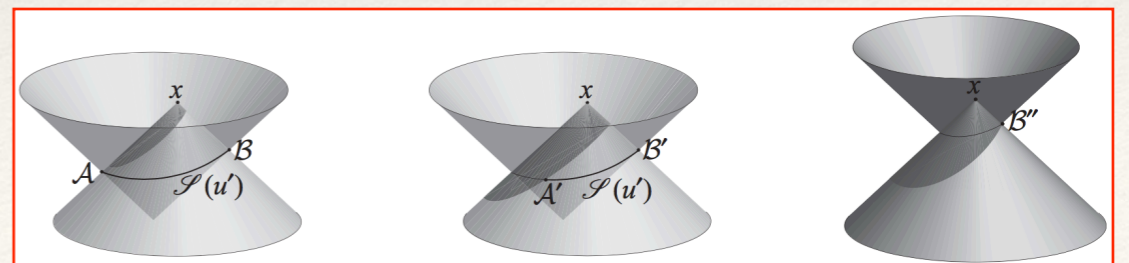


$$\begin{aligned} \psi_{\mathcal{W}}/h_{\mathcal{W}}^{\alpha\beta} &= \int_{-\infty}^u du' \oint_{\mathcal{J}(u')} \frac{\tau^A(u' + r', x')}{t - u' - n' \cdot x} r' (u', \theta', \phi')^2 d\Omega', \quad \tau^A(x') = \frac{1}{4\pi} \frac{f^A(\tau')}{r'^n} n'^{\langle L \rangle} \\ &= \frac{1}{2} n^{\langle L \rangle} \int_{-\infty}^u du' f(u') \int_{\mathcal{R}} \frac{P_l(\xi)}{r r'^{(n-1)}} dr' \end{aligned} \quad n = e_z$$

$$A(s, r) = \int_{\mathcal{R}} \frac{P_l(\cos \theta')}{r'^{(n-1)}} dr', \quad B(s, r) = \int_s^{r+s} \frac{P_l(\cos \theta')}{r'^{(n-1)}} dr', \quad s = \frac{1}{2}(u - u')$$

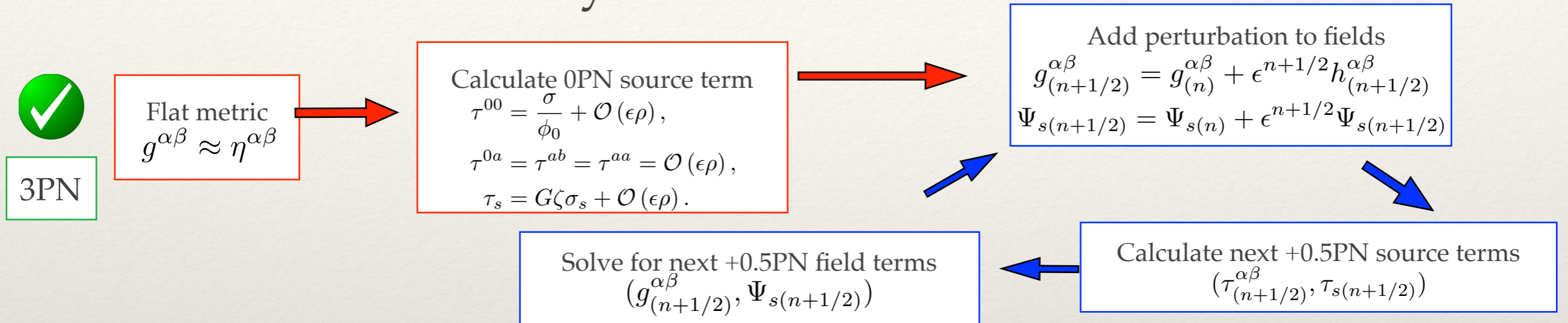
$$\psi_{\mathcal{W}}/h_{\mathcal{W}}^{\alpha\beta} = \frac{n^{\langle L \rangle}}{r} \left[\int_{\mathcal{R}-r}^{\mathcal{R}} ds f^A(\tau - 2s) A(s, r) + \int_{\mathcal{R}}^{\infty} ds f^A(\tau - 2s) B(s, r) \right] \text{Far zone field pt}$$

Near zone field pt



Direct Integration of Relaxed Einstein Equations

- ❖ Solve simultaneously for source terms $(\tau^{\alpha\beta}, \tau_s)$ and fields $(g^{\alpha\beta}, \Psi)$



- ❖ Stress energy tensor

$$T^{\alpha\beta} = \rho * (-g)^{-1/2} u^\alpha u^\beta (u^0)^{-1}, \quad \rho * = \sum_A m_A \delta^3(x - x_A),$$

$$\Rightarrow T^{\alpha\beta} = \rho * (-g)^{-1/2} v^\alpha v^\beta u^0 [1 + S(s; \Psi)]$$

$$\sigma = \rho * (-g)^{-1/2} u^0 (1 + v^2) [1 + S(s; \Psi)], \quad \sigma^a = \rho * (-g)^{-1/2} u^0 v^a [1 + S(s; \Psi)],$$

$$\sigma^{ab} = \rho * (-g)^{-1/2} u^0 v^a v^b [1 + S(s; \Psi)], \quad \sigma_s = \rho * (-g)^{-1/2} (u^0)^{-1} [1 - 2s + S_s(s; \Psi)].$$

$$S(s; \Psi) = \epsilon \Psi s + \frac{1}{4} \epsilon^2 \Psi^2 (2a_s - s) + \frac{1}{24} \epsilon^3 \Psi^3 (-6a_s + 3s + 4b_s)$$

$$+ \frac{1}{192} \epsilon^4 \Psi^4 (30a_s - 15s - 20b_s + 24c_s) + \mathcal{O}(\epsilon^5),$$

$$S_s(s; \Psi) = -2\epsilon (\Psi a_s) - \epsilon^2 \Psi^2 b_s - \epsilon^3 \Psi^3 c_s + \mathcal{O}(\epsilon^4)$$

$$a_s \equiv s^2 + s' - \frac{1}{2}s,$$

$$b_s \equiv a'_s - a_s - s a_s,$$

$$c_s \equiv \frac{1}{3} (b'_s - 2b_s - s b_s).$$

$$\sigma \equiv T^{00} + T^{ii},$$

$$\sigma^a \equiv T^{0a},$$

$$\sigma^{ab} \equiv T^{ab},$$

$$\sigma_s \equiv -T + 2\psi \frac{\partial T}{\partial \psi}.$$

$$s_A \equiv \left(\frac{d \ln M_A(\phi)}{d \ln \phi} \right)_0$$

To-do

