

4.5 Post-Newtonian order gravitational radiation

Tanguy Marchand - IAP
marchand@iap.fr

*T. Marchand, L. Blanchet, G. Faye 2016 (CQG)
arxiv:1607.07601*

IAP - 26th June 2017

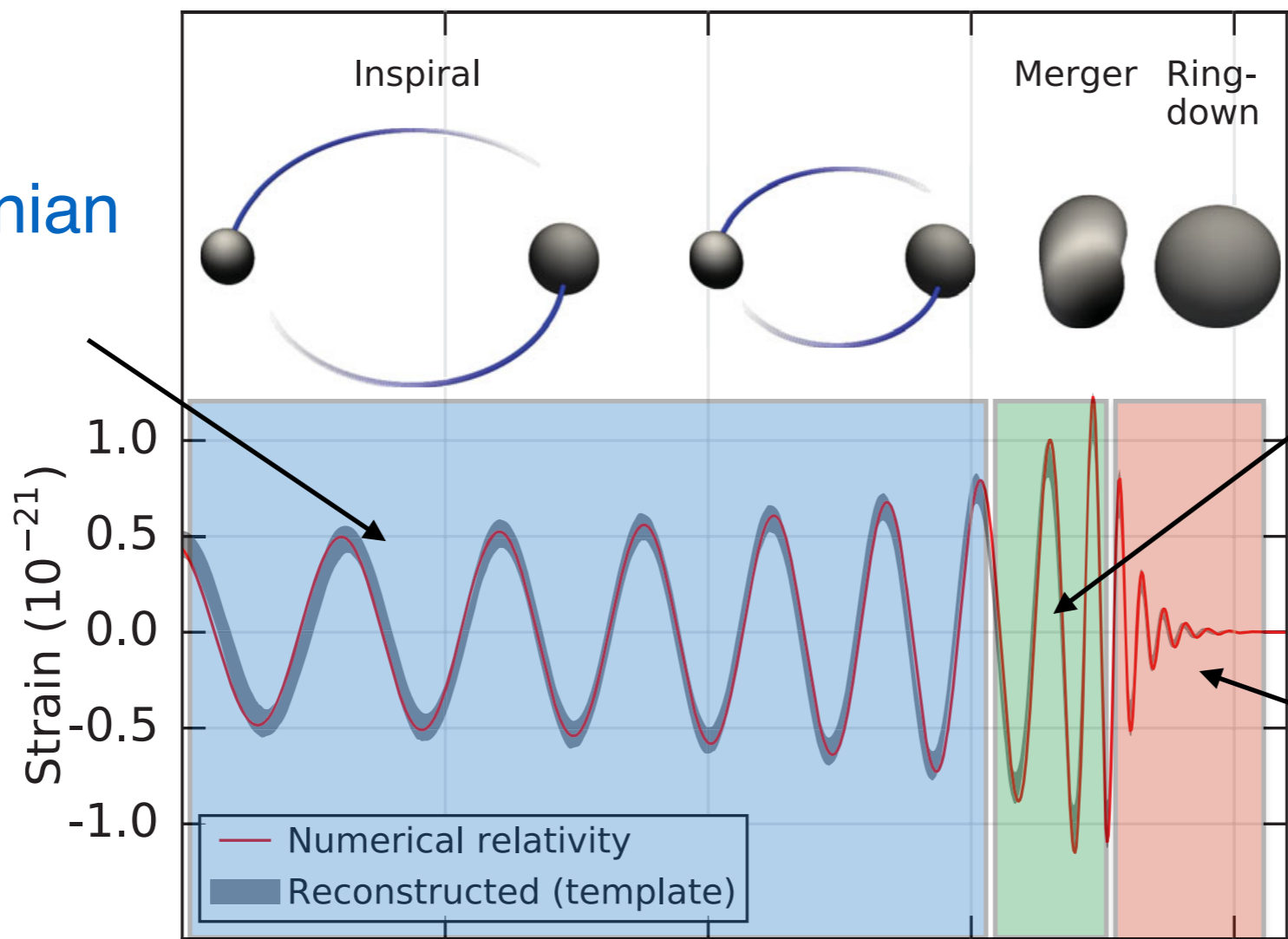
I. Introduction

II. The multipolar-post-Minkowskian algorithm

III. The 4.5PN project

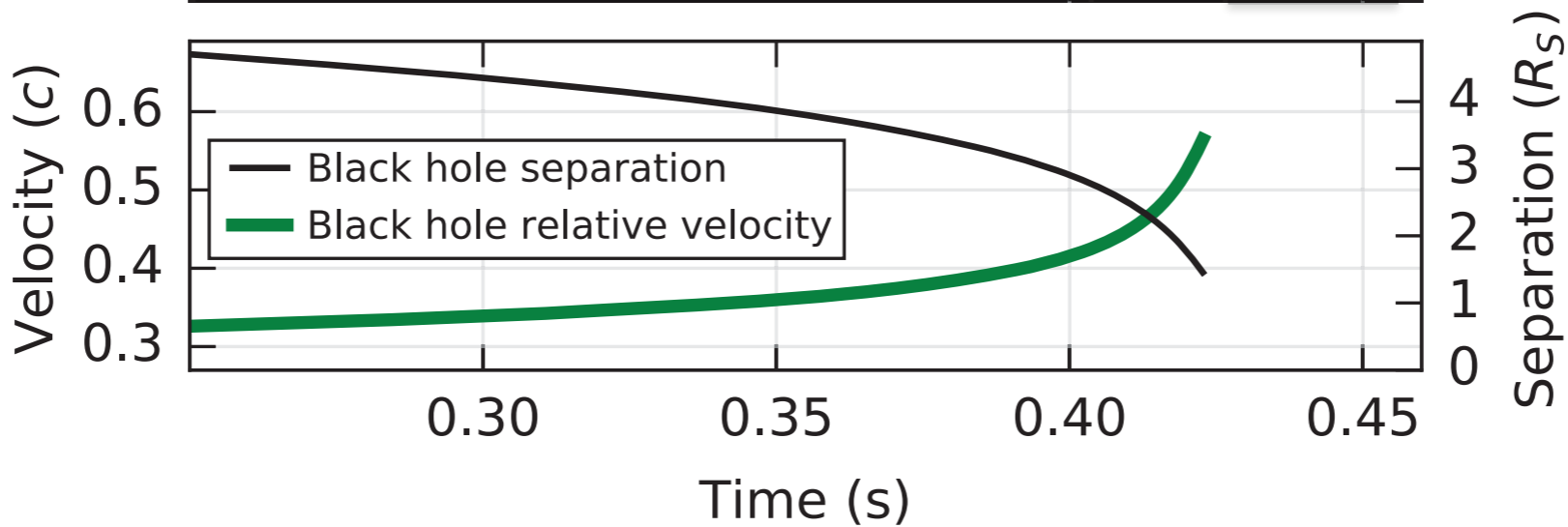
I. INTRODUCTION

Post-Newtonian theory



Numerical relativity

BH perturbations QNM



PRL 116, 061102 (2016)

Post-Newtonian theory

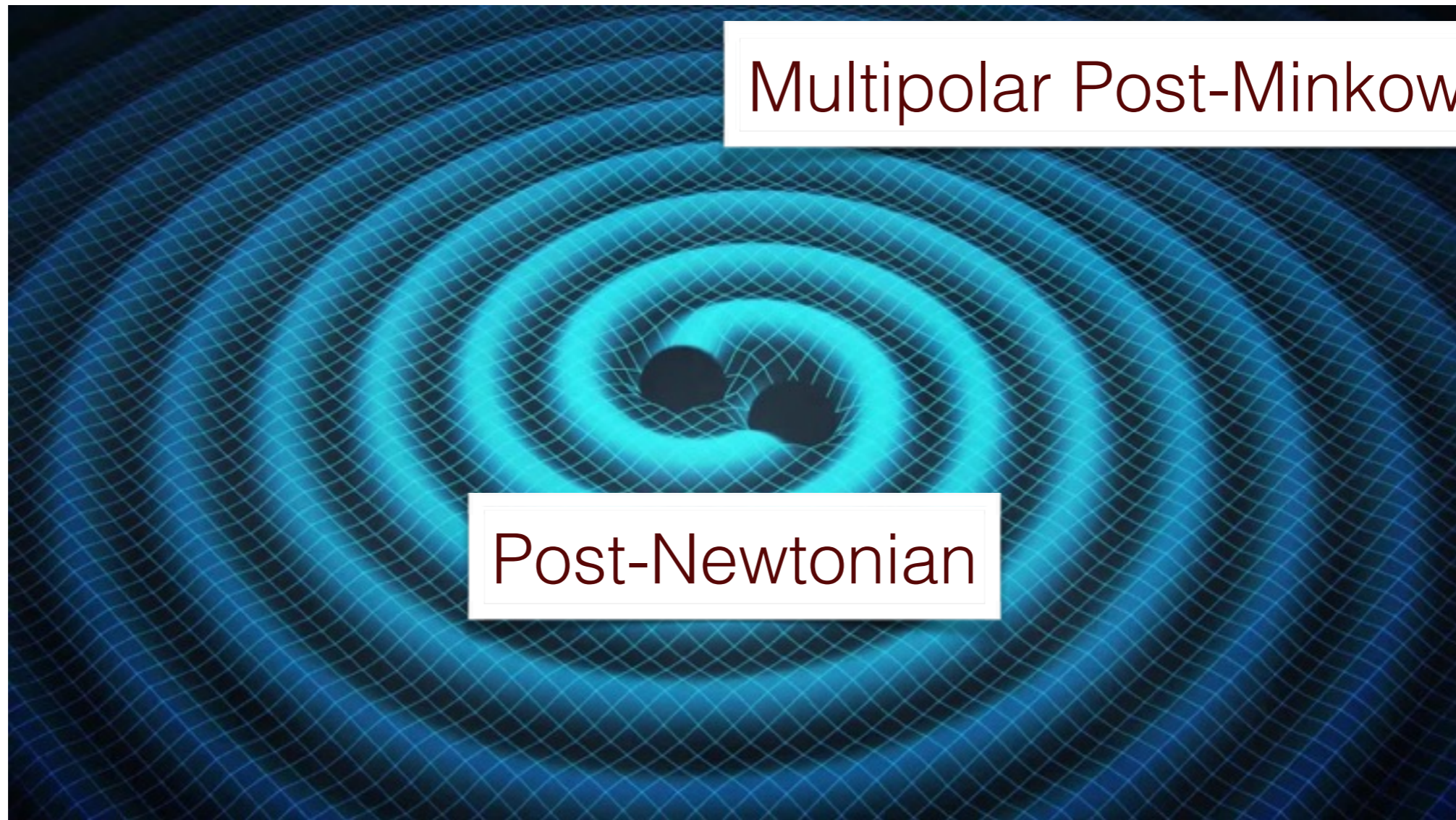
- ▶ Perturbative expansion of relativistic effects

- ▶ 1 PN $\rightarrow \left(\frac{v}{c}\right)^2$

- ▶ More and more difficulties appear as we go to higher orders

Blanchet-Damour-Iyer formalism

Blanchet-Damour-Iyer formalism



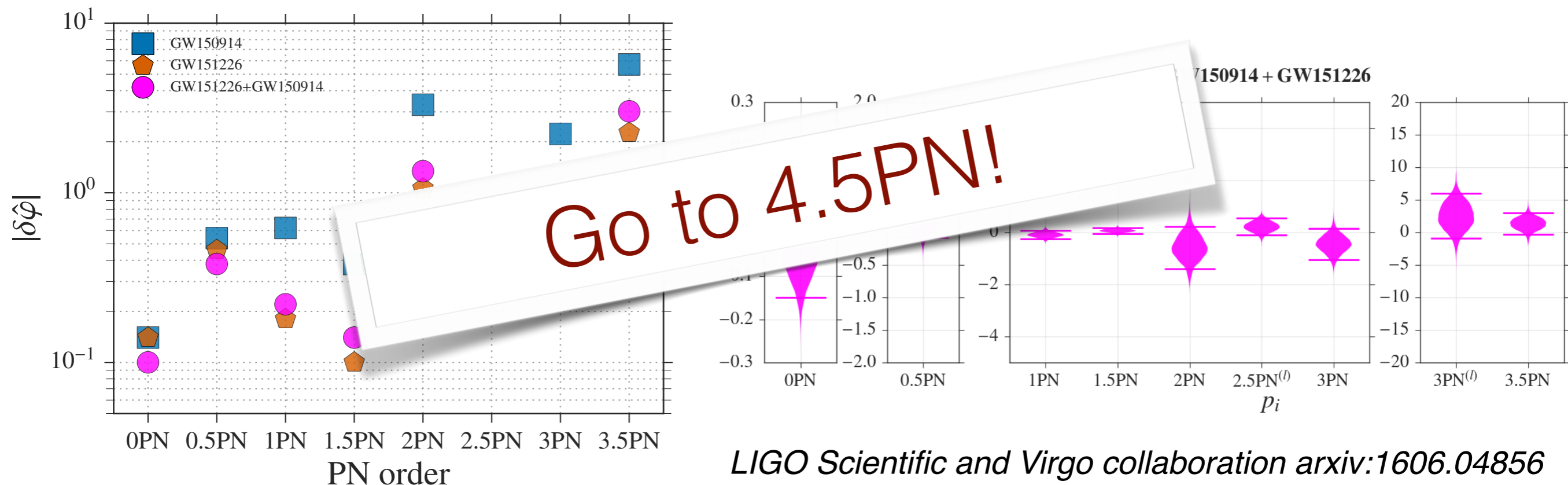
*[Matching equation Blanchet 1998]
figure: www.virgo-gw.eu*

Once the **equations of motions** and the **flux** is known at n-PN

We consider consecutive circular orbits of orbital frequencies Ω , energy E_n and emitting a flux \mathcal{F}_n

$$\frac{dE_n(x)}{dt} = \mathcal{F}_n(x) \quad \Rightarrow \quad \phi_n = f_n(x)$$

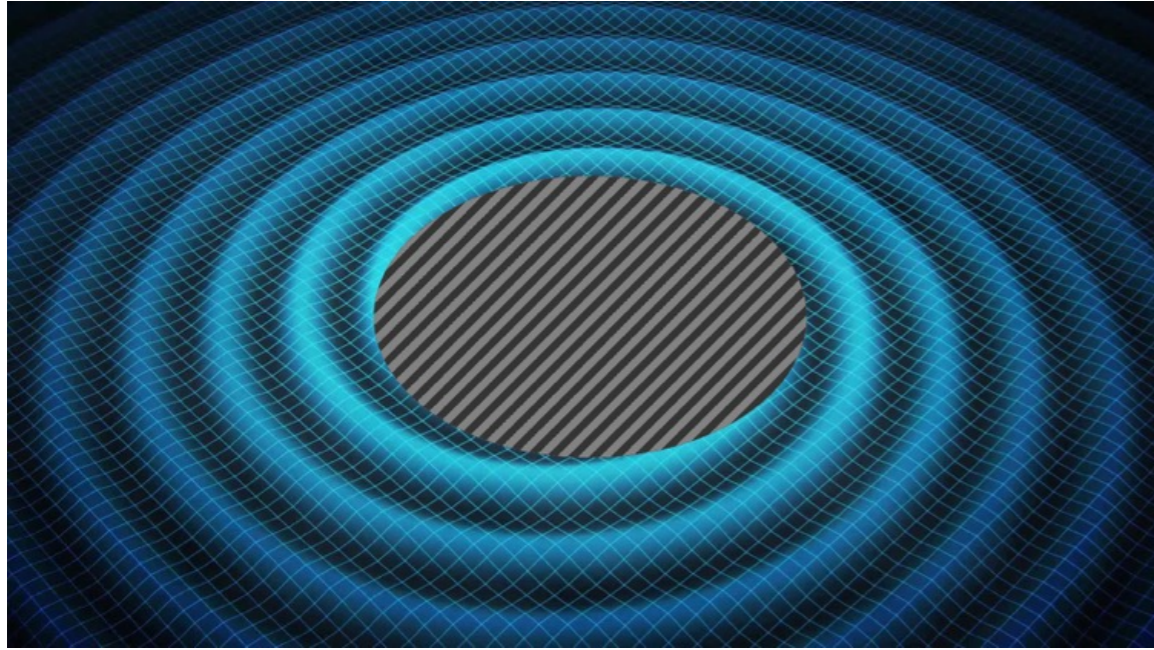
$$x = \left(\frac{Gm\Omega}{c^3} \right)^{2/3}$$



LIGO Scientific and Virgo collaboration arxiv:1606.04856

II. The multipolar post-Minkowskian (MPM) algorithm

The MPM algorithm

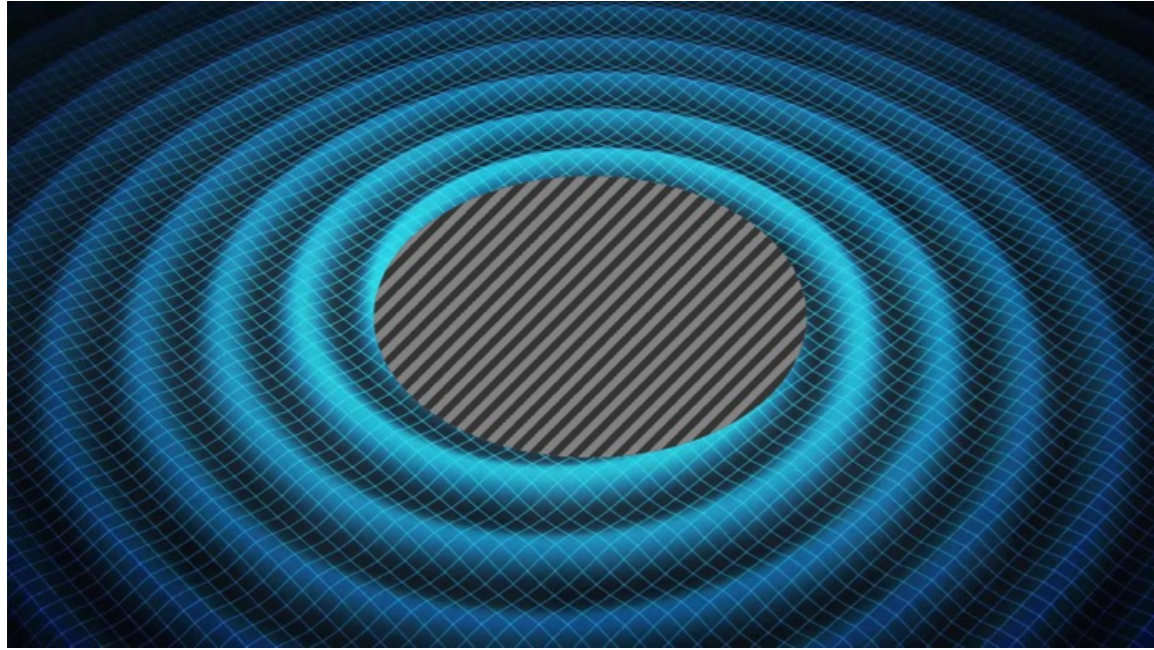


$$G_{\mu\nu}(g_{\alpha\beta}, \partial g_{\alpha\beta}, \partial^2 g_{\alpha\beta}) = 0$$

$$h^{\mu\nu} \equiv \sqrt{-g}g^{\mu\nu} - \eta^{\mu\nu} = \mathcal{G}h_{(1)}^{\mu\nu} + \mathcal{G}^2h_{(2)}^{\mu\nu} + \dots$$

$$\begin{cases} \square h_{(i)}^{\mu\nu} = \Lambda_{(i)}^{\mu\nu}(h_{(1)}^{\alpha\beta}, \dots, h_{(i-1)}^{\alpha\beta}) \\ \partial^\mu h_{(i)\mu\nu} = 0 \quad [harmonic\ coordinates] \end{cases}$$

The MPM algorithm

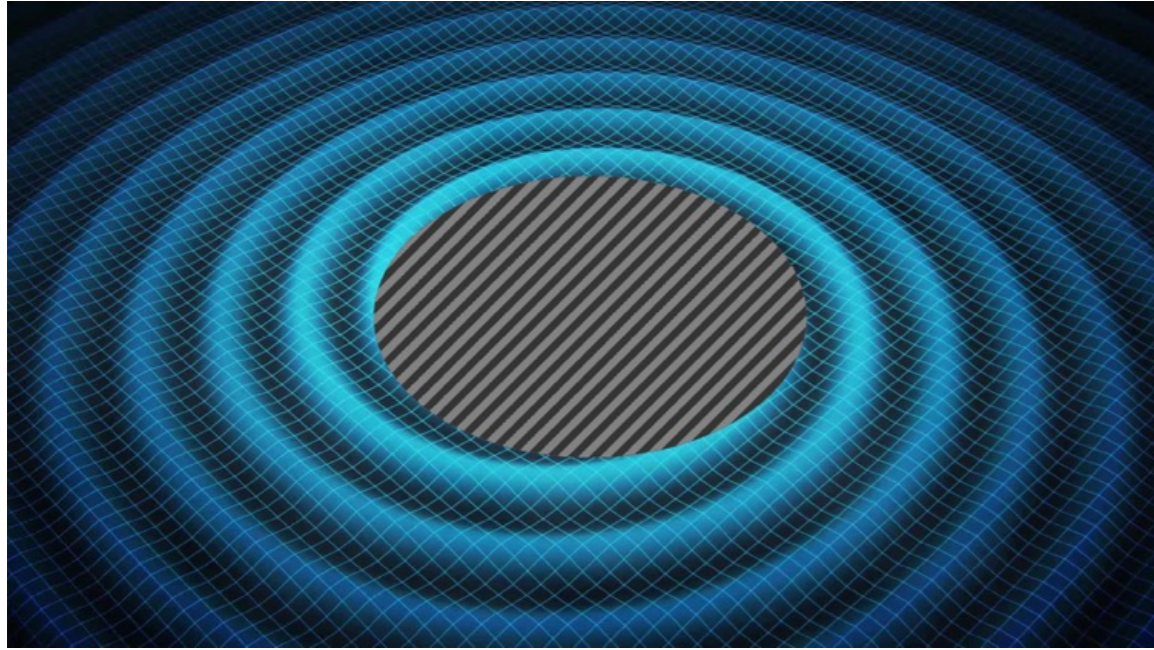


$$\begin{cases} \square h_{(i)}^{\mu\nu} = \Lambda_{(i)}^{\mu\nu}(h_{(1)}^{\alpha\beta}, \dots, h_{(i-1)}^{\alpha\beta}) \\ \partial^\mu h_{(i)\mu\nu} = 0 \end{cases}$$

$$\begin{aligned} h_{(1)}^{\mu\nu} &\sim \sum_{l \geq 0} \partial_{i_1, \dots, i_l} \left(\frac{M_{i_1 \dots i_l}(t-r)}{r} \right) + \sum_{l \geq 2} \partial_{i_1, \dots, i_l} \left(\frac{S_{i_1 \dots i_l}(t-r)}{r} \right) \\ &= h_M^1 + h_{M_{ij}}^1 + h_{M_{ijk}}^1 + \dots + h_{S_{ij}}^1 + h_{S_{ijk}}^1 + \dots \quad [\text{Thorne 80}] \end{aligned}$$

$$h_{(2)}^{\mu\nu} = h_{(2)M \times M}^{\mu\nu} + h_{(2)M \times M_{ij}}^{\mu\nu} + h_{(2)M_{ij} \times M_{ij}}^{\mu\nu} + \dots$$

First issue: regularization



$$\begin{cases} \square h_{(i)}^{\mu\nu} = \Lambda_{(i)}^{\mu\nu}(h_{(1)}^{\alpha\beta}, \dots, h_{(i-1)}^{\alpha\beta}) \\ \partial^\mu h_{(i)\mu\nu} = 0 \end{cases}$$

$$\square^{-1}\Lambda(x, t) = \int d^3x' \frac{\Lambda(x', t - |x - x'|)}{|x - x'|}$$

Issue: $\Lambda \sim_{r \rightarrow 0} \frac{1}{r^k}, k \geq 3$

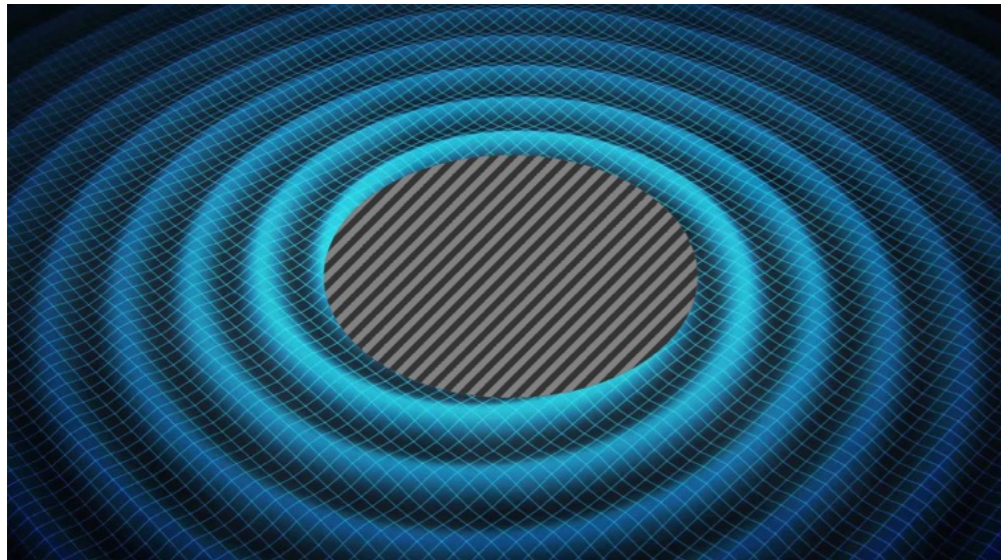


$$\text{FP}_{B=0} \square^{-1} \left[\left(\frac{r}{r_0} \right)^B \Lambda \right]$$

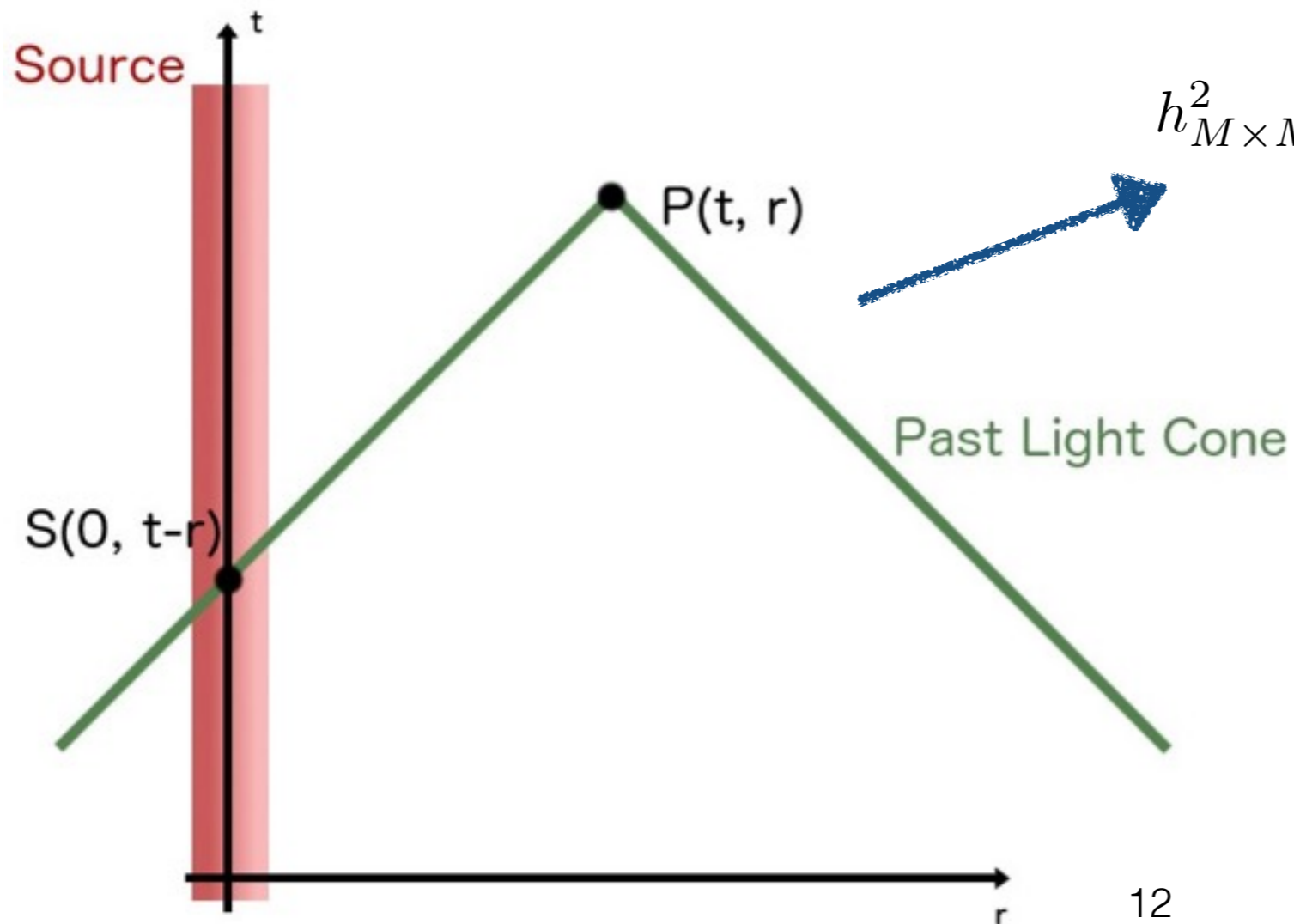
[Analytic continuation in $B \in \mathbb{C}$]

[Finite Part when $B \rightarrow 0$]

Second issue: tails



$$\begin{cases} \square h_{(i)}^{\mu\nu} = \Lambda_{(i)}^{\mu\nu}(h_{(1)}^{\alpha\beta}, \dots, h_{(i-1)}^{\alpha\beta}) \\ \partial^\mu h_{(i)\mu\nu} = 0 \end{cases}$$



$$h_{M \times M_{ij}}^2(t, r) \sim M \int_0^\infty d\tau M_{ij}(t - r - \tau) \mathcal{Q}(\tau)$$

4.5PN project

Ultimate Goal: compute the flux up to 4.5PN

Done so far: compute all the 4.5PN contributions of the tails:

$$h_{(2)M \times M_{ij}}^{\mu\nu}, h_{(3)M \times M \times M_{ij}}^{\mu\nu}, h_{(4)M \times M \times M \times M_{ij}}^{\mu\nu}$$

$$h_{(2)M \times M_{ijk}}^{\mu\nu}, h_{(3)M \times M \times M_{ijk}}^{\mu\nu}$$

$$h_{(2)M \times M_{ijkl}}^{\mu\nu}, \dots$$

[TM Blanchet Faye (2016) 1607.07601]

► Required new analytical formulae

$$\begin{aligned} & \text{FP}_{B=0} \square^{-1} \left[\hat{n}_L \left(\frac{r}{r_0} \right)^B r^{-k} \int_1^\infty dy Q_m(y) F(t - ry) \right] \\ &= -\hat{n}_L \int_1^\infty ds F^{(k-2)}(t - rs) \left(Q_l(s) \int_1^s dy Q_m^{(-k+2)}(y) P_l(y) + P_l(s) \int_s^\infty dy Q_m^{(-k+2)}(y) Q_l(y) \right) \end{aligned}$$

► Implementing the algorithm into Mathematica

[xAct - xTensor]

Going to radiative coordinate

$$h_{M^3 \times M_{ij}}^{00} = \frac{M^3 \hat{n}_{ab}}{r} \int_0^{+\infty} d\tau M_{ab}^{(6)} \left\{ -\frac{8}{3} \ln^3 \left(\frac{\tau}{2r} \right) + \frac{148}{21} \ln^2 \left(\frac{\tau}{2r} \right) + \frac{232}{21} \ln \left(\frac{r}{r_0} \right) \ln \left(\frac{\tau}{2r} \right) \right. \\ \left. + \frac{1016}{2205} \ln \left(\frac{\tau}{2r} \right) + \frac{104}{15} \ln \left(\frac{r}{r_0} \right) + \frac{16489}{1575} - \frac{232\pi^2}{63} \right\} + \mathcal{O} \left(\frac{1}{r^{2-\epsilon}} \right)$$



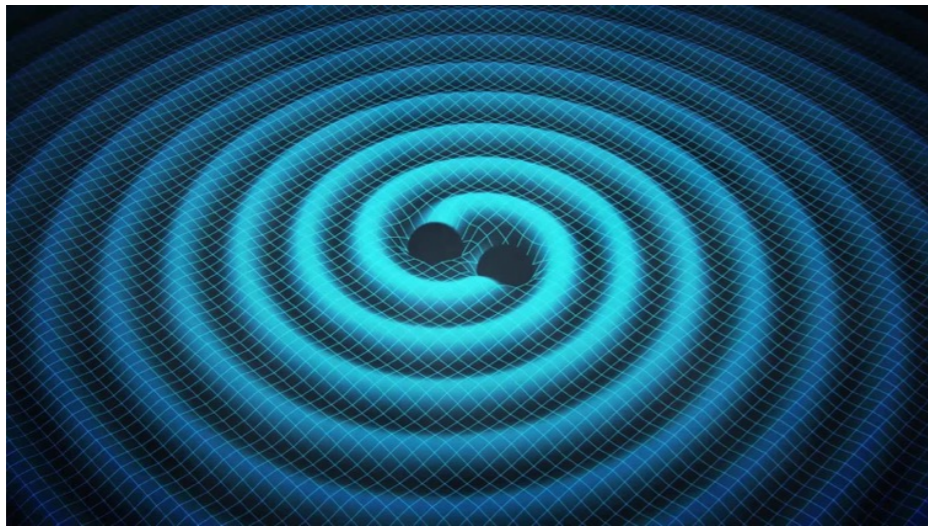
$$X^\mu = x^\mu + \xi^\mu(x) \quad \begin{cases} \xi^0 = -2M \ln \left(\frac{r}{b_0} \right) \\ \xi^i = 0 \end{cases}$$

$$H_{M^3 \times M_{ij}}^{00} = \frac{M^3 \hat{N}_{ab}}{R} \int_0^{+\infty} d\tau M_{ab}^{(6)} \left\{ -\frac{8}{3} \ln^3 \left(\frac{\tau}{2b_0} \right) - 4 \ln^2 \left(\frac{\tau}{2b_0} \right) + \frac{232}{21} \ln \left(\frac{\tau}{2b_0} \right) \ln \left(\frac{\tau}{2r_0} \right) \right. \\ \left. - \frac{14272}{2205} \ln \left(\frac{\tau}{2b_0} \right) + \frac{104}{15} \ln \left(\frac{\tau}{2r_0} \right) + \frac{16489}{1575} - \frac{232\pi^2}{63} \right\} + \mathcal{O} \left(\frac{1}{R^2} \right)$$

We can compute the flux directly from $H^{\mu\nu}$

Matching equation to compute the multipole moments

[e.g. Blanchet 1998]



Tails

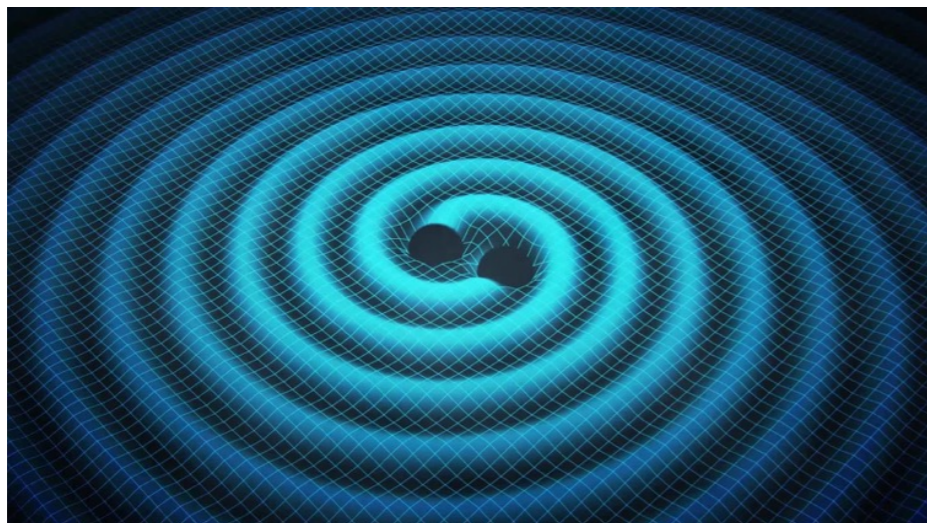
$$\nu = \frac{m_1 m_2}{(m_1 + m_2)^2}$$

$$\gamma = \frac{Gm}{rc^2} \quad (\text{circular orbits})$$

$$\mathcal{F}_{\text{tails}} = \frac{32c^5}{5G} \nu^2 \gamma^5 \left\{ 4\pi\gamma^{3/2} + \dots \right. \\ \left. + \left(\frac{9997778801}{106444800} - \frac{6848}{105} \ln\left(\frac{r}{r_0}\right) + \left[-\frac{8058312817}{2661120} + \frac{287}{32}\pi^2 + \frac{572}{3} \ln\left(\frac{r}{r'_0}\right) \right] \nu \right. \right. \\ \left. \left. - \frac{12433367}{13824} \nu^2 - \frac{1026257}{266112} \nu^3 \right) \pi\gamma^{9/2} + \mathcal{O}\left(\frac{1}{c^{11}}\right) \right\},$$

Tails-of-tails-of-tails and Tails-of-tails \times tails

$$\mathcal{F}_{\text{T-T-T and T-T*T}} = \frac{32c^5}{5G} \nu^2 \gamma^5 \left\{ \left(-\frac{467044}{3675} - \frac{3424}{105} \ln(16\gamma) + \frac{6848}{105} \ln\left(\frac{r}{r_0}\right) - \frac{6848}{105} \gamma_E \right) \pi\gamma^{9/2} \right. \\ \left. + \mathcal{O}\left(\frac{1}{c^{11}}\right) \right\}.$$



$$\nu = \frac{m_1 m_2}{(m_1 + m_2)^2}$$

$$x = \left(\frac{Gm\Omega}{c^3} \right)^{2/3} = \mathcal{O} \left(\frac{1}{c^2} \right)$$

$$\begin{aligned} \mathcal{F}_{\text{total}} = \frac{32c^5}{5G} \nu^2 x^5 \left\{ 1 + \left(-\frac{1247}{336} - \frac{35}{12} \nu \right) x + 4\pi x^{3/2} + \left(-\frac{44711}{9072} + \frac{9271}{504} \nu + \frac{65}{18} \nu^2 \right) x^2 \right. \\ + \left(-\frac{8191}{672} - \frac{583}{24} \nu \right) \pi x^{5/2} + \left[\frac{6643739519}{69854400} + \frac{16}{3} \pi^2 - \frac{1712}{105} \gamma_E \right. \\ \left. - \frac{856}{105} \ln(16x) + \left(-\frac{134543}{7776} + \frac{41}{48} \pi^2 \right) \nu - \frac{94403}{3024} \nu^2 - \frac{775}{324} \nu^3 \right] x^3 \\ + \left(-\frac{16285}{504} + \frac{214745}{1728} \nu + \frac{193385}{3024} \nu^2 \right) \pi x^{7/2} + (\text{unknown coefficients}) x^4 \\ + \left(\frac{265978667519}{745113600} - \frac{6848}{105} \gamma_E - \frac{3424}{105} \ln(16x) + \left[\frac{2062241}{22176} + \frac{41}{12} \pi^2 \right] \nu \right. \\ \left. - \frac{133112905}{290304} \nu^2 - \frac{3719141}{38016} \nu^3 \right) \pi x^{9/2} + \mathcal{O}(x^5) \left. \right\}. \end{aligned}$$

[Tanaka et al. gr-qc/9701050]



[See also Messina&Nagar 2017]

arxiv:1607.07601

Conclusion

- ▶ Goal: reach 4.5PN Flux in order to have the 4.5PN phase
- ▶ Done so far: Tail effects up to 4.5PN, and the 4.5PN term in the flux of circular orbits.
- ▶ What's next: Computing M_{ij} at 4PN and S_{ij} at 3PN.

Ultimately, 4.5PN phase and polarization

Thank you !