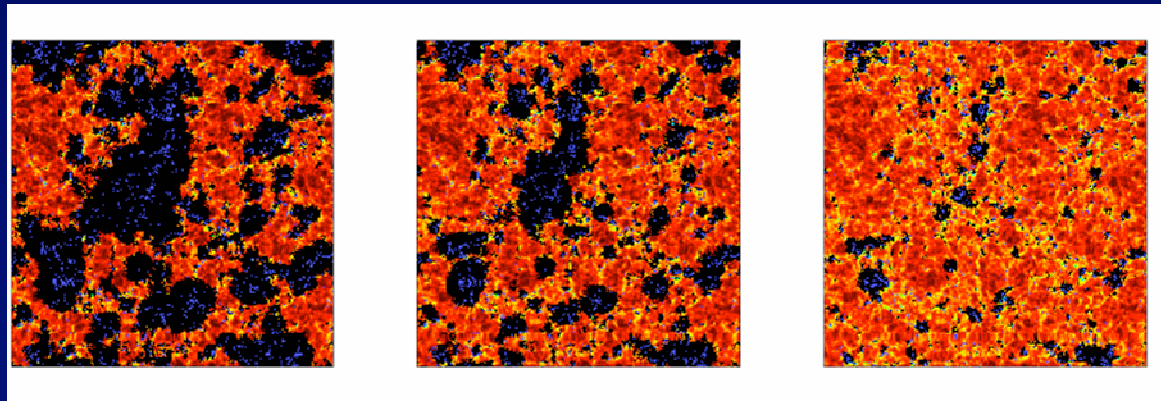


IAP Colloquium, Paris
July 7, 2008

The Physics of Reionization: Common Assumptions

- * *Standard model of particle physics*
- * *Initial conditions from inflation*
- * *Weakly-interacting Cold Dark Matter*



Surprises could signal unexpected new physics

WMAP Cosmological Parameters

Model: Λ cdm

Data: all

$10^2 \Omega_b h^2$	=	$2.19^{+0.06}_{-0.08}$
A	=	$0.67^{+0.04}_{-0.05}$
$A_{0.002}$	=	$0.81^{+0.04}_{-0.05}$
$\Delta_{\mathcal{R}}^2$	=	$(20 \times 10^{-10} \pm 1 \times 10^{-10}) \times 10^{-10}$
$\Delta_{\mathcal{R}}^2(k = 0.002/\text{Mpc})$	=	$(24 \times 10^{-10} {}^{+1 \times 10^{-10}}_{-2 \times 10^{-10}}) \times 10^{-10}$
h	=	$0.71^{+0.01}_{-0.02}$
H_0	=	71^{+1}_{-2} km/s/Mpc
ℓ_A	=	$303.0^{+0.9}_{-1.3}$
n_s	=	$0.938^{+0.013}_{-0.018}$
$n_s(0.002)$	=	$0.938^{+0.012}_{-0.023}$
Ω_b	=	$0.044^{+0.002}_{-0.003}$
$\Omega_b h^2$	=	$0.0220^{+0.0006}_{-0.0008}$
Ω_c	=	$0.22^{+0.01}_{-0.02}$
Ω_Λ	=	0.74 ± 0.02
Ω_m	=	$0.26^{+0.01}_{-0.03}$
$\Omega_m h^2$	=	$0.131^{+0.004}_{-0.010}$
r_s	=	148^{+1}_{-2} Mpc
b_{SDSS}	=	$0.95^{+0.05}_{-0.06}$
σ_8	=	$0.75^{+0.03}_{-0.04}$
$\sigma_8 \Omega_m^{0.6}$	=	$0.34^{+0.02}_{-0.03}$
A_{SZ}	=	$0.78^{+0.23}_{-0.78}$
t_0	=	$13.8^{+0.1}_{-0.2}$ Gyr
τ	=	$0.069^{+0.026}_{-0.029}$
θ_A	=	0.594 ± 0.002 °
z_{eq}	=	3135^{+85}_{-159}
z_γ	=	$9.3^{+2.8}_{-2.0}$

The initial conditions of the Universe can be summarized on a single sheet of paper, yet thousands of books cannot fully describe the complex structures we see today...

THE DARK AGES of the Universe

Astronomers are trying to fill in the blank pages in our photo album of the infant universe

By Abraham Loeb

When I look up into the sky at night, I often wonder whether we humans are too preoccupied with ourselves. There is much more to the universe than meets the eye on earth. As an astrophysicist I have the privilege of being paid to think about it, and it puts things in perspective for me. There are things that I would otherwise be bothered by—my own death, for example. Everyone will die sometime, but when I see the universe as a whole, it gives me a sense of longevity. I do not care so much about myself as I would otherwise, because of the big picture.

Cosmologists are addressing some of the fundamental questions that people attempted to resolve over the centuries through philosophical thinking, but we are doing so based on systematic observation and a quantitative methodology.

Perhaps the greatest triumph of the past century has been a model of the universe that is supported by a large body of data. The value of such a model to our society is sometimes underappreciated. When I open the daily newspaper as part of my morning routine, I often see lengthy descriptions of conflicts between people about borders, possessions or liberties. Today's news is often forgotten a few days later.

But when one opens ancient texts that have appealed to a broad audience over a longer period of time, such as the Bible, what does one often find in the opening chapter? A discussion of how the constituents of the universe—light, stars, life—were created. Although humans are often caught up with mundane problems, they are curious about the big picture. As citizens of the universe we cannot help but wonder how the first sources of light formed, how life came into existence and whether we are alone as intelligent beings in this vast space. Astronomers in the 21st century are uniquely positioned to answer these big questions.

What makes modern cosmology an empirical science is that we are literally able to peer into the past. When you look at your image reflected off a mirror one meter

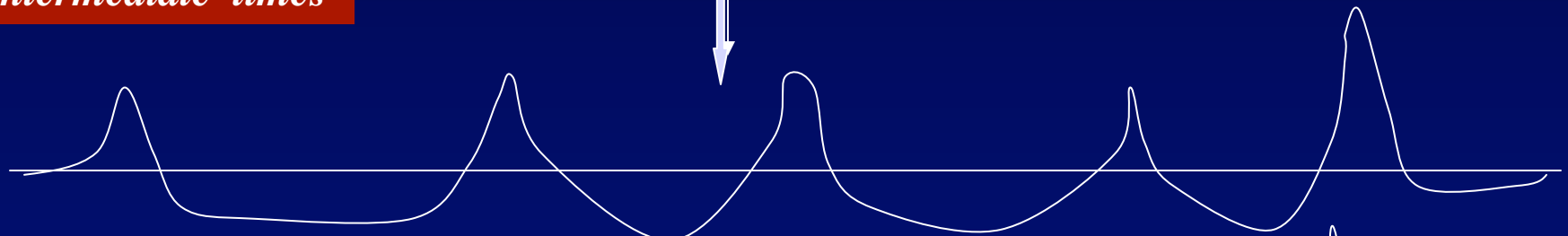
On small scales the universe is clumpy

Early times

Density perturbation

*Mean
Density*

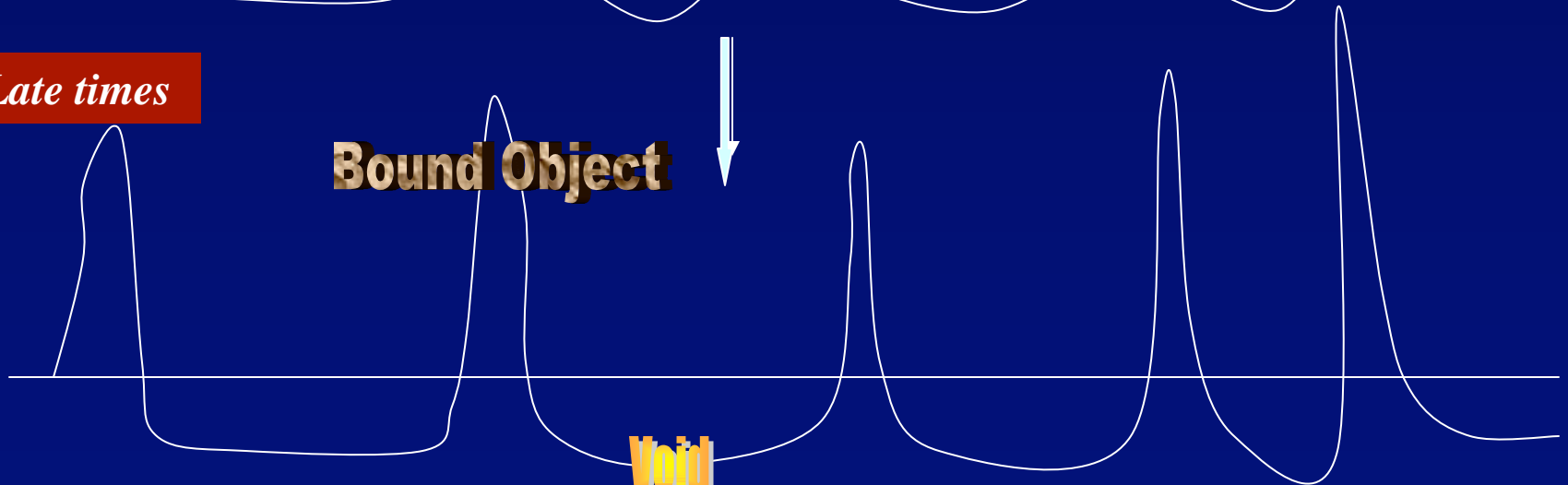
Intermediate times



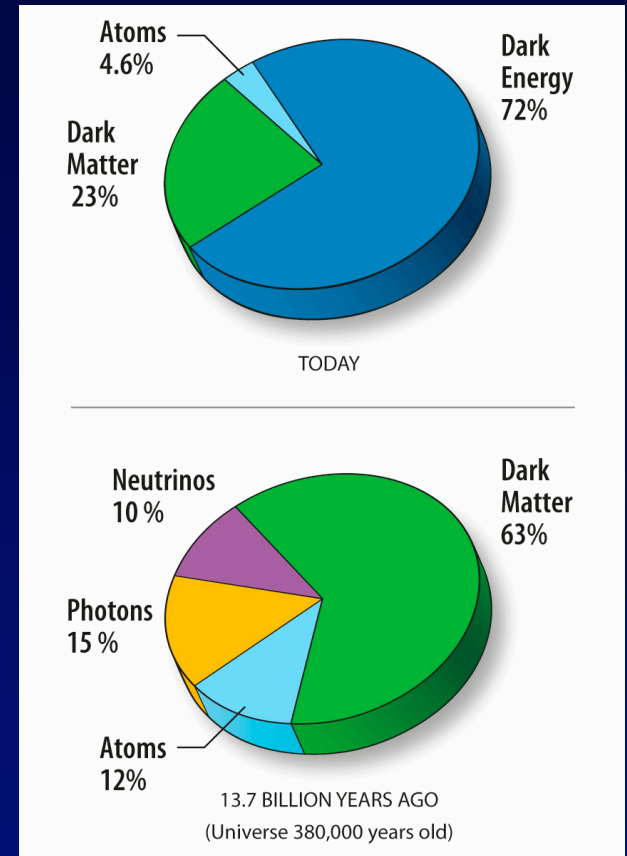
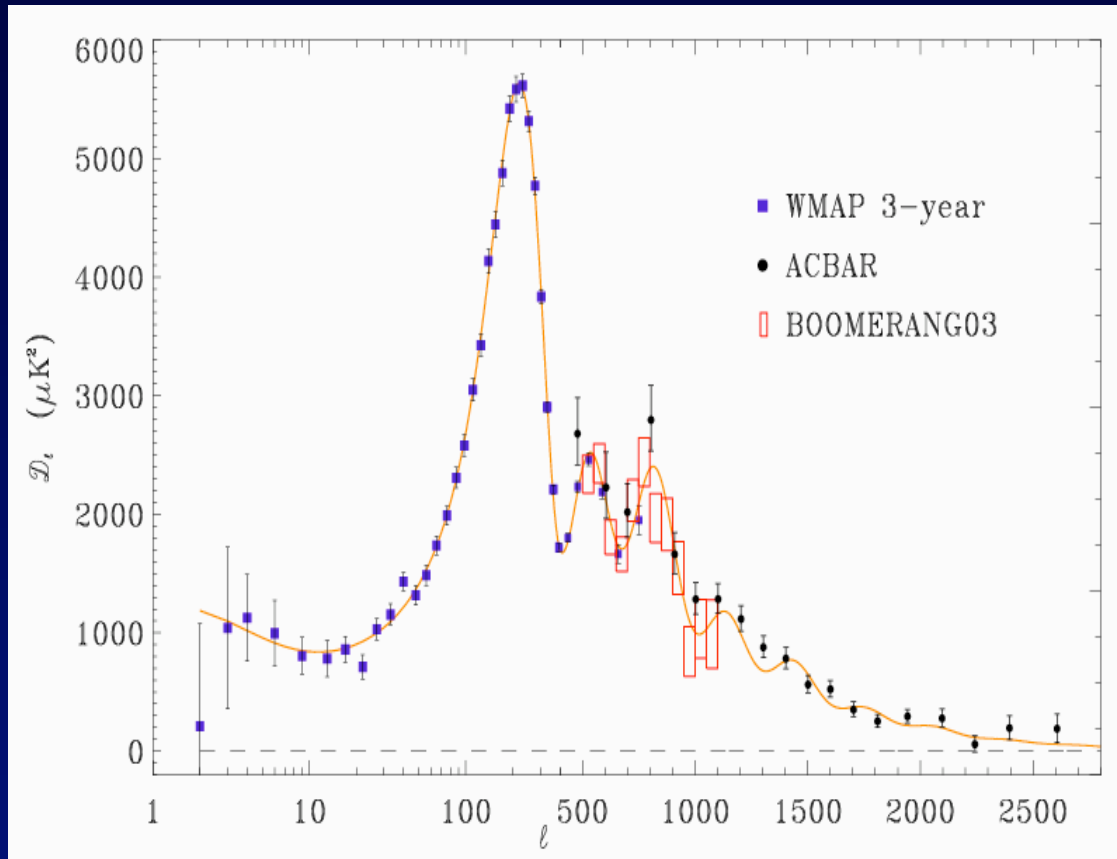
Late times

Bound Object

Void



Evidence that Most Matter is EM Dark



Diffusion damping of small-scale fluctuations in the baryon-photon fluid prior to cosmic recombination implies that galaxies could not have formed in our Universe without dark matter!

The First Dwarf Galaxies Form at $z \sim 30$

The distribution of matter can be mapped through:

- (i) Surveys of galaxies*
- (ii) Surveys of the diffuse (intergalactic) gas*

*molecular hydrogen in Jeans mass objects
($\sim 10^5 M_{\odot}$)*

Yoshida et al. 2003

Close-up images of some of the most distant galaxies in the Hubble Ultra Deep Field

Galaxies at very early times tend to be very small and often show signs of interactions. The HUDF contains nearly 50 galaxies at redshifts 5–6, compared to a few tentative identifications in earlier, shallower observations.

galaxies $z \sim 3-4$ Lookback time 11.4–12 billion years (312 objects)



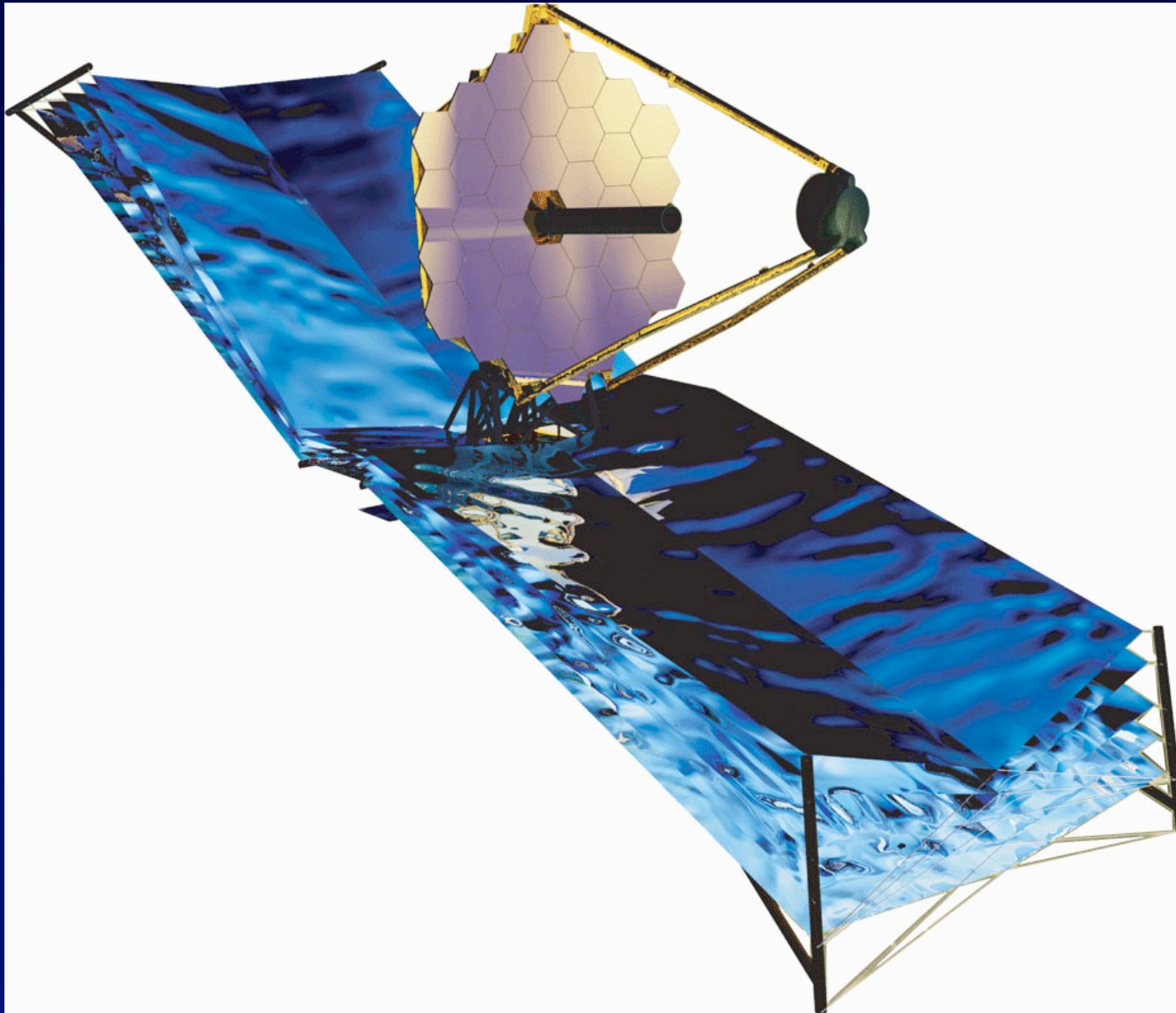
galaxies $z \sim 4-5$ Lookback time 12–12.3 billion years (79 objects)



galaxies $z > 5$ Lookback time 12.3–12.6 billion years (45 objects)



Searching for the First Galaxies: *James Webb Space Telescope*



*Mirror diameter: 6.5
meter*

Material: beryllium

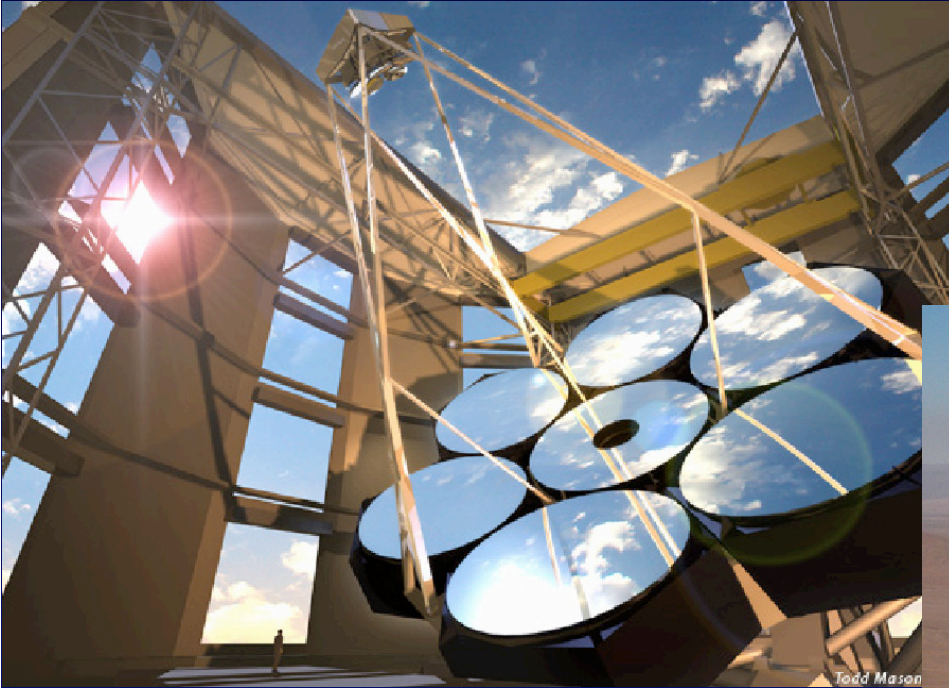
18 segments

*Wavelength coverage:
0.6-28 micron*

L2 orbit

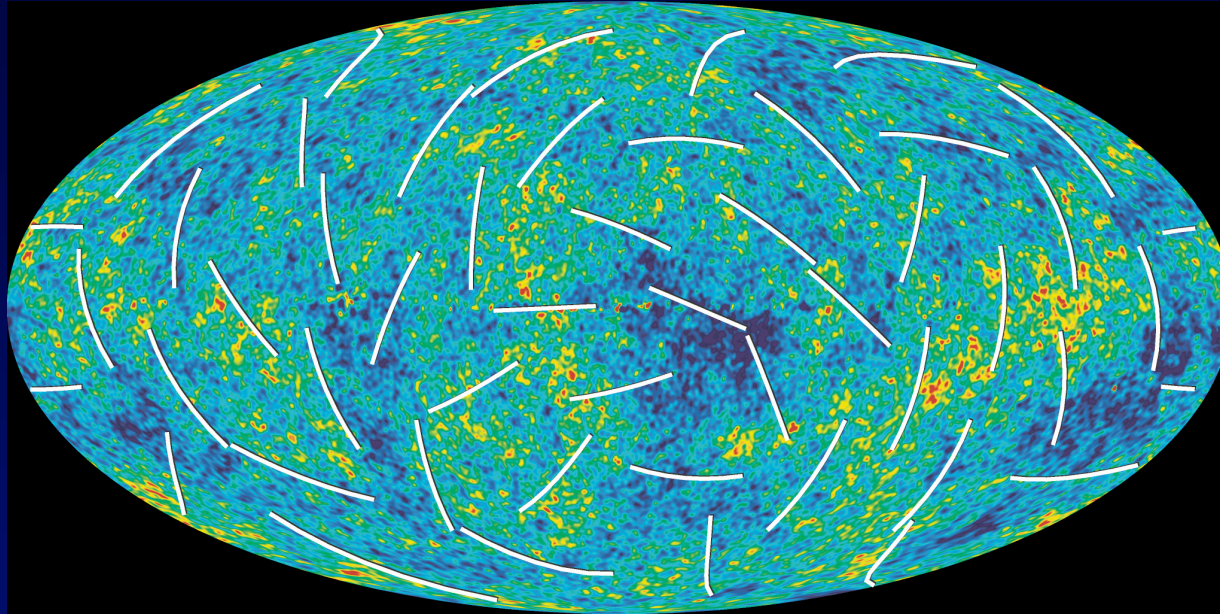
Launch date: 2013

Extremely Large Telescopes (20-40 meters)



- GMT=Seven mirrors, each 8.4m in diameter
- TMT, EELT – segmented 20-40m aperture

Cosmic Microwave Background (*WMAP5*)

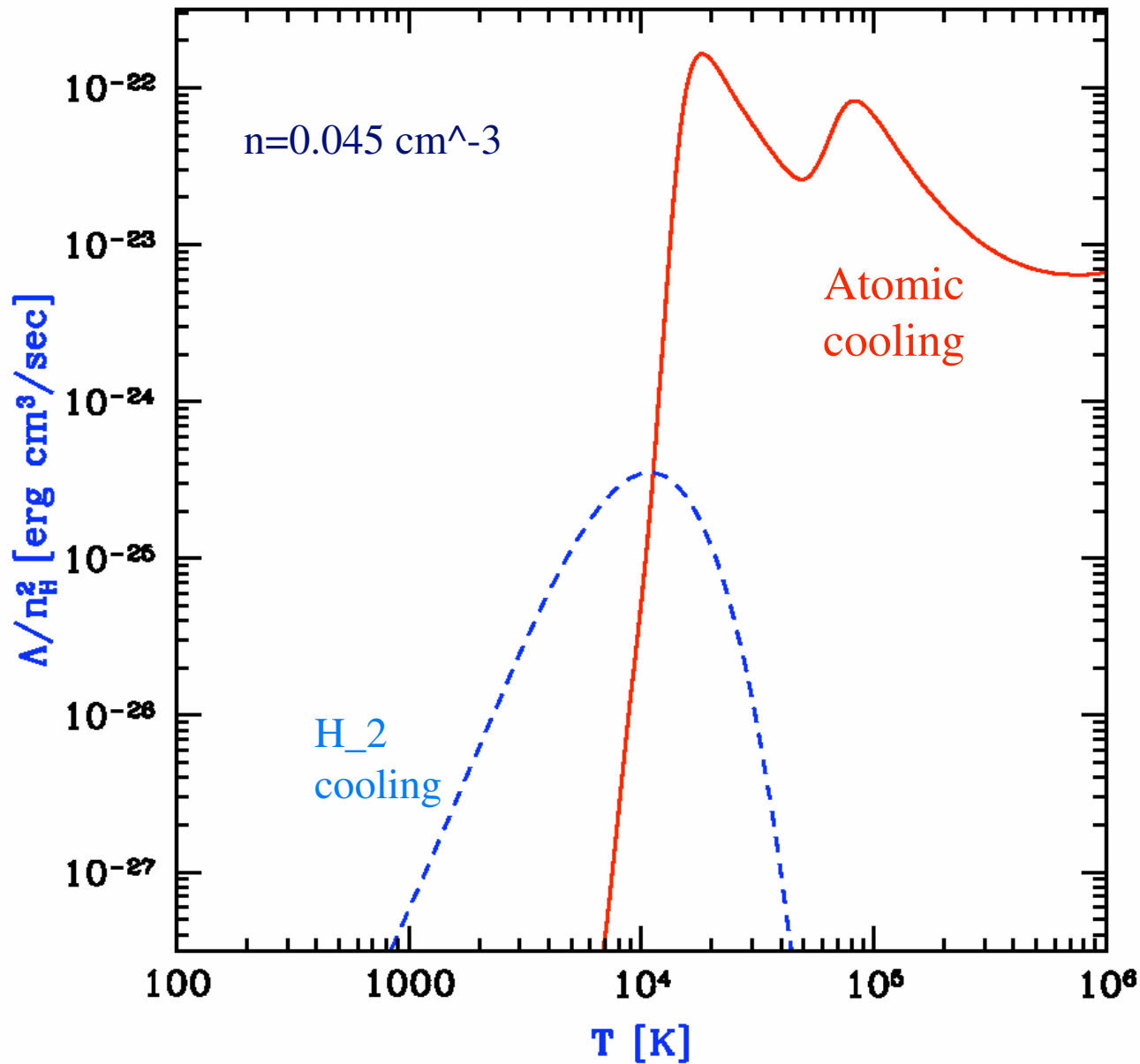


$$\tau = 0.09 \pm 0.02$$

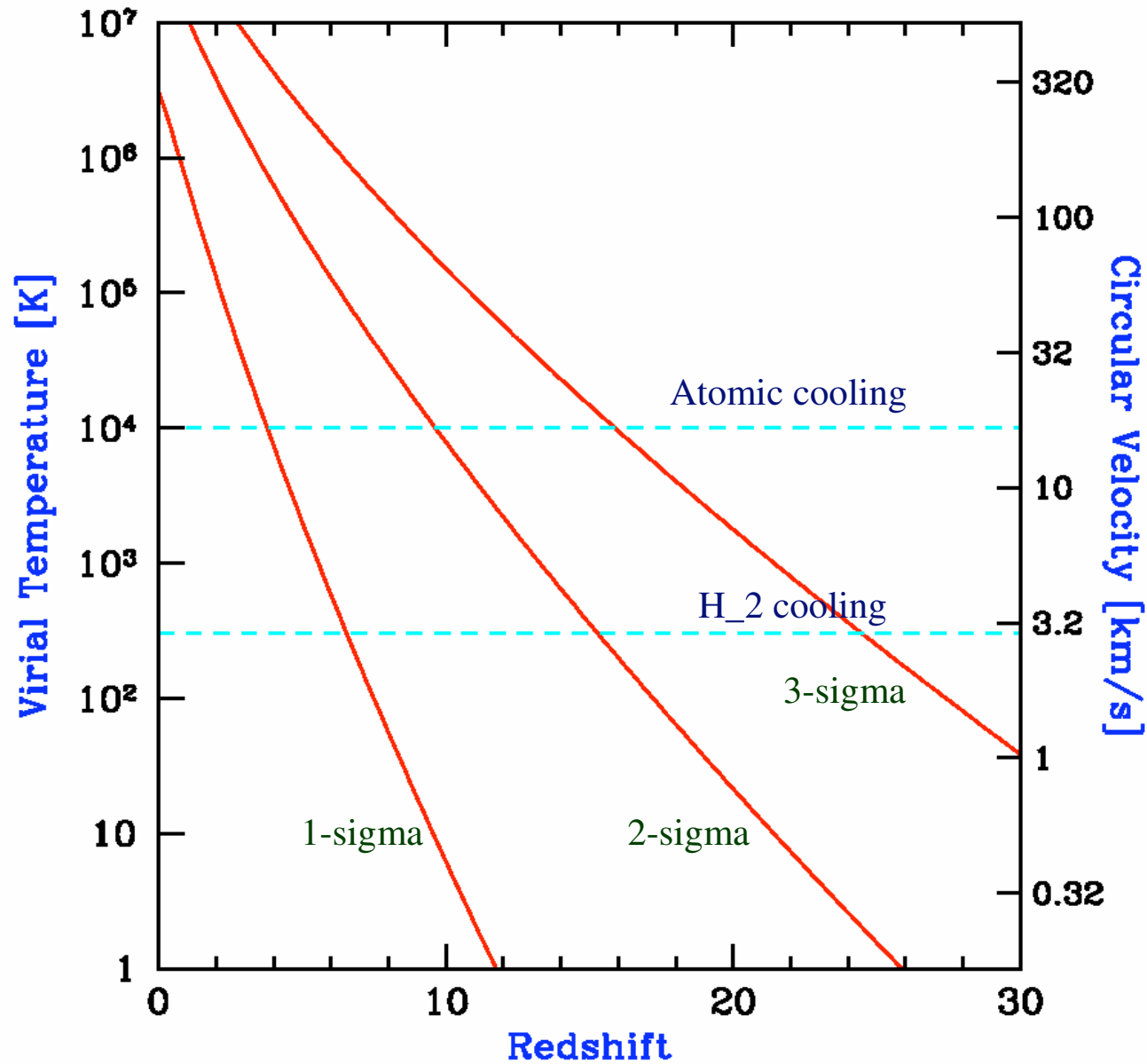
The polarization data indicates that the first stars must have formed 400 million years after the big bang, when the universe was only a few percent of its current age!

Dunkley et al. 2008

Cooling Rate of Primordial Gas



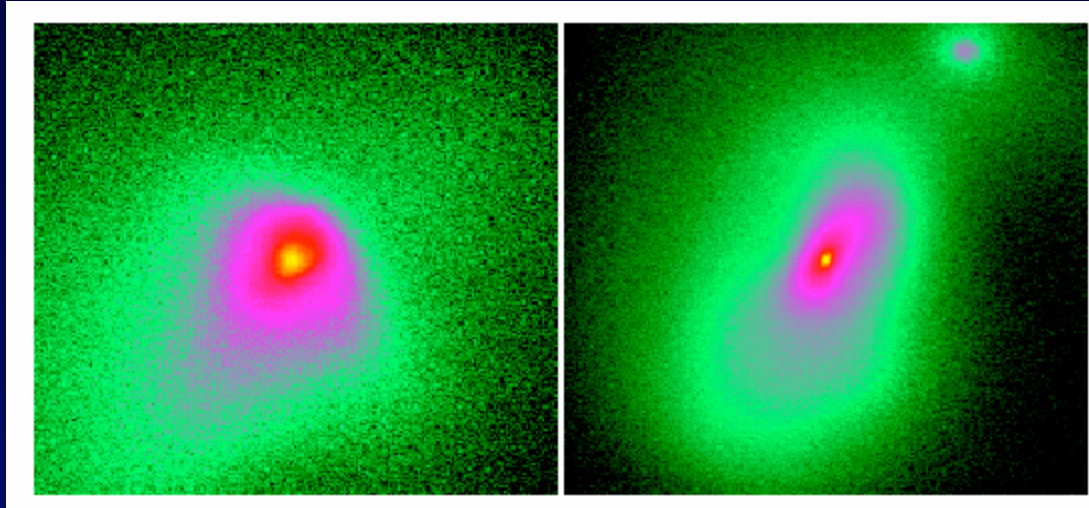
Virial Temperature of Halos



Massive Accretion by Pop-III Proto-Stars

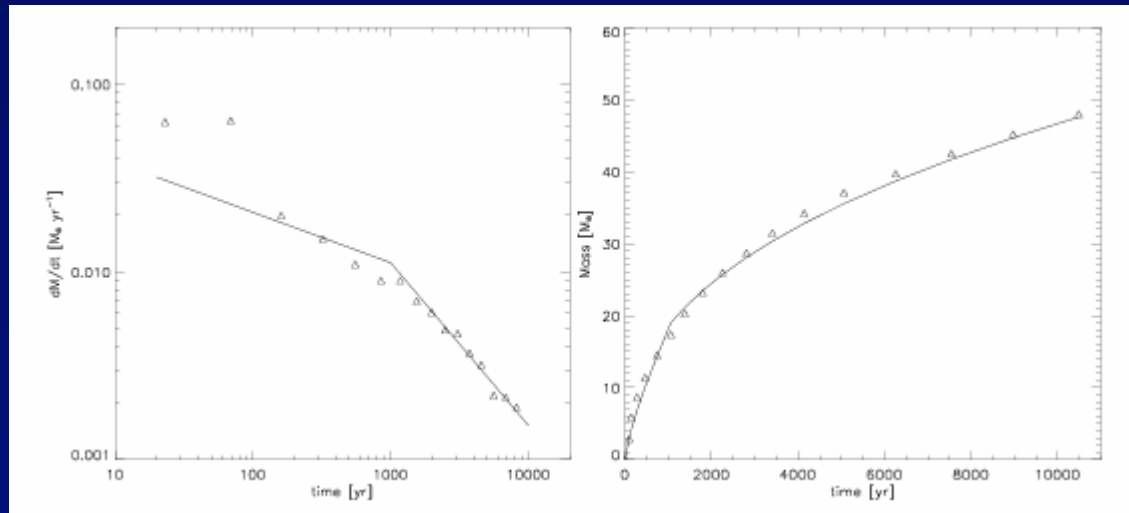
Bromm & Loeb,
astro-ph/0312456

*Resolving
accretion
flow down
to ~ 0.03 pc*



23.5pc

0.5pc



$$\dot{M} \sim c_s^3 / G$$

$T_{\min} \sim 200\text{K}$
for : H_2 – cooling

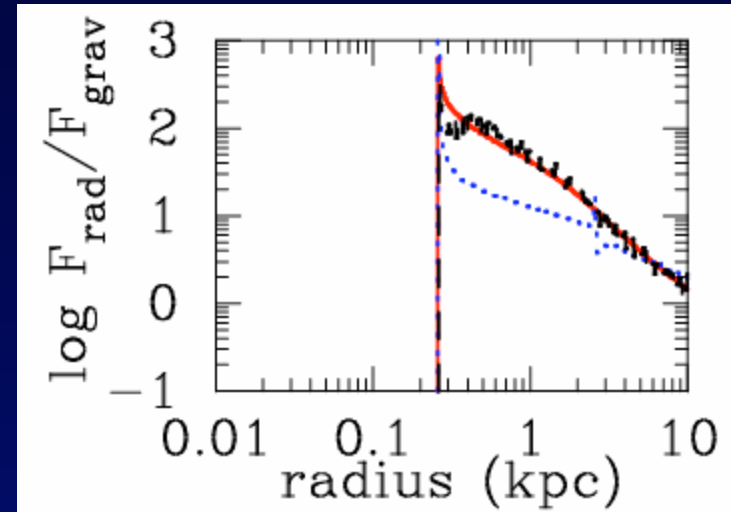
Final stellar mass is feedback limited (radiation, wind)

Outflows Driven by Ly α Radiation Pressure

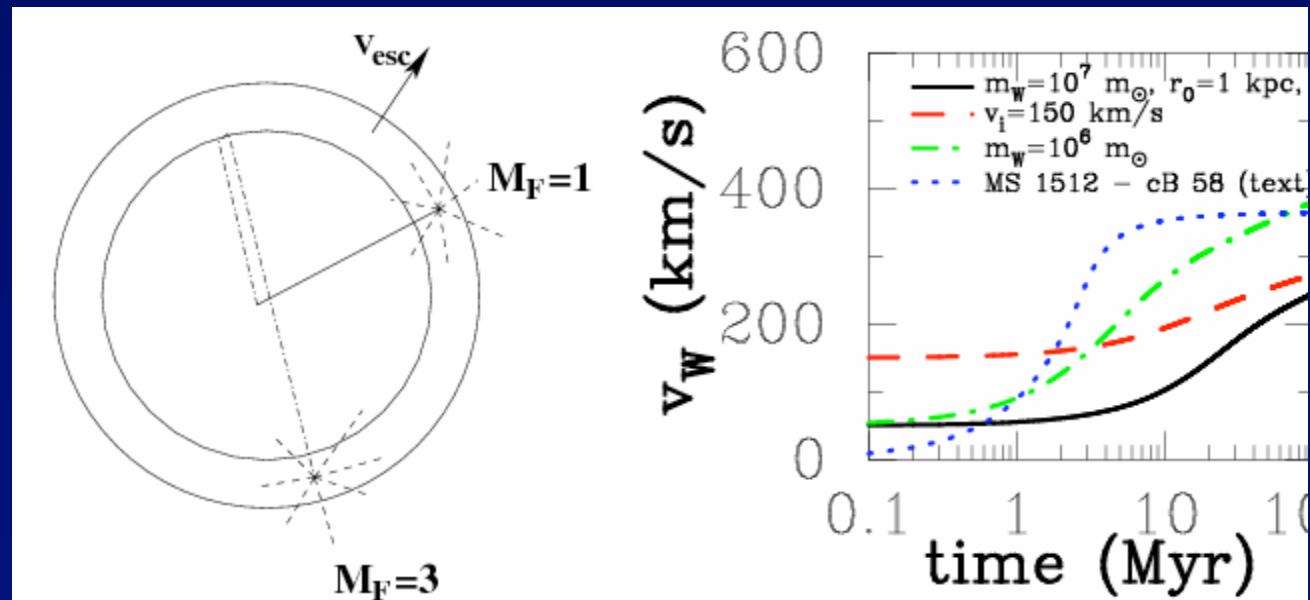
- IGM around the mini-halos hosting the first stars

$$M_{\star} = 100M_{\odot}$$

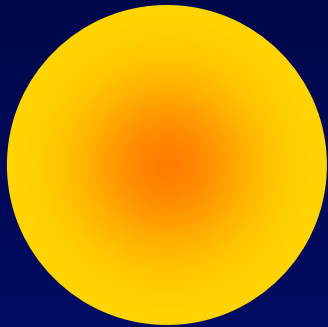
$$M_{\text{halo}} = 10^6M_{\odot}$$



- Supershells around starburst galaxies



Number of ionizing photons ($>13.6\text{eV}$) per baryon incorporated into stars:



Massive, metal free stars

$$T_{\text{eff}} \sim 10^5 \text{K}$$
$$L = L_E \propto M$$

$M > 300M_{\odot}$	$\sim 100,000$
$M > 100M_{\odot}$	$\sim 40,000$

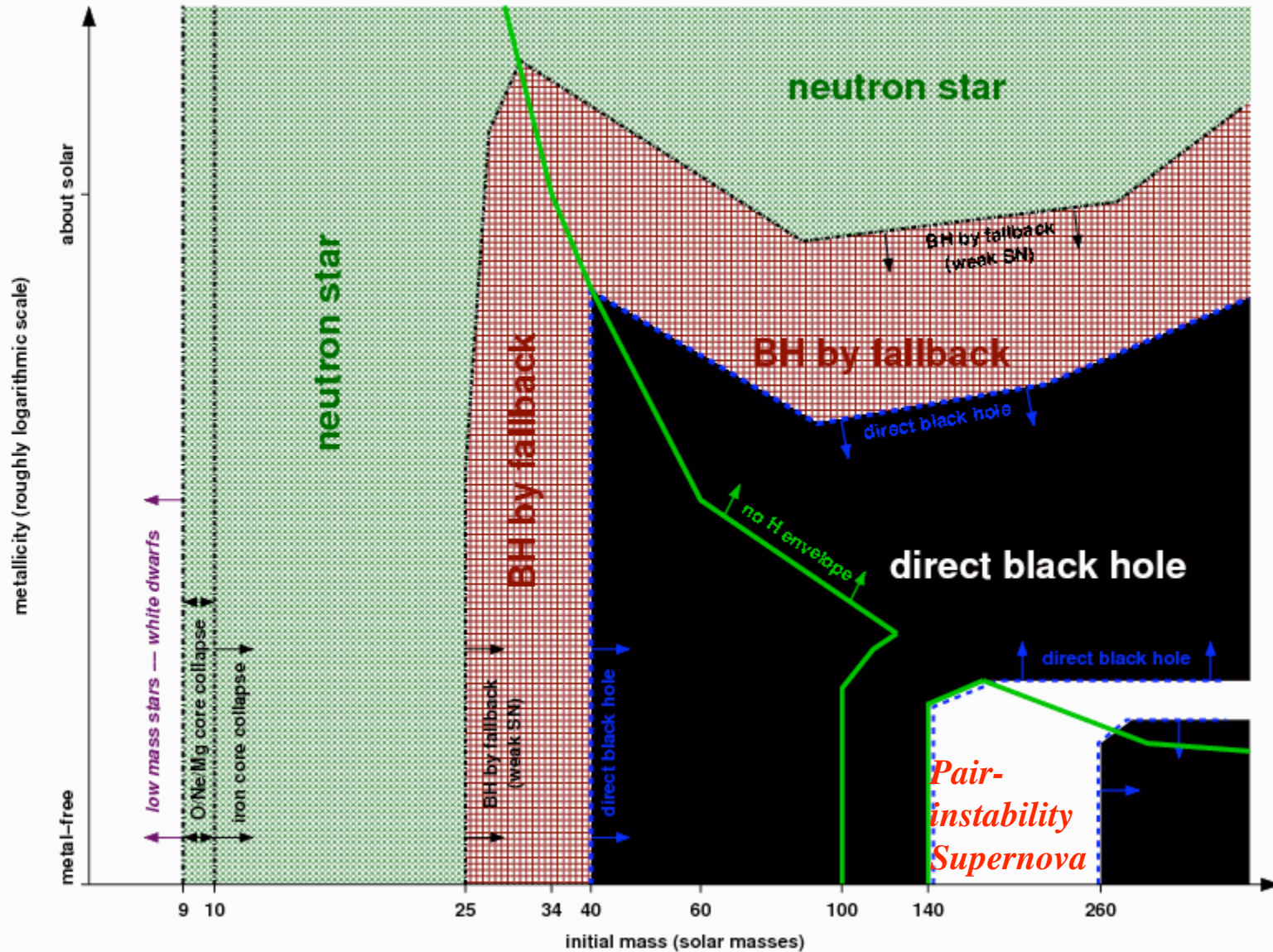
Gain by up to a factor of ~30!



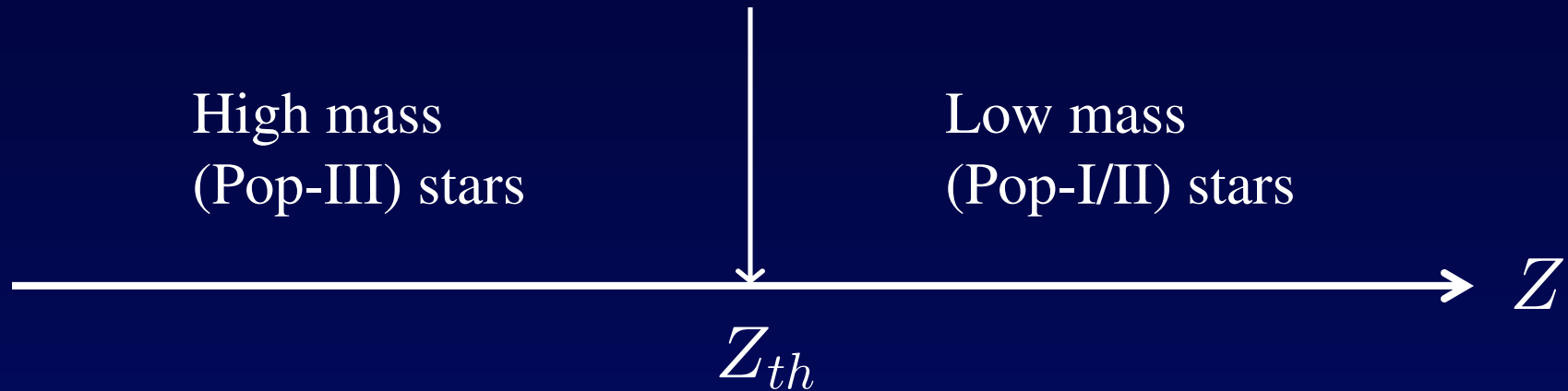
Salpeter mass function

Metal free	$\sim 7,000$
$Z = 0.01Z_{\odot}$	$\sim 3,500$

How do massive stars end their life?



Threshold Metallicity



Atomic cooling (CII, OI) → $Z_{th} \sim 10^{-3} Z_{\odot}$

Dust, molecules (CO) → $Z_{th} < 10^{-5} Z_{\odot}$

Milky-Way halo metal-poor stars: $Z_{Fe} \sim 10^{-5} Z_{\odot}$

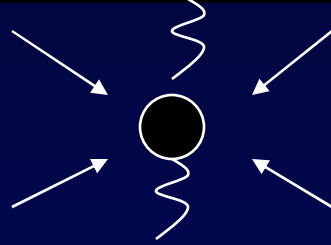
but: $Z_{C,O} > 10^{-3} Z_{\odot}$

Early Galaxies: *Parametrizing Our Ignorance*

Star formation

- Mass function of dark matter halos: *N-body simulations*.
- Minimum halo mass for star formation:
 - 200K – H₂ cooling
 - 10⁴ K -- atomic H cooling
 - 10⁵ K -- assembly of gas from a photo-ionized IGM
- Star formation efficiency: $f_{\star} \sim 10\%$
- # of ionizing photons/baryon in stars:
 - ~4000 – Pop I/II; ~ 10⁵ -- Pop III (metal free)

Nuclear Black Holes



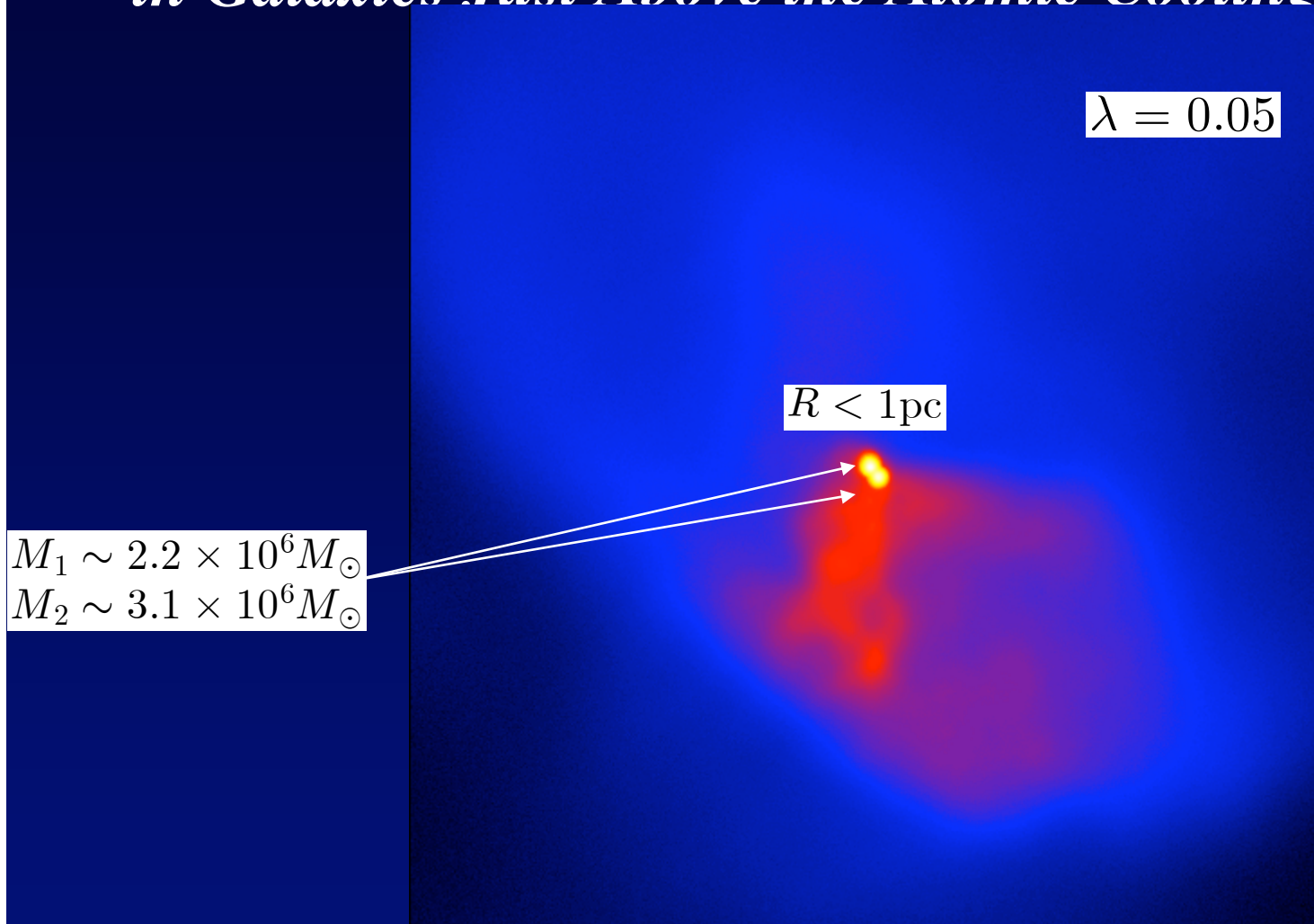
$$L = \epsilon \dot{M} c^2 \quad t_E = M / \dot{M} = 4 \times 10^7 \frac{(\epsilon/10\%)}{(L/L_E)} \text{years}$$

$$M \propto \exp\{t/t_E\}$$

Stellar mass seed requires ~billion years to grow to an SDSS quasar ($10^9 M_\odot$)

...But a billion year is the Hubble time at $z \sim 6$, and feedback from star formation and quasar activity is likely to suppress continuous accretion, and also...

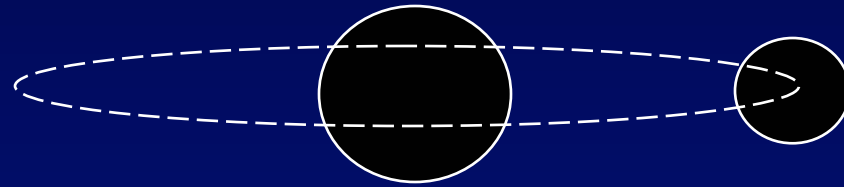
Massive Black Hole Seeds: *Suppressed Fragmentation in Galaxies Just Above the Atomic Cooling Threshold*



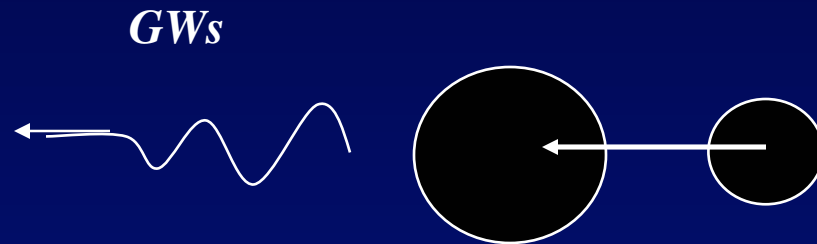
Unusual environments: H_2 suppressed ; binary black holes may form –LISA sources

Low-spin systems: Eisenstein & Loeb 1995; Numerical simulations: Bromm & Loeb 2002

Gravitational Wave Recoil

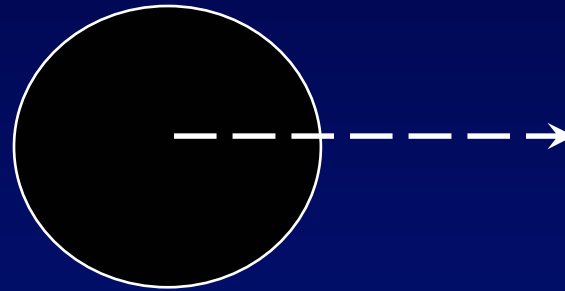


Gravitational Wave Recoil



*Anisotropic emission of gravitational waves →
momentum recoil*

Gravitational Wave Recoil



Recoil speed (\sim tens-4000 km/s) is independent of remnant black hole mass \rightarrow *low-mass halos may easily lose their low-mass seeds after several mergers*

*First Galaxies Were Strongly Clustered on
Scales of up to ~100 comoving Mpc*

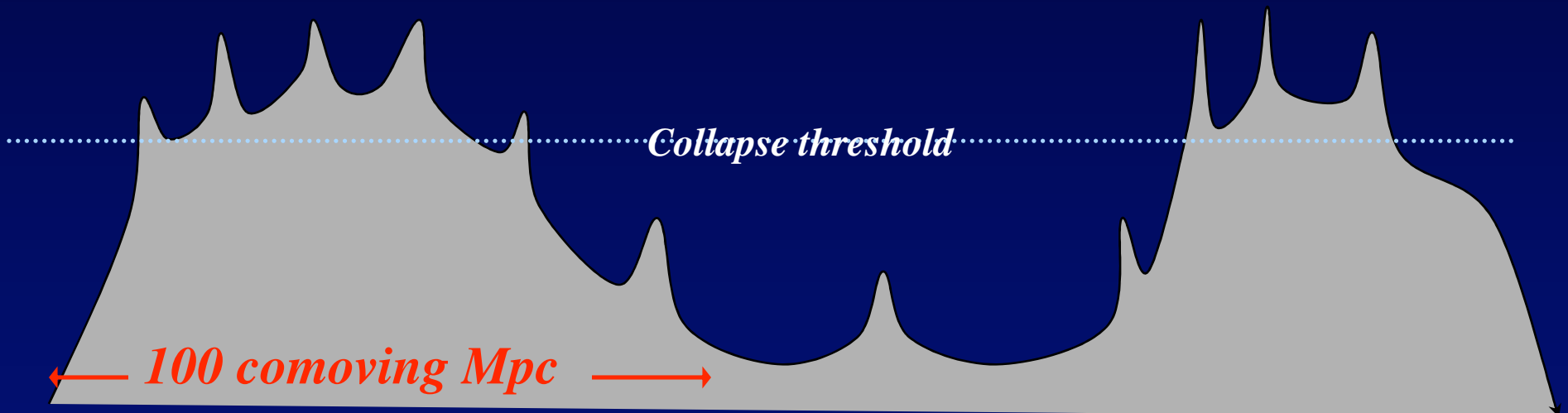
$z=20$

Collapse threshold



First Galaxies Were Strongly Clustered on Scales of up to ~100 comoving Mpc

$z=10$



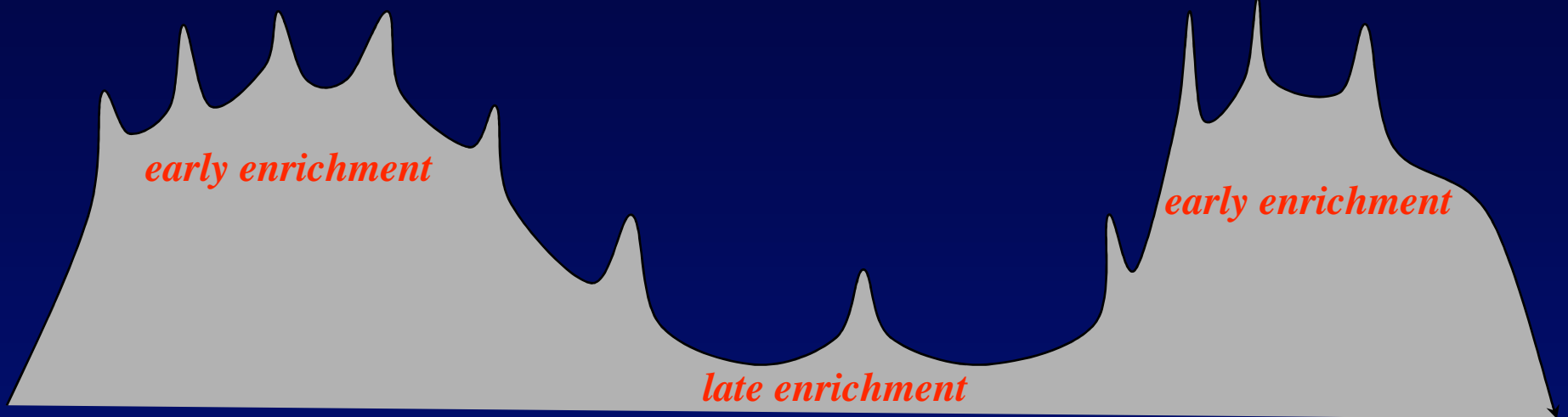
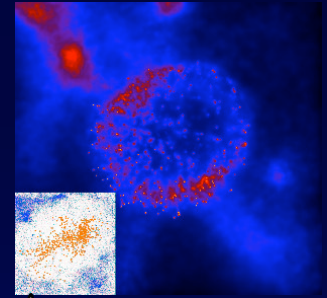
Challenges for numerical simulations of reionization:

**Resolving dwarf galaxies as sources of ionizing photons*

**Simulation box >100 comoving Mpc on a side*

**Following gravity, hydrodynamics, radiative transfer and their interaction*

→ Enrichment of Primordial Gas with Heavy Elements was Highly Inhomogeneous



early Pop-III

late Pop-III

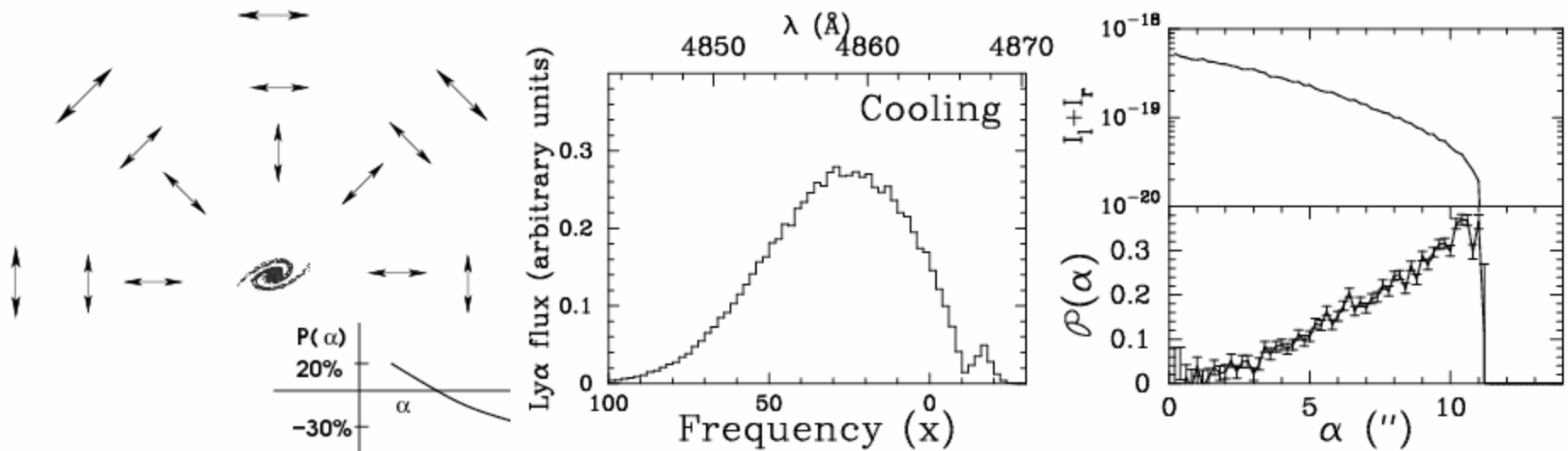
early Pop-III

Searches for high- z Galaxies:

- Lyman-break
- $\text{Ly}\alpha$
- Other lines ($\text{H}\alpha$, CO, CII, OI, He)

A future frontier: polarized $\text{Ly}\alpha$ halos

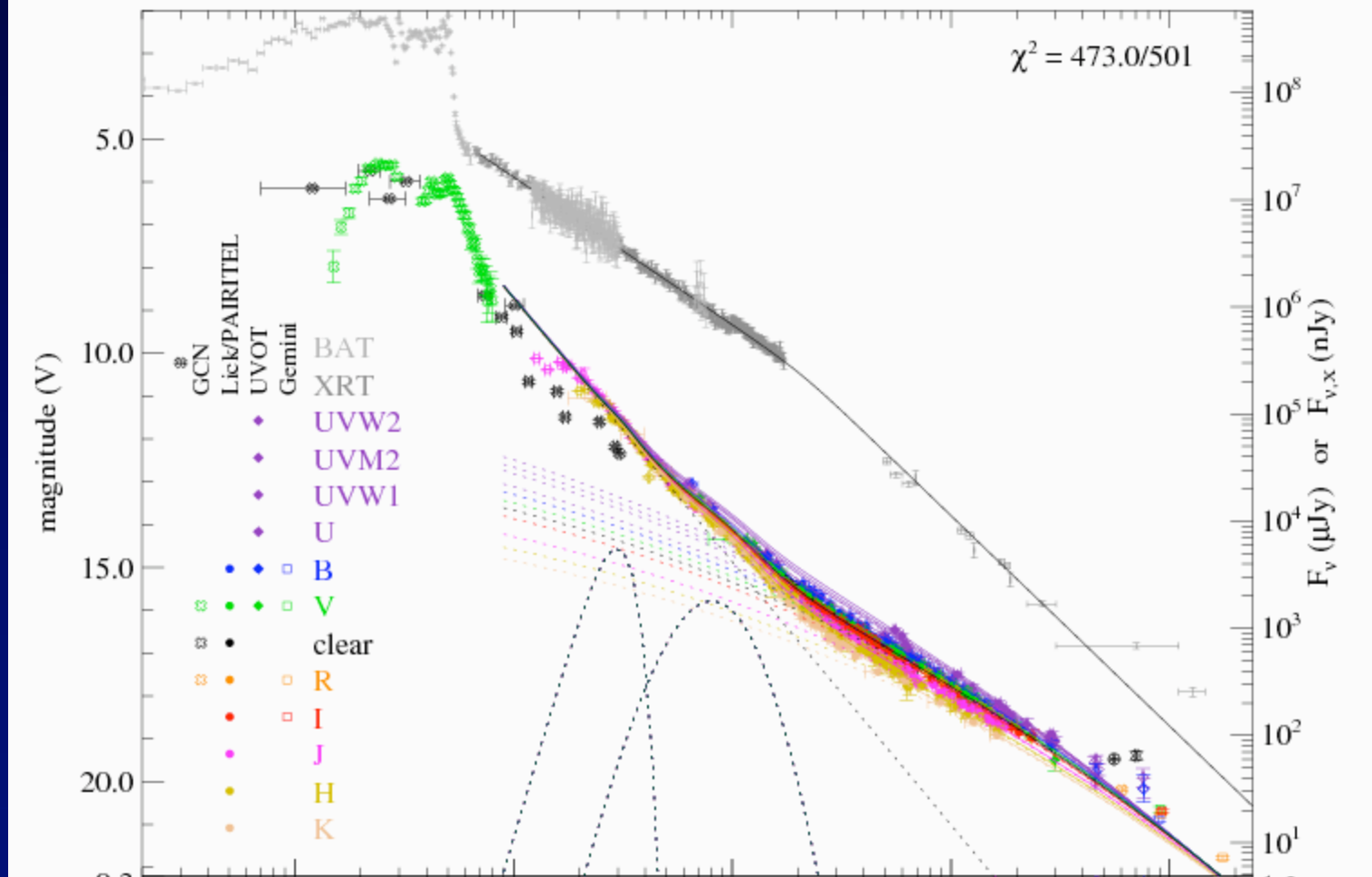
Collapsing gas cloud



Rybicki & Loeb 1999; Dijkstra & Loeb arxiv:0711.2312

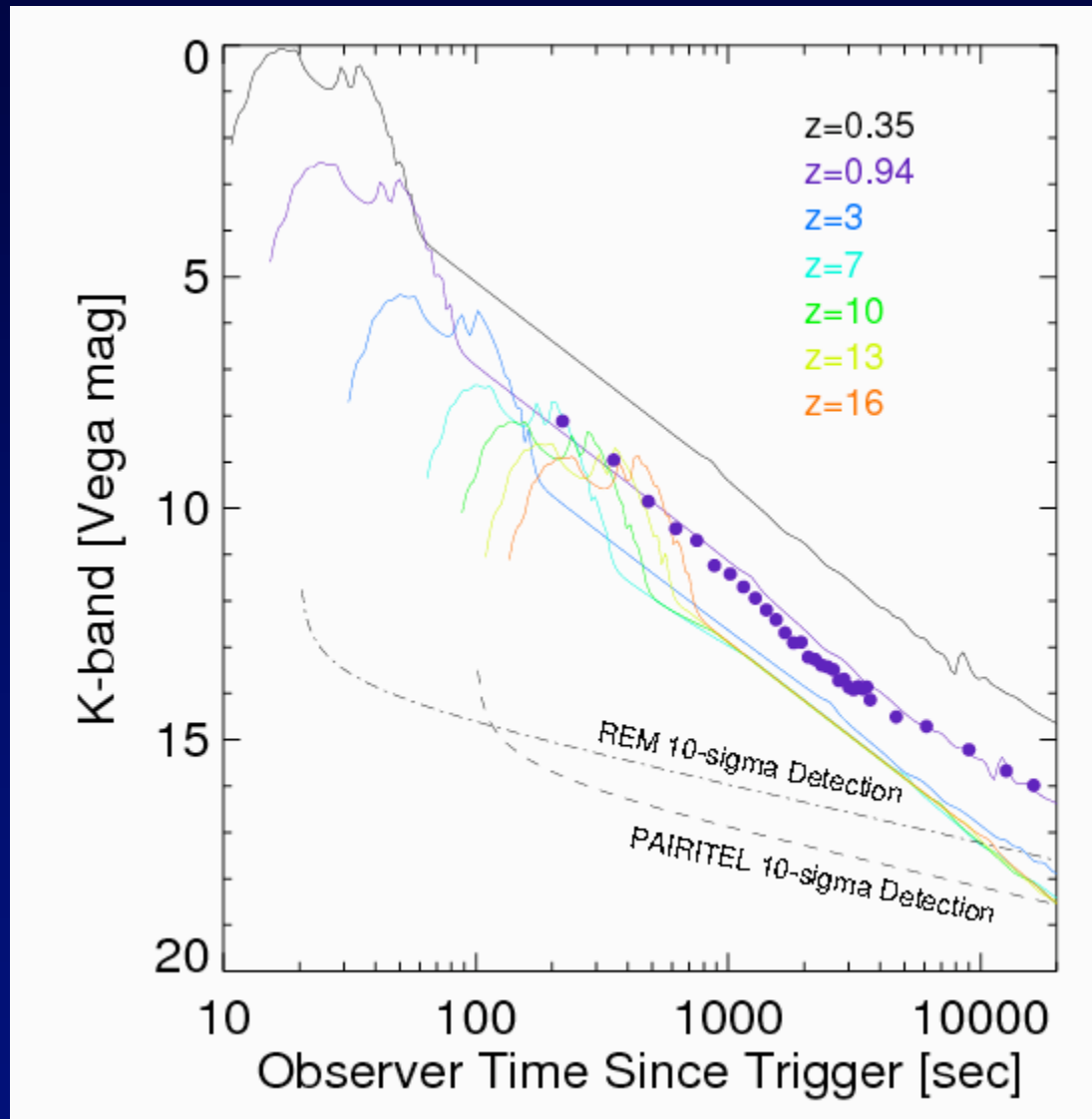
Long Gamma-Ray Bursts: Observing One Star at a Time

GRB080319B – Visible to the naked eye at the edge of the Universe



Existing finder: Swift; Proposed: JANUS, EXIST (high-z GRBs)

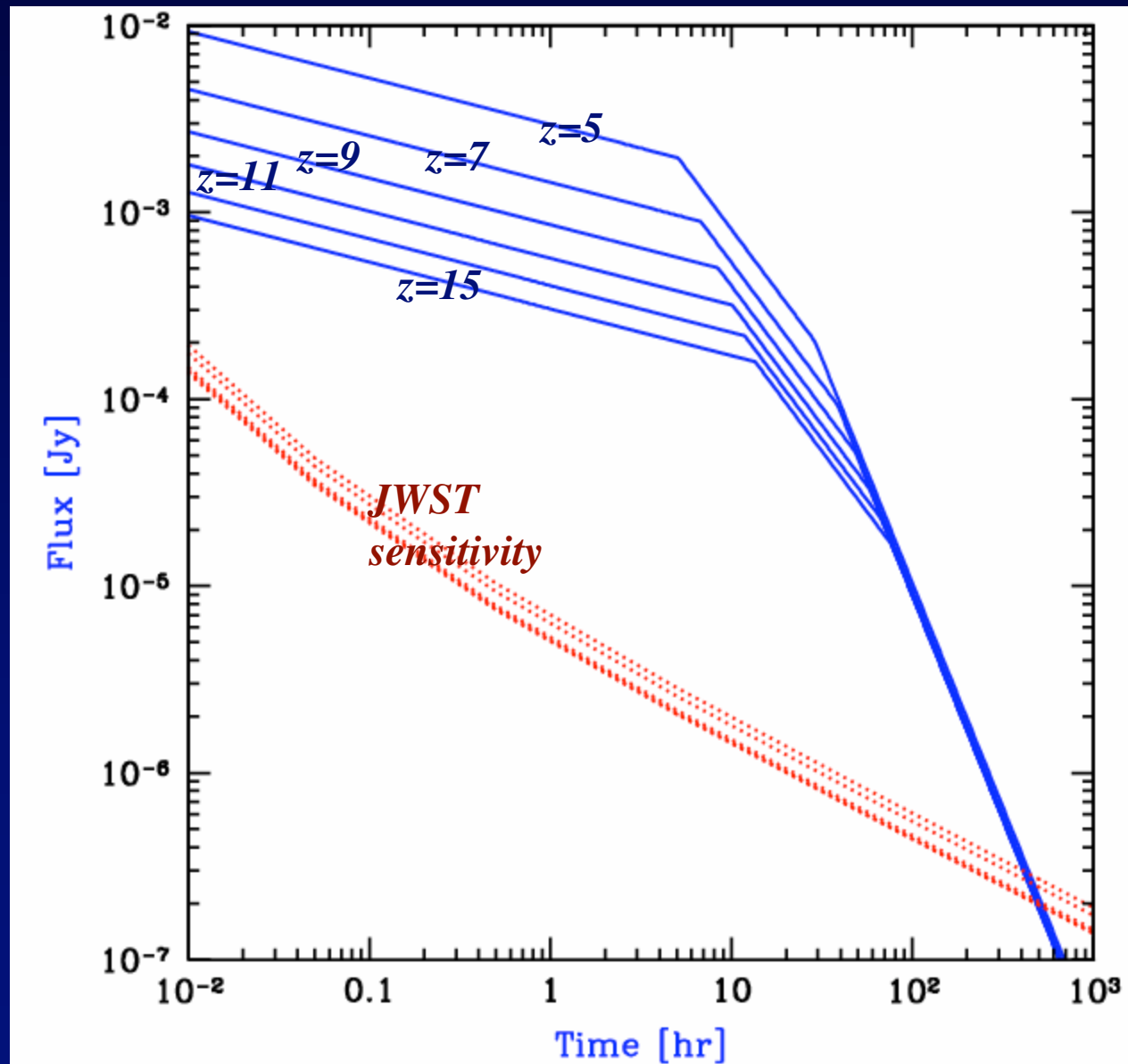
GRB080319B at higher redshifts



Bloom et al.
arXiv:0803.3215

Detectability of Afterglow Emission Near the Ly α Wavelength

Photometric redshift identification: based on the Ly α trough

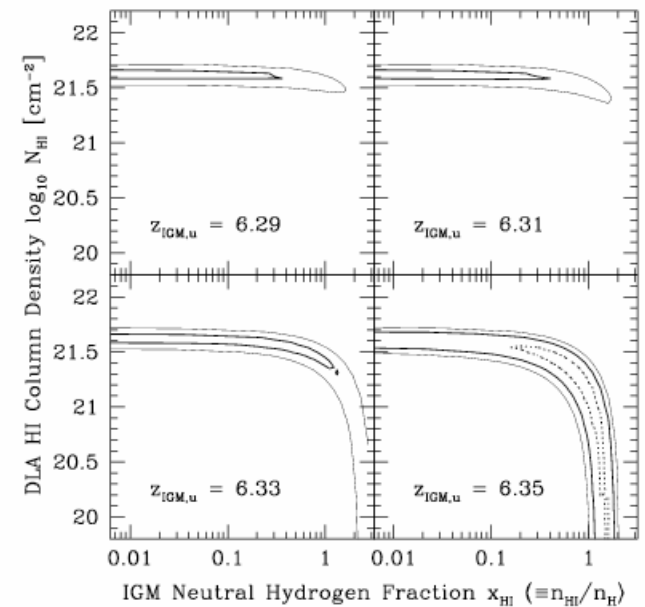
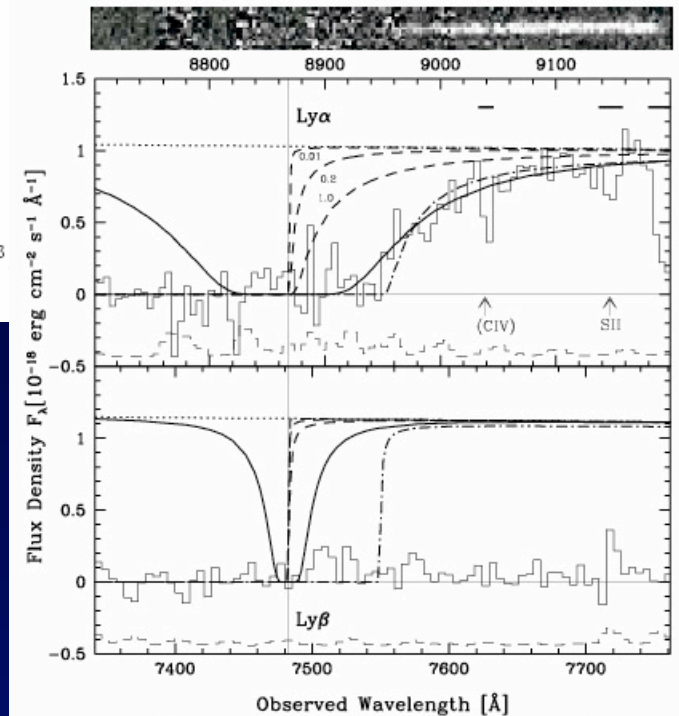
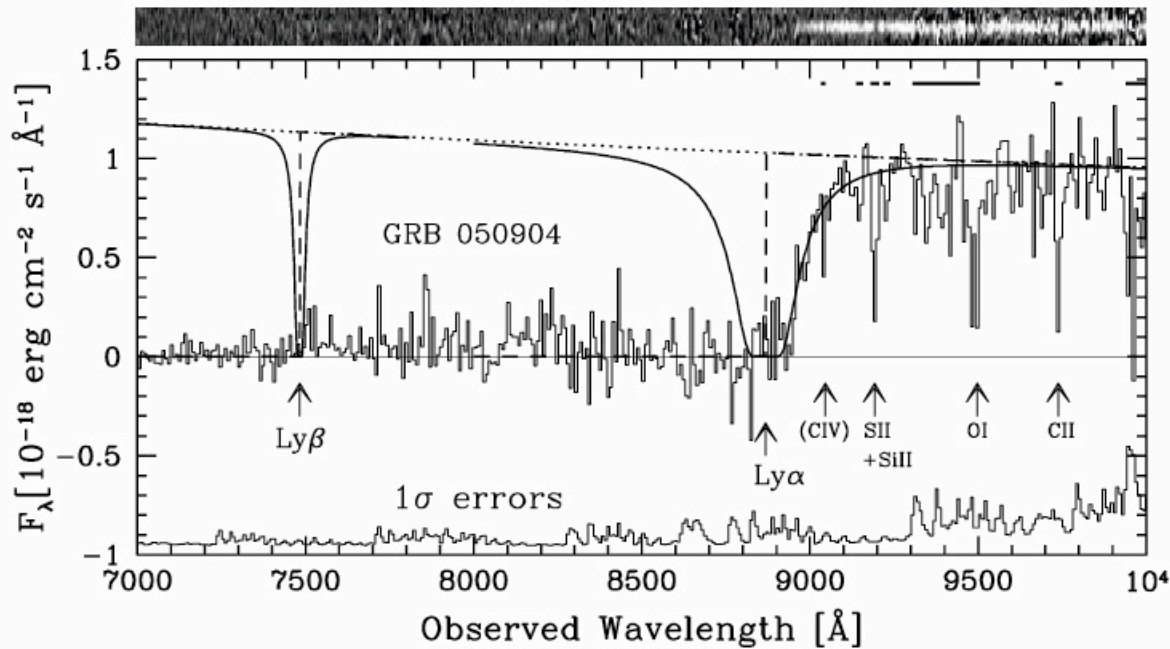


*Barkana &
Loeb 2003*

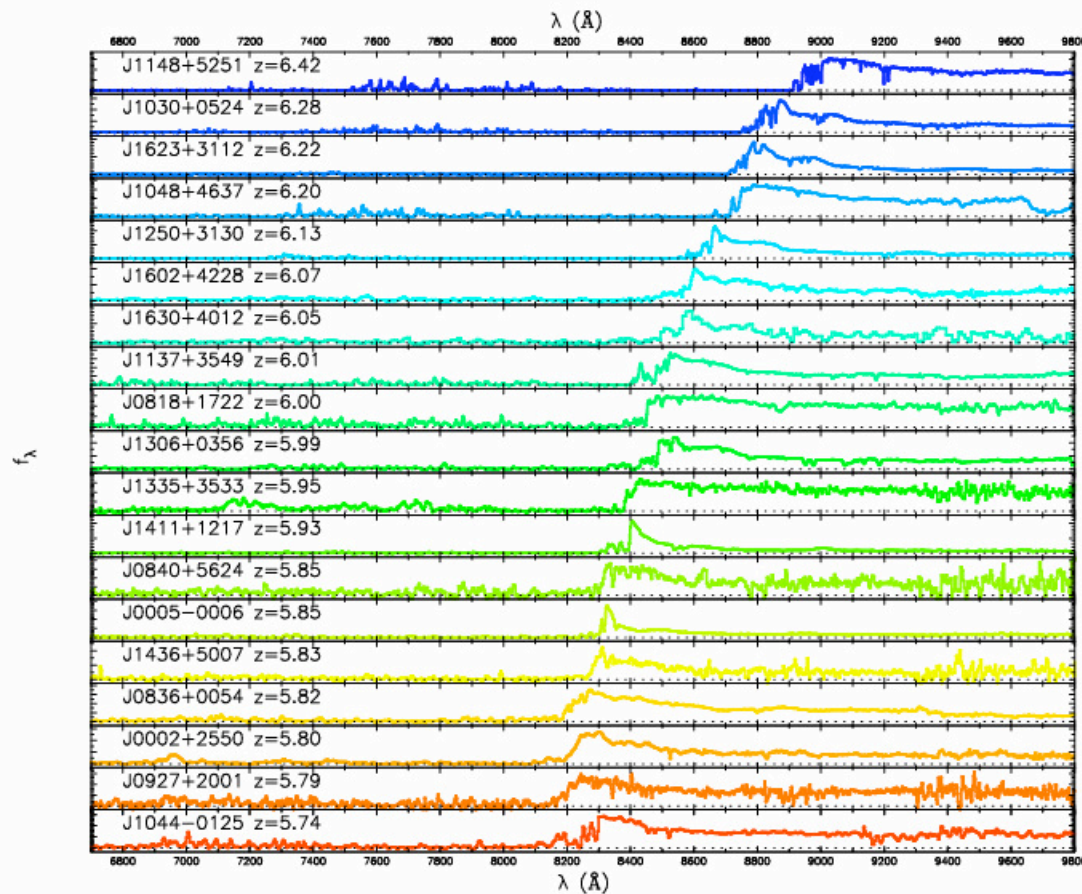
astro-ph/0305470

Implications for the Cosmic Reionization from the Optical Afterglow Spectrum of the Gamma-Ray Burst 050904 at $z = 6.3^*$

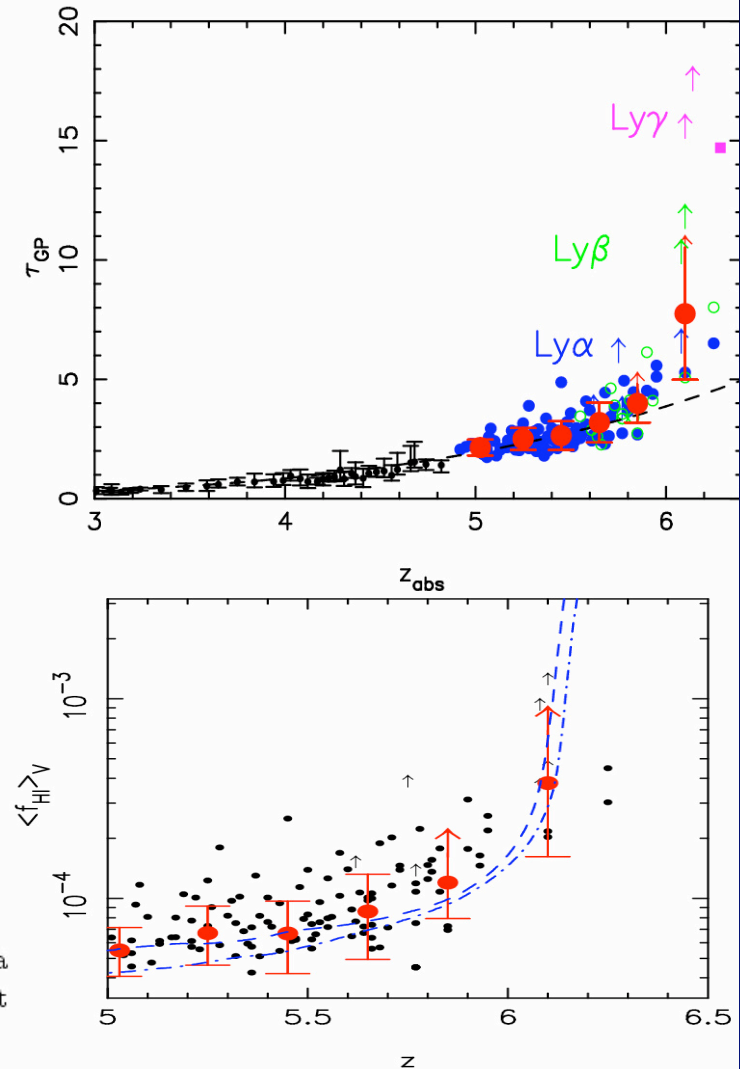
Tomonori TOTANI¹, Nobuyuki KAWAI², George KOSUGI³, Kentaro AOKI⁴, Toru YAMADA³
 Masanori IYE³, Kouji OHTA¹, and Takashi HATTORI⁴



So far, the hydrogen was only probed by quasars



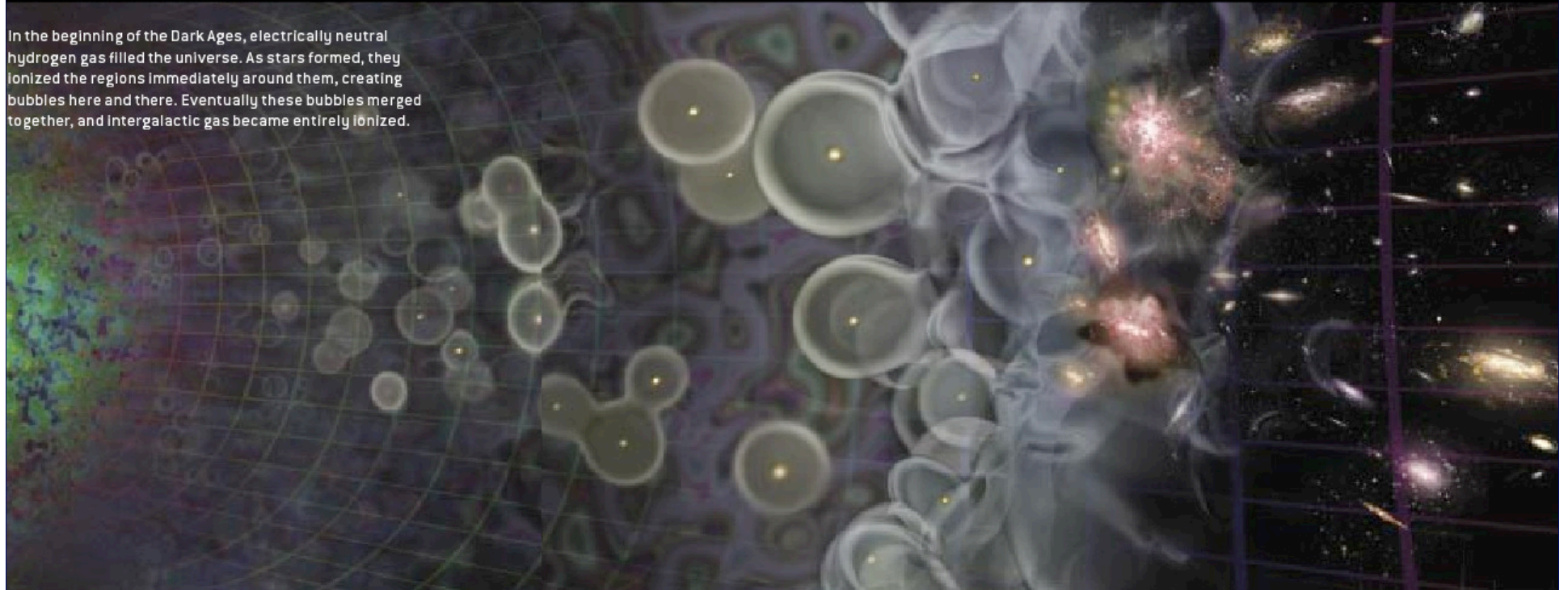
Spectra of our sample of nineteen SDSS quasars at $5.74 < z < 6.42$. Twelve of the spectra were taken with Keck/ESI, while the others were observed with the MMT/Red Channel and Kitt Peak 4-meter/MARS spectrographs. See Table 1 for detailed information.



21cm Mapping of Cosmic History

LIGHTING UP THE COSMOS

In the beginning of the Dark Ages, electrically neutral hydrogen gas filled the universe. As stars formed, they ionized the regions immediately around them, creating bubbles here and there. Eventually these bubbles merged together, and intergalactic gas became entirely ionized.



Time:
Width of frame:
Observed wavelength:

210 million years
2.4 million light-years
4.1 meters

All the gas is neutral. The white areas are the densest and will give rise to the first stars and quasars.



290 million years
3.0 million light-years
3.3 meters

Faint red patches show that the stars and quasars have begun to ionize the gas around them.



370 million years
3.6 million light-years
2.8 meters

These bubbles of ionized gas grow.



460 million years
4.1 million light-years
2.4 meters

New stars and quasars form and create their own bubbles.



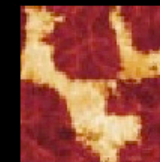
540 million years
4.6 million light-years
2.1 meters

The bubbles are beginning to interconnect.



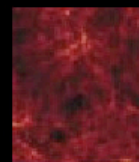
620 million years
5.0 million light-years
2.0 meters

The bubbles have merged and nearly taken over all of space.



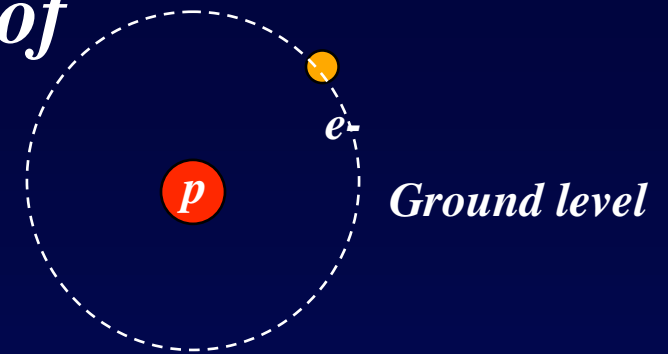
710 million years
5.5 million light-years
1.8 meters

The only remaining neutral hydrogen is concentrated in galaxies.



Simulated images of 21-centimeter radiation show how hydrogen gas turns into a galaxy cluster. The amount of radiation (*white is highest; orange and red are intermediate; black is least*) reflects both the density of the gas and its degree of ionization: dense, electrically neutral gas appears white; dense, ionized gas appears black. The images have been rescaled to remove the effect of cosmic expansion and thus highlight the cluster-forming processes. Because of expansion, the 21-centimeter radiation is actually observed at a longer wavelength; the earlier the image, the longer the wavelength.

Mapping the Cosmic Distribution of Hydrogen

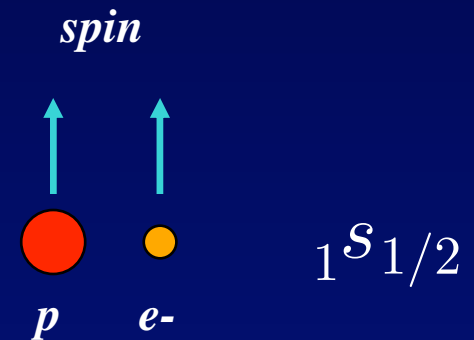
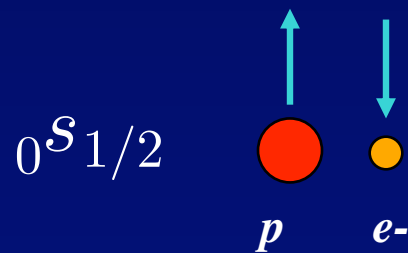


excitation rate = (atomic collisions) + (radiative coupling to CMB)

Couple T_s to T_k

Couples T_s to T_γ

$$21\text{cm} = (1.4\text{GHz})^{-1}$$

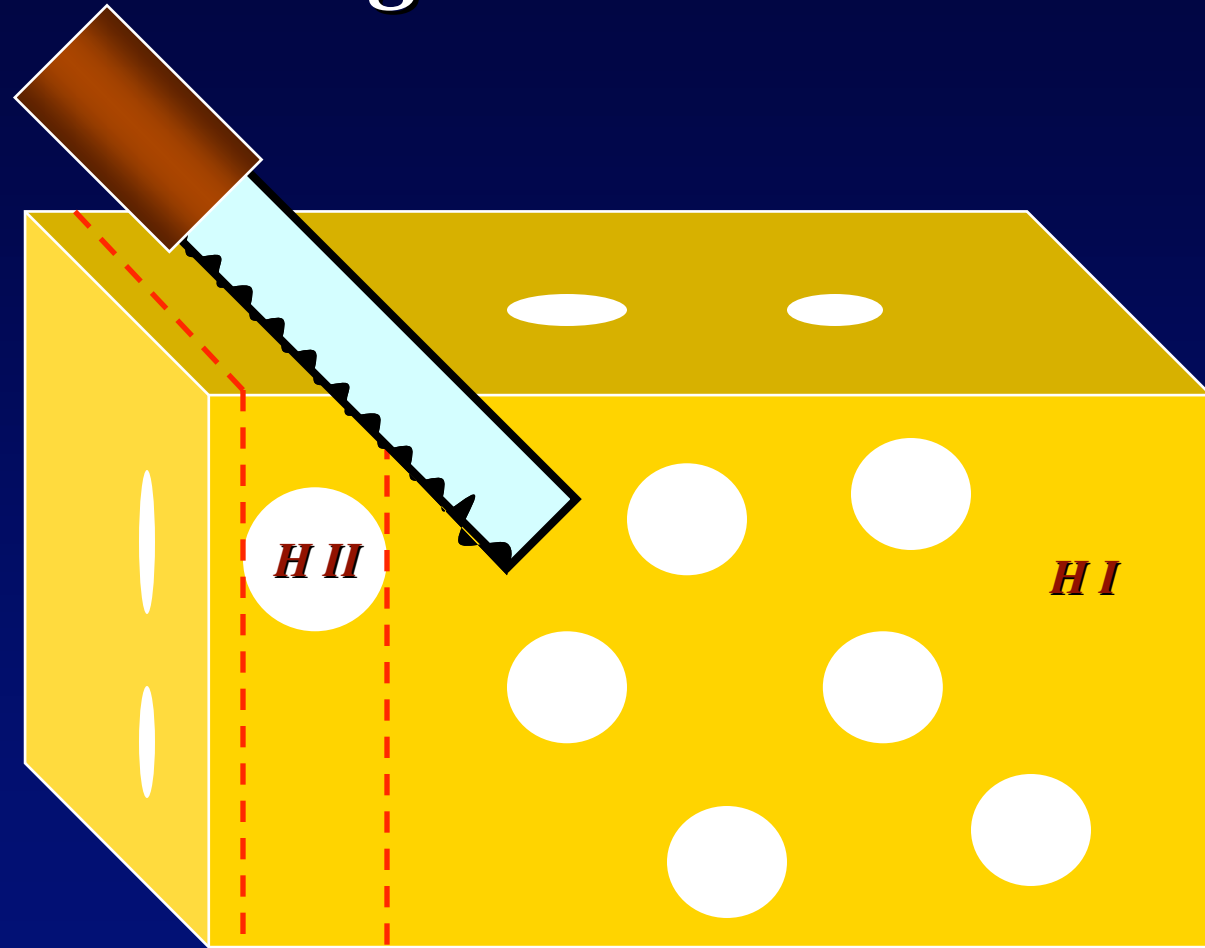


Spin Temperature

$$\frac{n_1}{n_0} = \frac{g_1}{g_0} \exp\left\{-\frac{0.068\text{K}}{T_s}\right\} \quad (g_1/g_0) = 3$$

Predicted by Van de Hulst in 1944; Observed by Ewen & Purcell in 1951 at Harvard

*21cm Tomography of Ionized Bubbles During Reionization is like
Slicing Swiss Cheese*



Observed wavelength \Leftrightarrow distance

$$21\text{cm} \times (1 + z)$$

Separating the Physics from the Astrophysics

Physics: initial conditions from inflation;
nature of dark matter and dark energy

Astrophysics: consequences of star formation

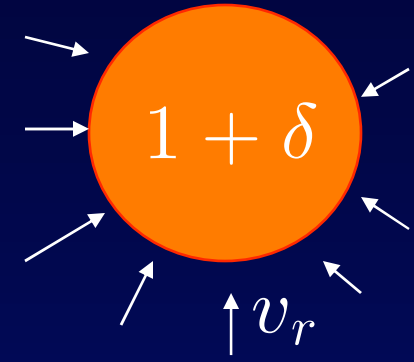
Three epochs:

- Before the first galaxies ($z > 25$): mapping of density fluctuations through 21cm absorption
- During reionization: anisotropy of the 21cm power spectrum due to peculiar velocities
- After reionization ($z < 6$): dense pockets of residual hydrogen (DLAs) trace large scale structure

Line-of-Sight Anisotropy of 21cm Flux Fluctuations

$$T_b = \tau \left(\frac{T_s - T_\gamma}{1+z} \right)$$

Peculiar velocity changes $\tau \propto \frac{n_{\text{HI}}}{dv_r/dr} = \bar{n}(1+\delta) \sim \bar{H}(1 - \frac{1}{3}\delta)$



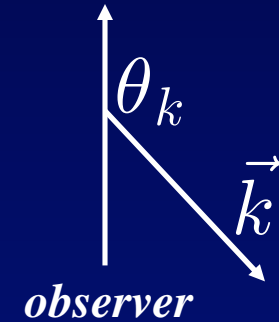
→ *Power spectrum is not isotropic* (“Kaiser effect”)

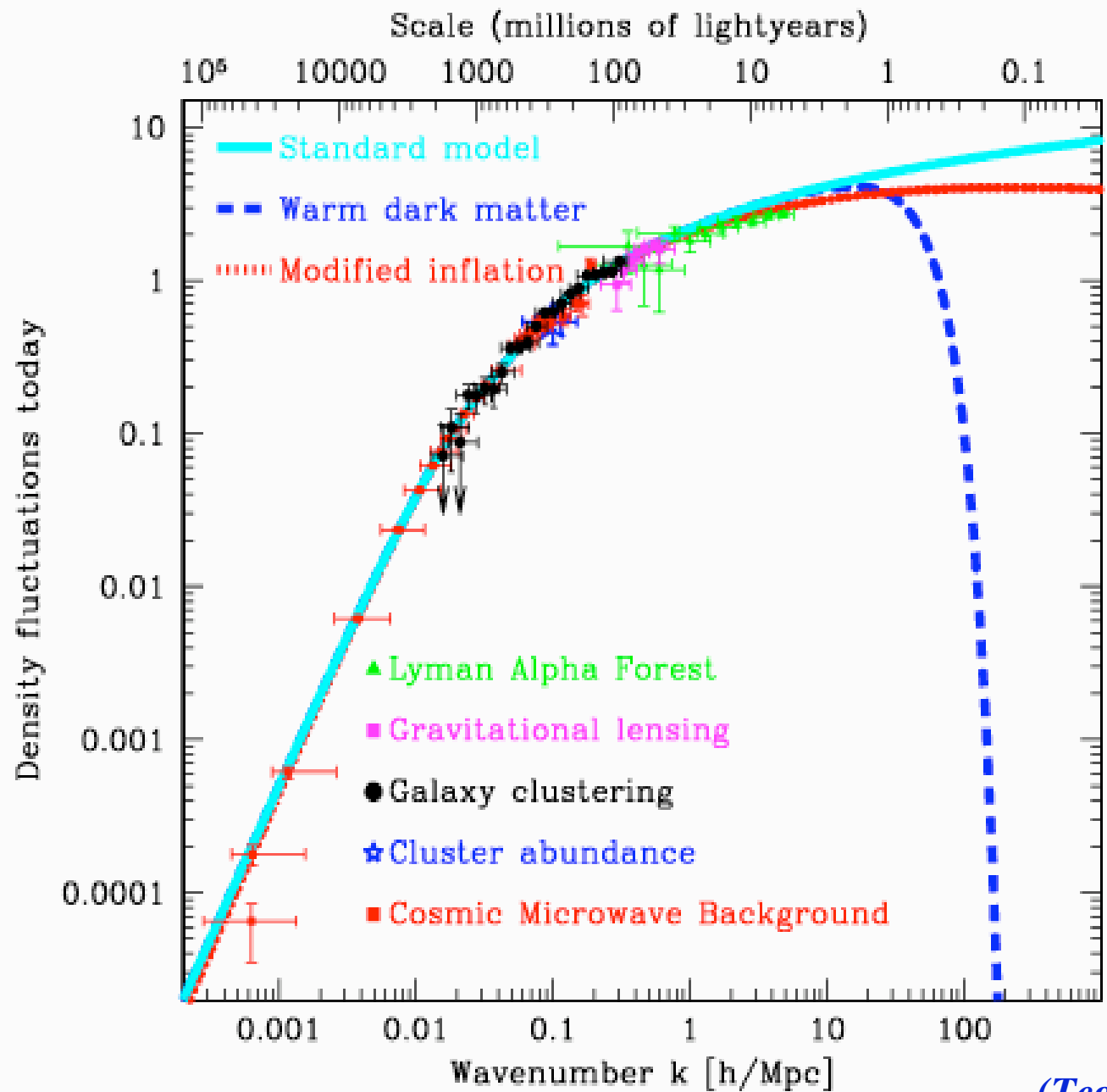
$$\frac{dv_r}{dr} \rightarrow \delta_v(\vec{k}) = -\cos^2 \theta_k \times \delta(\vec{k})$$

$$P_{T_b} = [\cos^2 \theta_k \delta(\vec{k}) + \delta_{\text{iso}}(\vec{k})]^2$$

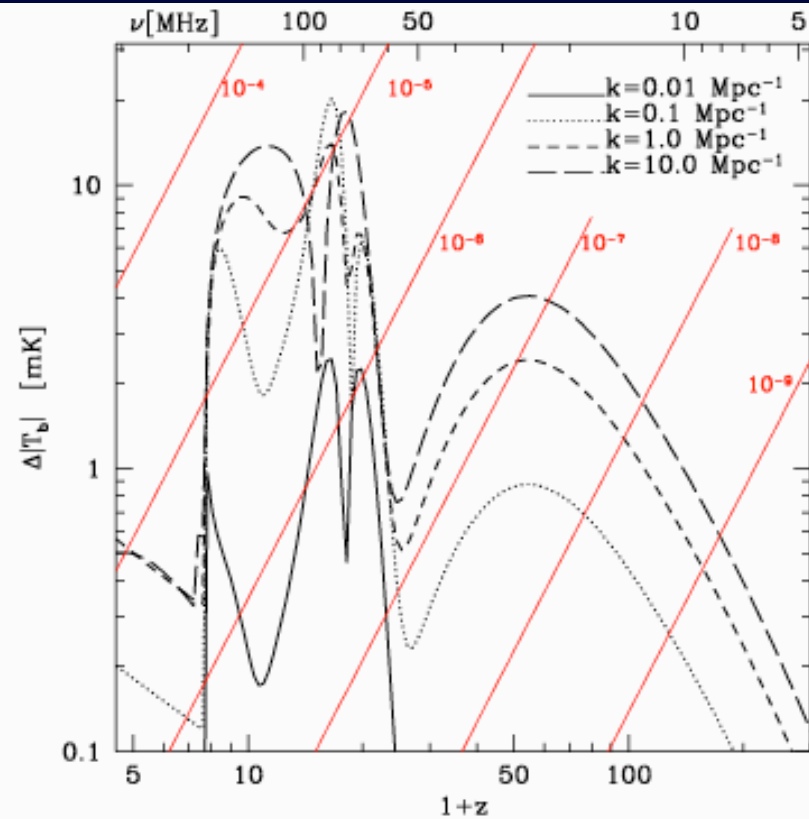
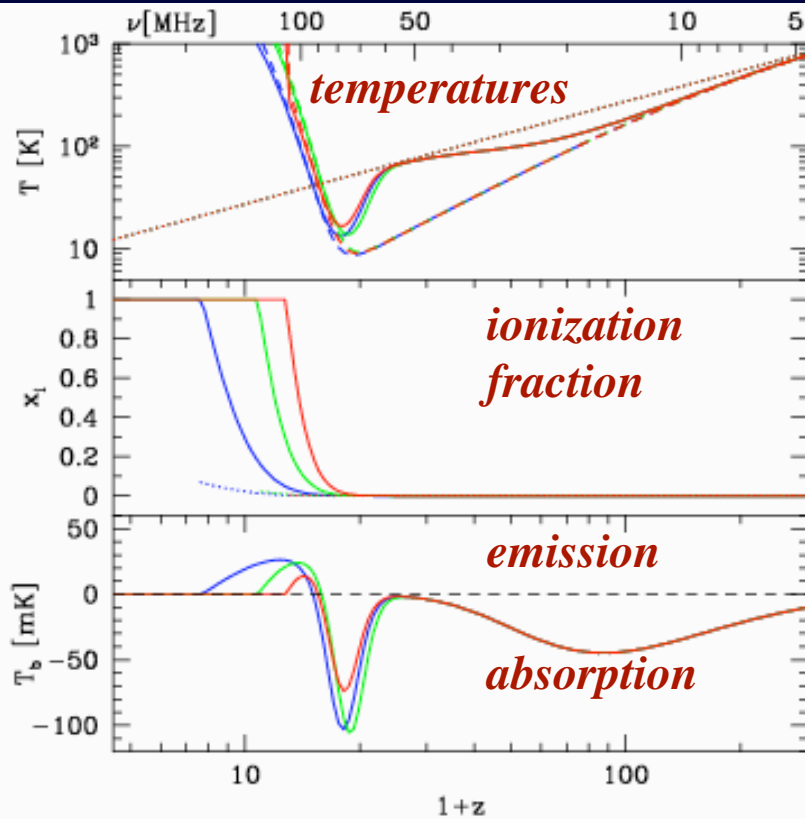
$$\delta_{\text{iso}} = \beta\delta + \delta_{x_{\text{HI}}} + \delta_T + \dots$$

$\cos^4 \theta_k, \cos^2 \theta_k, \cos^0 \theta_k$ *terms allow separation of powers*





(Tegmark 2008)



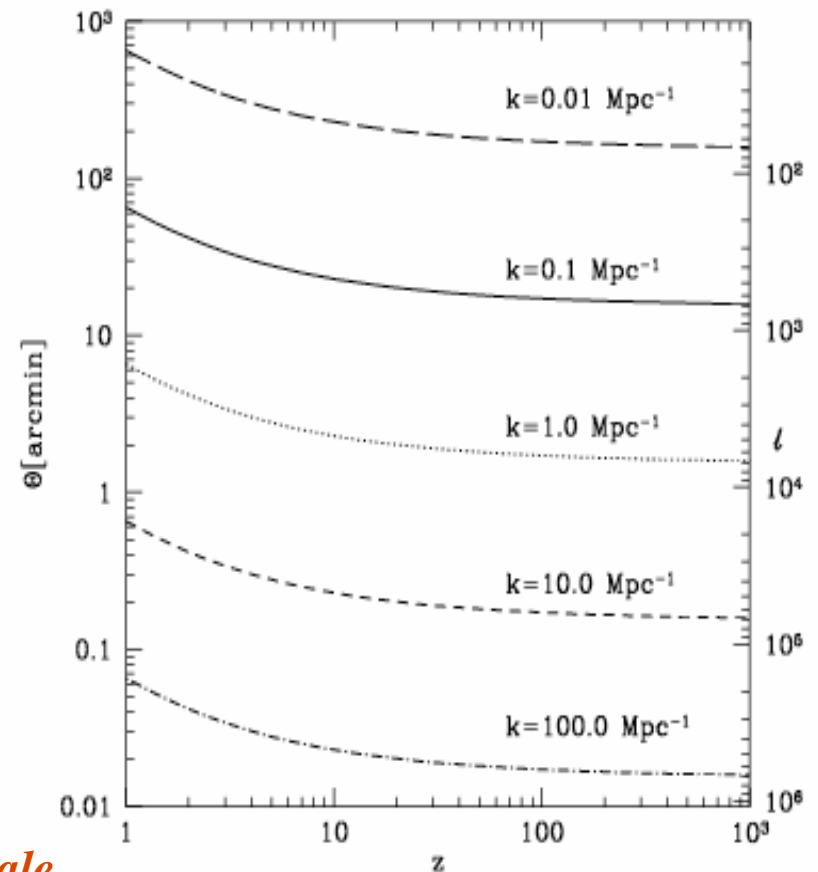
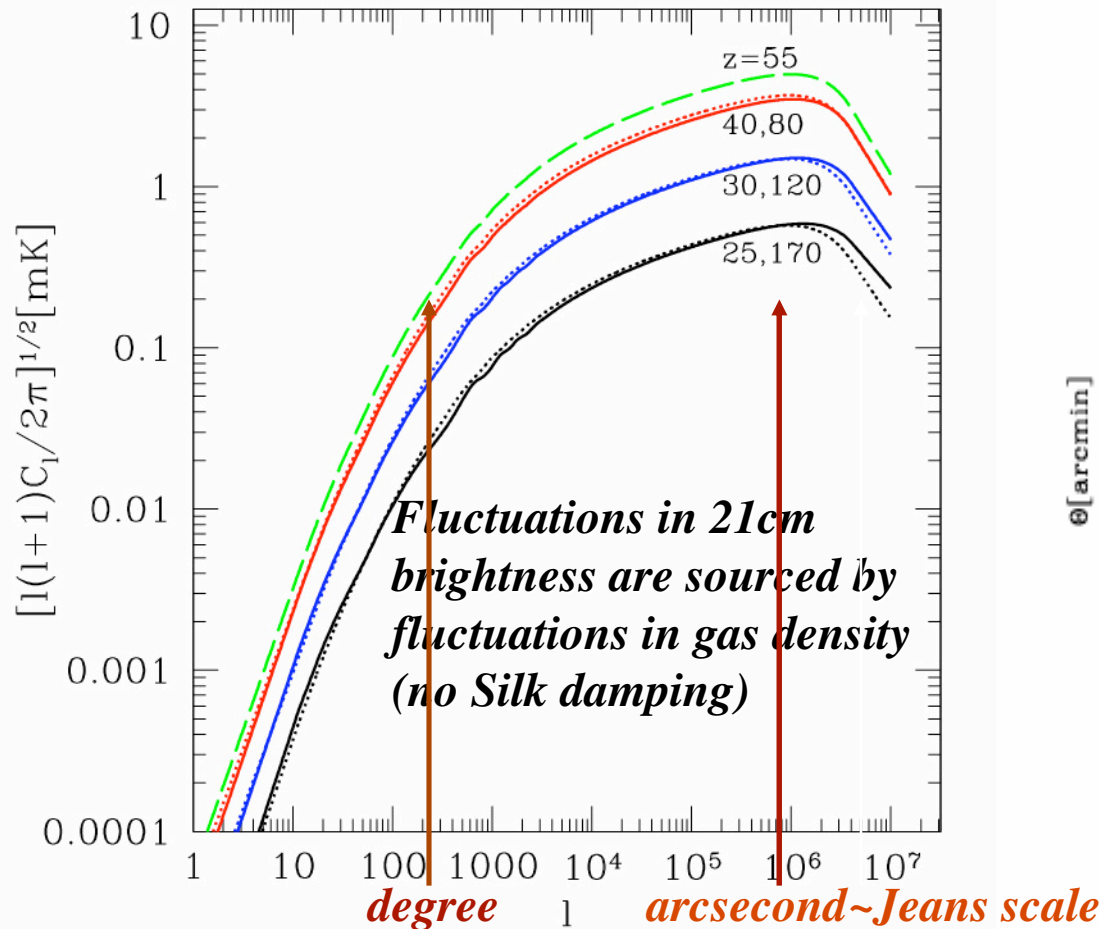
Left: Top panel: Evolution of the mean CMB (dotted curve), intergalactic medium (IGM, dashed curve), and spin (solid curve) temperatures. Middle panel: Evolution of the filling fraction of ionized bubbles (solid curve) and electron fraction outside the bubbles (dotted curve). Bottom panel: Evolution of mean 21cm brightness temperature. Three different astrophysical models are plotted, corresponding to the -1σ (red curve), best-fit (green curve), and $+1\sigma$ (blue curve) optical depth values derived from WMAP [1]. Right: Redshift evolution of the angle-averaged 21cm power spectrum $\bar{\Delta T}_b$ in the -1σ model for wave-numbers $k = 0.01$ (solid curve), 0.1 (dotted curve), 1.0 (short dashed curve), and 10.0 Mpc^{-1} (long dashed curve). Diagonal lines indicate the foreground brightness of the sky $T_{\text{sky}}(\nu)$ times a factor r ranging from 10^{-4} to 10^{-9} , indicative of the level of foreground subtraction required [2].

(Pritchard & Loeb 2008)

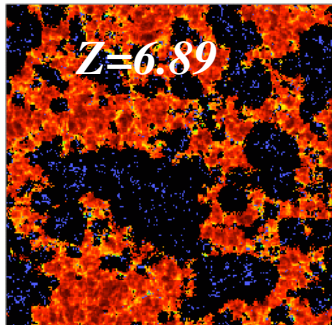
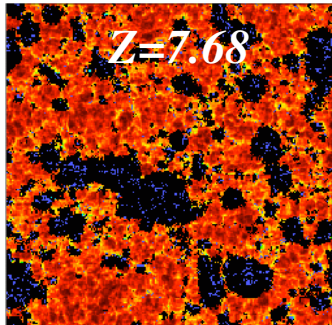
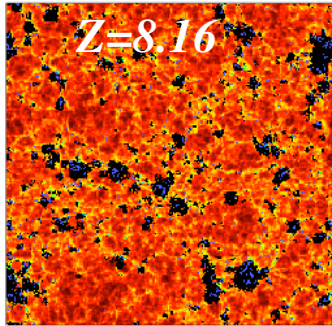
21 cm Absorption During the Dark Ages

$$T_b = \tau \left(\frac{T_s - T_\gamma}{1+z} \right)$$

$$T_b = 28\text{mK} \left(\frac{1+z}{10} \right)^{1/2} \left(\frac{T_s - T_\gamma}{T_s} \right)$$

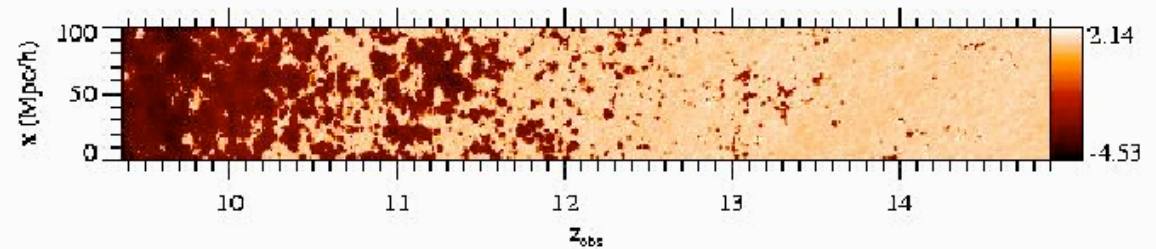
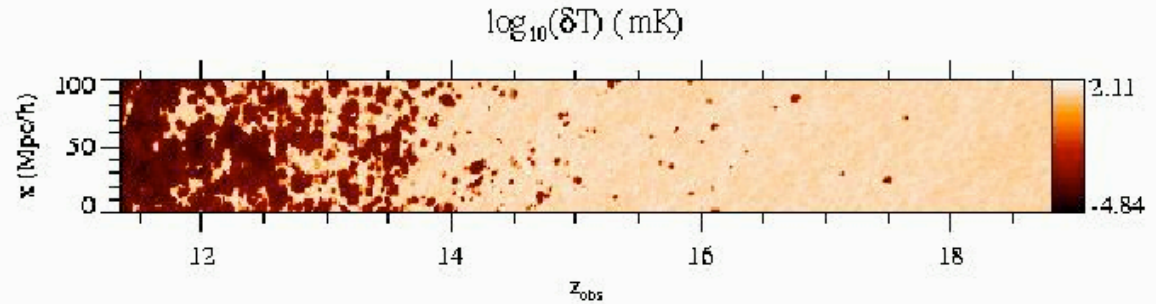


HI Density

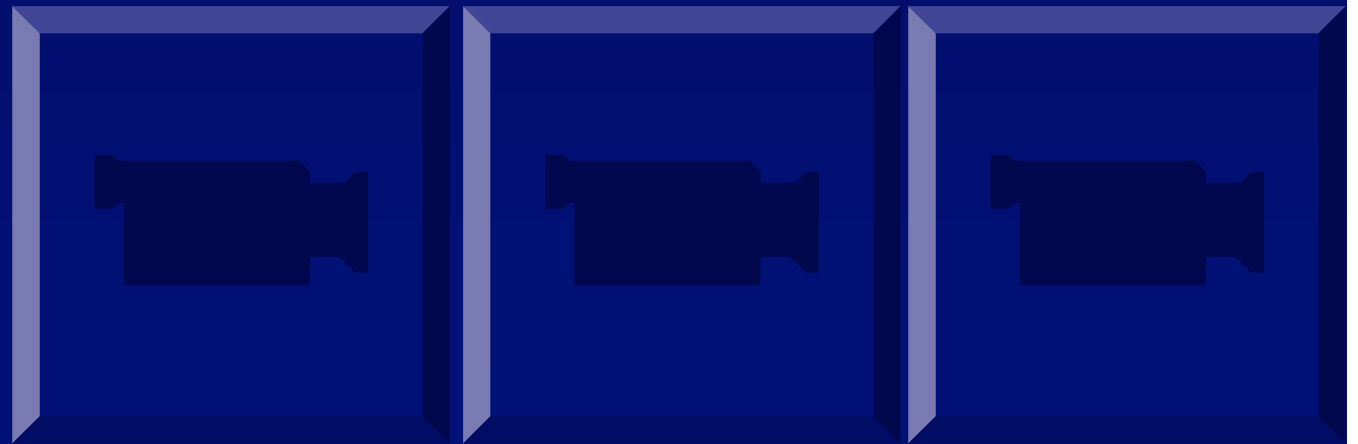


Zahn et al. 2006

21cm Mapping of Epoch of Reionization



Mellema et al. 2006



Trac, Cen, & Loeb 2008

Experiments

**MWA (Murchison Wide-Field Array)*

MIT/U.Melbourne, ATNF, ANU/CfA/Raman I.

**LOFAR (Low-frequency Array)*

Netherlands

**21CMA (formerly known as PAST)*

China

**PAPER*

UCB/NRAO

**GMRT (Giant Meterwave Radio Telescope)*

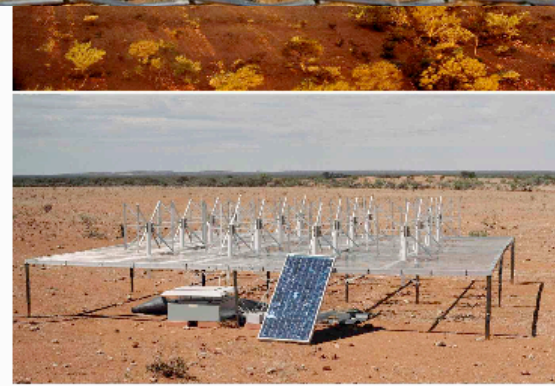
India/CITA/Pittsburg

**SKA (Square Kilometer Array)*

International

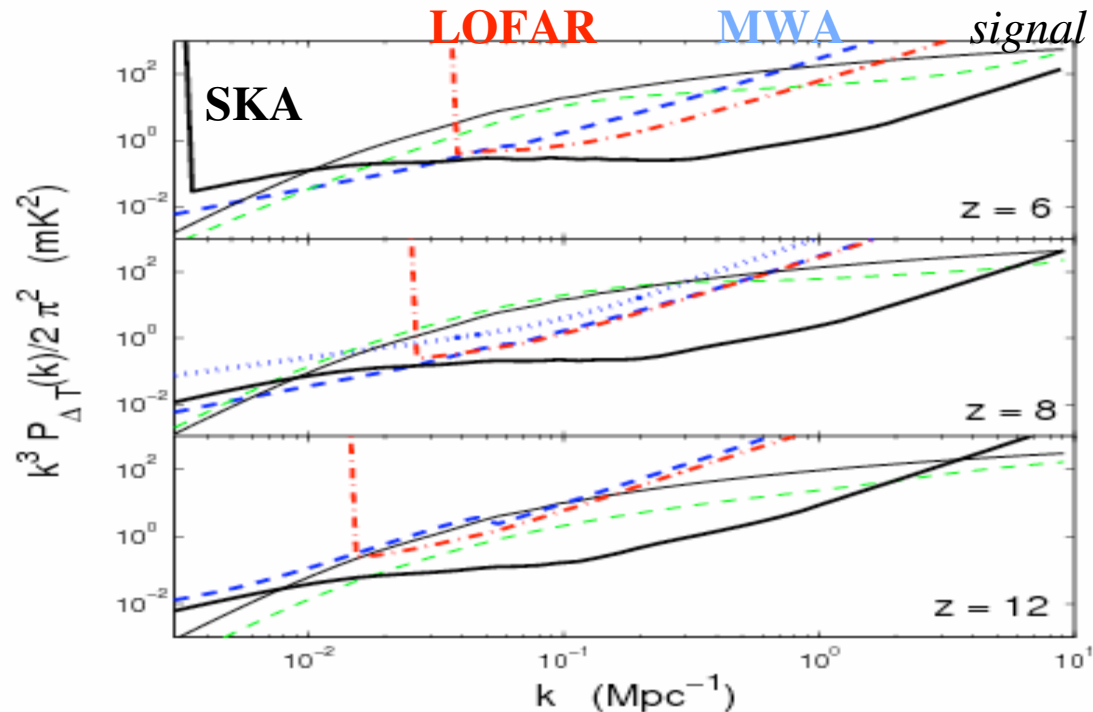


Murchison Wide-Field Array: mapping cosmic hydrogen through its 21cm emission



- 4mx4m tiles of 16 dipole antennae, 80-300MHz
- 500 antenna tiles with total collecting area 8000 sq.m. at 150MHz across a 1.5km area; few arcmin resolution

Power-Spectrum Sensitivity



Isotropic power spectrum sensitivity, in logarithmic bins with $\Delta k = k/2$, for several experimental configurations. In each panel, the thin solid and dashed curves show estimates of the signal with and without reionization. The thick solid, dashed, and dot-dashed curves show error estimates for 1000 hour observations over 6 MHz with the SKA, MWA, and LOFAR, respectively. Each assumes perfect foreground removal. The dotted curve in the middle panel assumes a flat antenna distribution for the MWA. From *McQuinn et al. 2006*

$$T_{\text{sky}} \sim 180 \left(\frac{\nu}{180 \text{ MHz}} \right)^{-2.6} \text{ K}$$

$$\Delta T^N|_{\text{int}} \sim 2 \text{ mK} \left(\frac{A_{\text{tot}}}{10^5 \text{ m}^2} \right)^{-1} \left(\frac{10'}{\Delta\theta} \right)^2 \left(\frac{1+z}{10} \right)^{4.6} \left(\frac{\text{MHz}}{\Delta\nu} \frac{100 \text{ hr}}{t_{\text{int}}} \right)^{1/2}$$

Cross-correlation between 21cm brightness and galaxy density

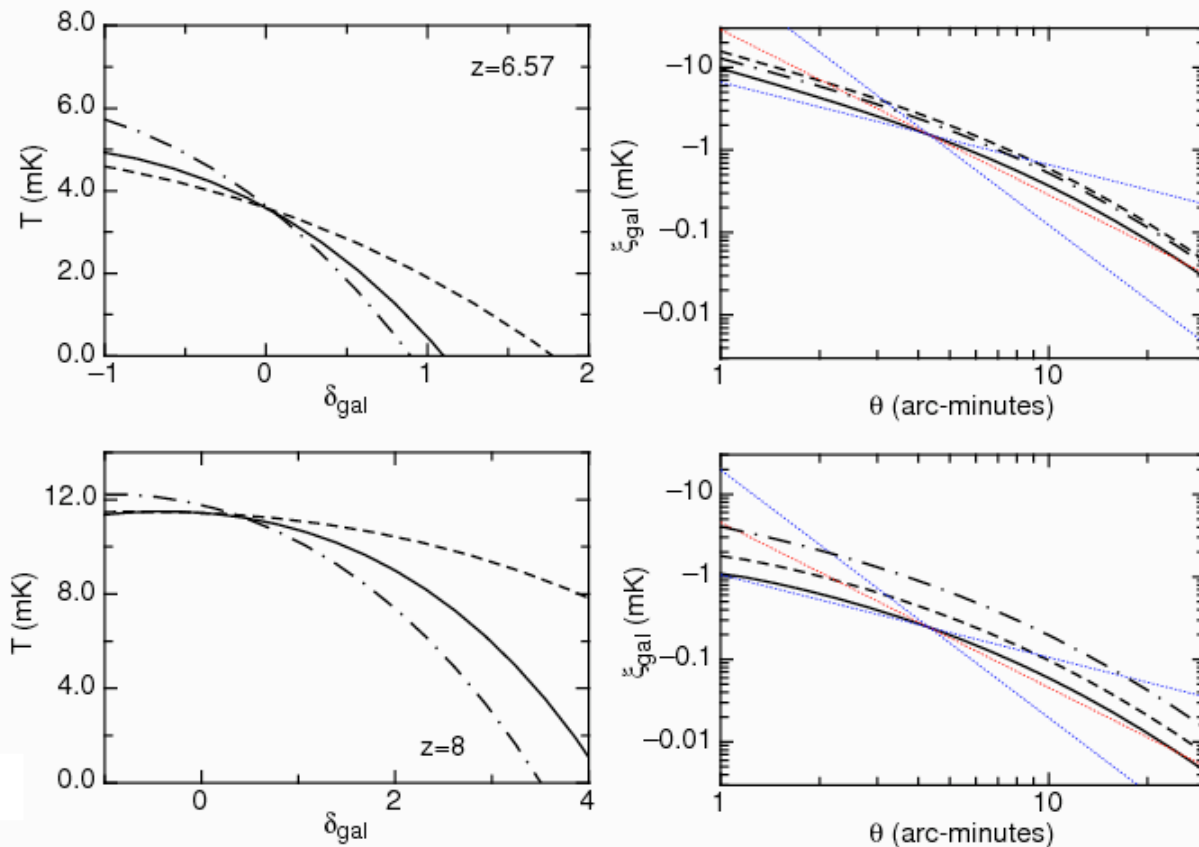


Figure 4. *Left:* 21cm brightness temperature as a function of δ_{gal} . Two values of galaxy mass are assumed for a clumping of $C = 10$, $M = 10^{10} M_{\odot}$ (solid line) and $M = 10^{11} M_{\odot}$ (dashed line). The dot-dashed line shows $C = 2$ with $M = 10^{10} M_{\odot}$. *Right:* The cross-correlation function $\xi_{\text{gal}} = \langle \delta_{\text{gal}}(T - \langle T \rangle) \rangle$ for the IGM smoothed on various angular scales (θ). The function is presented assuming $C = 10$ for masses of $M = 10^{10} M_{\odot}$ (solid line) and $M = 10^{11} M_{\odot}$ (dashed line). The dot-dashed line represents $C = 2$ with $M = 10^{10} M_{\odot}$. The lines show power-laws of slope $d(\log \xi_{\text{gal}})/d(\log \theta) = -1, -2$ and -3 . The upper and lower rows correspond to observations at $z = 6.57$ and $z = 8$ respectively.

Infrared imaging

Galaxy/Quasar

HI hole

21cm mapping

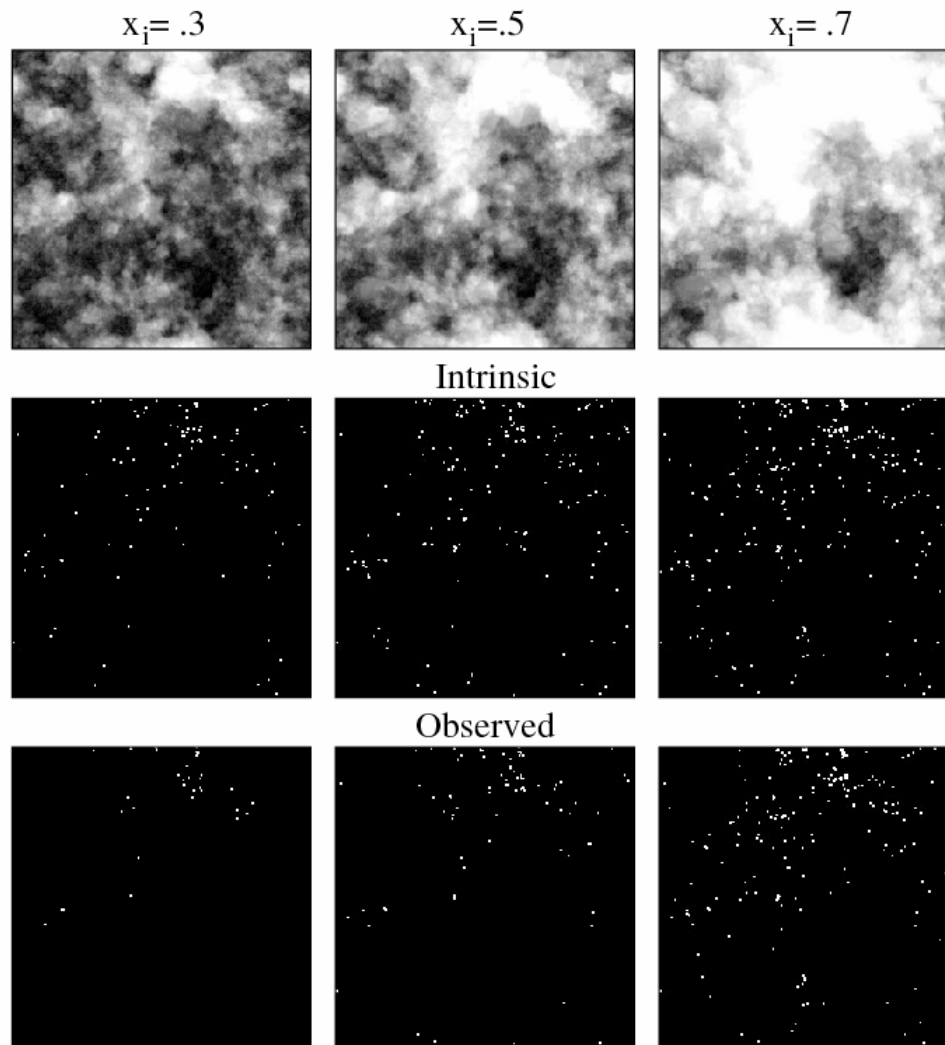


Figure 2. Top panels show the projection of \bar{x}_i in the survey volume. In the white regions the projection is fully ionized and in black it is neutral. The left, middle, and right panels are for $z = 8.2$ ($\bar{x}_i = 0.3$), $z = 7.7$ ($\bar{x}_i = 0.5$), and $z = 7.3$ ($\bar{x}_i = 0.7$). The middle and bottom rows are the intrinsic and observed Ly α emitters maps, respectively, for $f_E = 0.25$ and assuming that we can observe unobscured emitters with $m \exp(-\tau_\alpha(\nu_0)) > 7 \times 10^{10} M_\odot$. (Note that $L_{\text{int,E}} \propto m$.) The observed distribution of emitters is modulated by the location of the HII regions (compare bottom panels with corresponding top panels). Each panel is 94 Mpc across (or 0.6 degrees on the sky), roughly the area of the current Subaru Deep Field (SDF) at $z = 6.6$ (Kashikawa et al. 2006). The depth of each panel is $\Delta\lambda = 130 \text{ \AA}$, which matches the FWHM of the Subaru 9210 \AA narrow band filter. The number densities of Ly α emitters for the panels in the middle row are few times larger than the number density in the SDF photometric sample of $z = 6.6$ LAEs.

Clustering of Ly α Emitters

McQuinn et al.
arXiv:0704.2239

The Imprint of Reionization on Galaxy Clustering

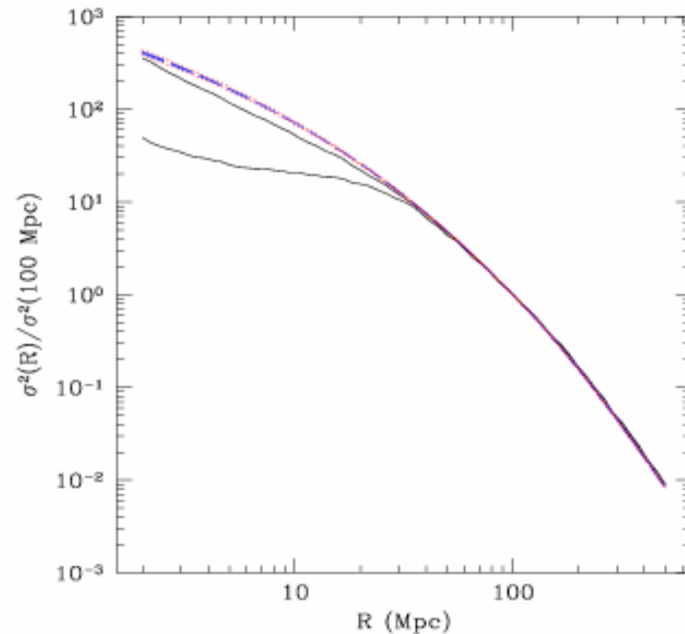


FIG. 5.— The normalized $\sigma^2(R)$ for $z_R = 15$ and $z_R = 6$ (upper and lower solid black curves, respectively); the limiting cases of the best current constraints of n (red, dotted curves) and α (blue, dashed curves).

Inhomogeneous photo-ionization heating to $\sim 10^4\text{K}$ modulates the minimum mass of galaxies on scales of tens of comoving Mpc

Babich & Loeb 2006, ApJ, 640, 1

Highlights

- *Large-aperture infrared telescopes and radio arrays will probe reionization over the coming decade.*
- *Simulations of reionization require large ($>100\text{Mpc}$) boxes and high resolution for source identification.*
- *21cm brightness fluctuations are expected to be anti-correlated with infrared galaxies during reionization*
- *Reionization leaves an imprint on the clustering of star-forming galaxies at intermediate redshifts. This might compromise the use of these galaxies for precision cosmology (e.g. acoustic oscillations/dark energy).*