

Merger Influence on the Thermal SZ Effect

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1. Model and characteristics:
 - integrated pressure
 - spectral function
 - (transonic merger)
2. An application: maps for subsonic mergers

Why look at mergers?

- increasing observational evidence from X-rays:
dimension: $r \sim 10\text{-}600$ kpc
Mach number: 0.4 – 2.5
(early stage/late stage, subsonic/transonic)
- source of non-thermal physics (spectral fct.)
- reveal the cluster substructure
- information on structure formation

Why look at mergers with SZ?

- compare with X-rays, provide an (independent) tool (in particular for high z)
- observe mergers which are not seen in X-rays (position and direction of motion), possibly extract parameters from a simple model
- investigate (high z) substructure
- source of error for parameters based on SZ flux, e.g. H_0

MODEL: framework

- focus on (common) minor mergers
 - ◊ DM profile not changed
(major mergers excluded)
 - $r_{\text{sub}} \leq 0.6 r_l$, $M_{\text{sub}} \leq 10\text{-}20 \% M_{\text{cl}}$
 - simple gas dynamics: gas spread around subcluster,
redistribution maintaining β -profile
- ◊ goal: build a model with characteristic merger features

Model and characteristics: subsonic

$$\Delta I(x) \propto g(x) \int_{cl} \left(\frac{k_B T_g}{m_e c^2} \right) \sigma_T n_e dl_{cl}$$

Subsonic:

- \sim incompressible, but P_{ex}
- Moving body through gas
- No shock waves $\diamond g(x)$ unchanged

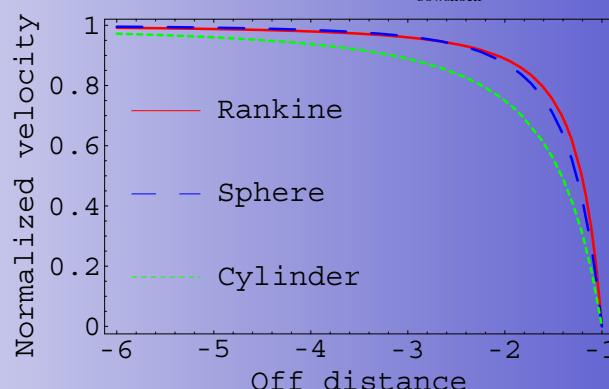
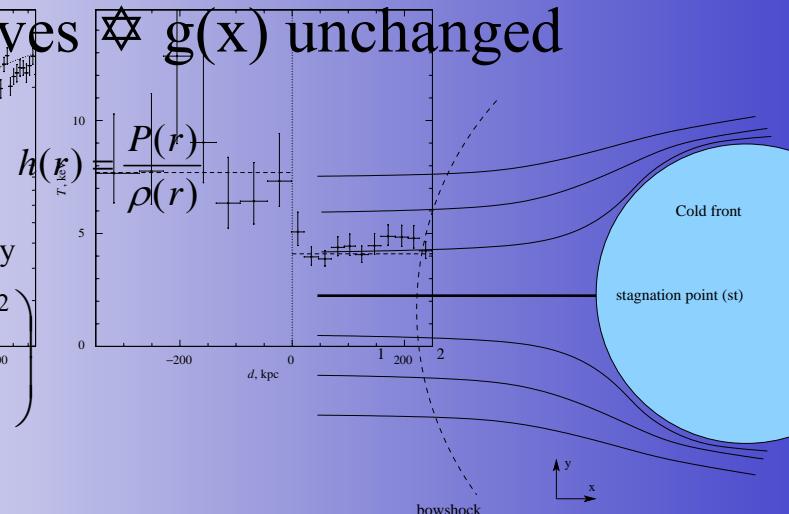
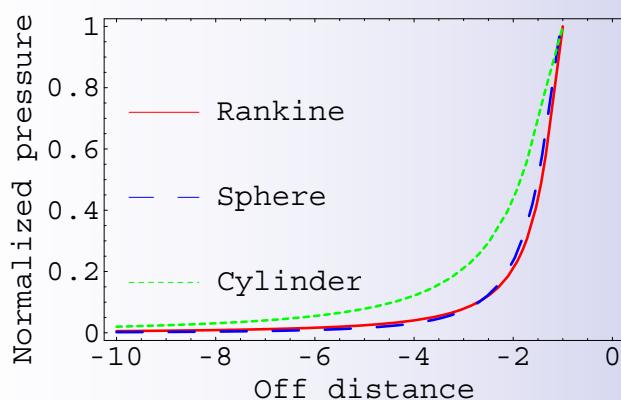
Bernoulli eq.

$$h(r) + \frac{1}{2} v_{\infty}^2(r) + \phi(r) = \text{const.}$$

Potential flow theory

$$(Stagnation point l.o.s) P_{ex}(r) \propto \rho_{gas}^{-3}(r) \frac{1}{2} \left(1 - \frac{1}{a^3} \left(1 + \frac{a^3}{z^3} \right)^2 \right)$$

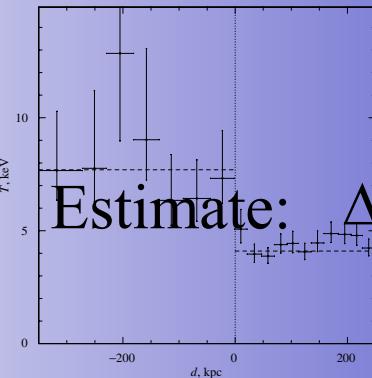
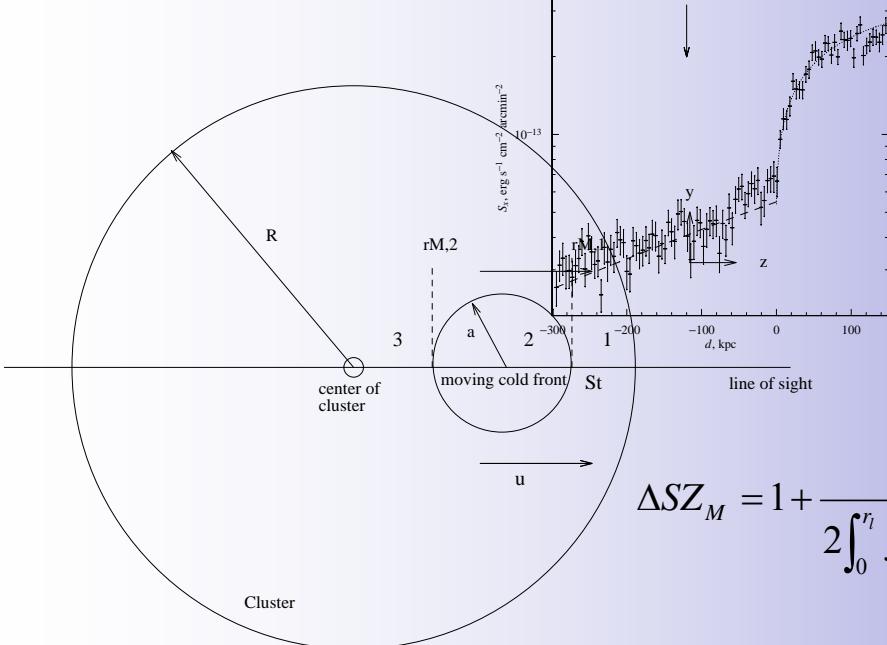
(Stagnation point l.o.s)



Modified SZ signal (1)

$$\Delta I(x) \propto g(x) \int_{cl} \left(\frac{k_B T_g}{m_e c^2} \right) \sigma_T n_e dl_{cl}$$

- no change in $g(x)$
- 3 regions along l.o.s: P_{ex}, subcluster, turbulent region behind



Estimate: $\Delta SZ_M = \frac{\Delta I_M(x)}{\Delta I(x)}$

$$\Delta SZ_M = 1 + \frac{1}{2 \int_{r_l}^{r_{M,1}} f(r) dr} \left[\int_{r_l}^{r_{M,1}} f(r) \frac{1}{2} M^2 \left(1 - \frac{v(z)^2}{u_\infty^2} \right) dr - \int_{r_{M,1}}^{r_{M,2}} f(r) dr + a \frac{n_{e,M} T_M}{n_{e,0} T} \right]$$

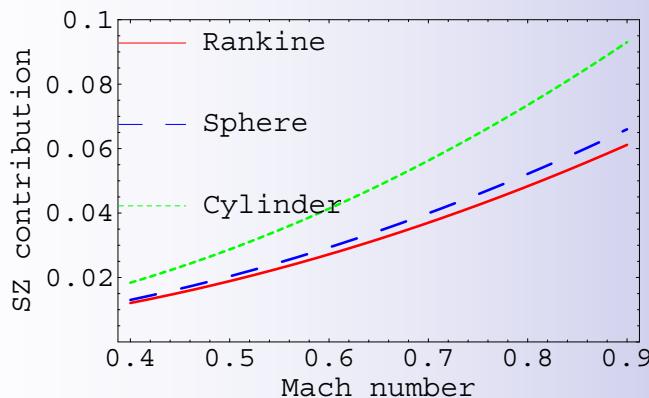
◊ Many parameters: r_c , M , subcluster

Modified SZ signal (2)

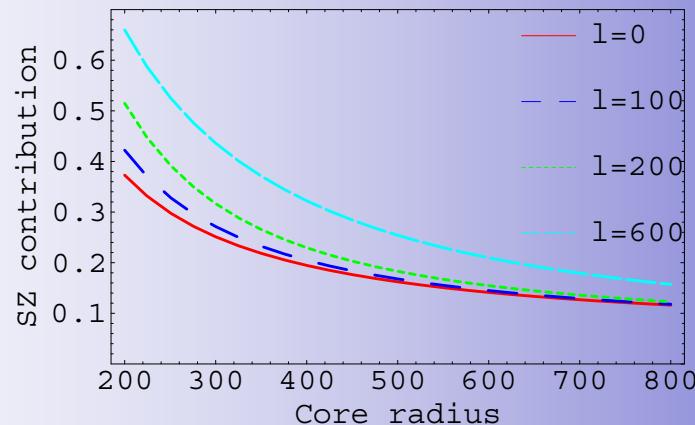
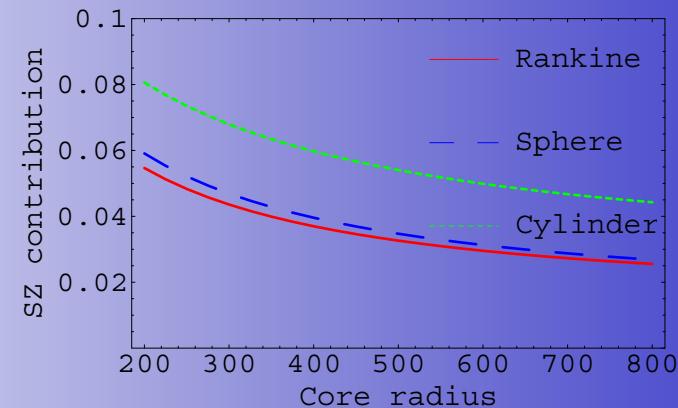
$$\Delta SZ_M = \frac{\Delta I_M(x)}{\Delta I(x)}$$

$$\Delta SZ_M = 1 + \frac{1}{2 \int_0^r f(r) dr} \left[\int_{r_l}^{r_{M,1}} f(r) \frac{1}{2} M^2 \left(1 - \frac{v(z)^2}{u_\infty^2} \right) dr - \int_{r_{M,1}}^{r_{M,2}} f(r) dr + a \frac{n_{e,M} T_M}{n_{e,0} T} \right]$$

Excess pressure contribution

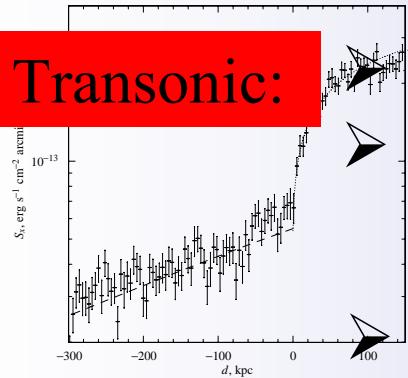


Excess pressure contribution



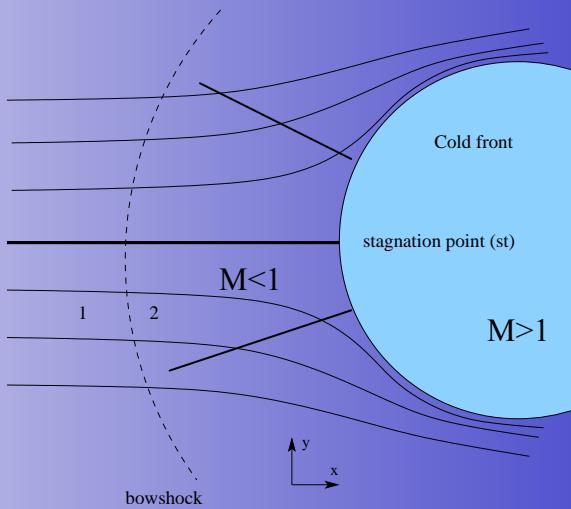
subcluster contribution
depending on location l

Model and characteristics: transonic



➤ bow shock ⚫ change microphysics, $g(x)$
 ➤ Rankine-Hugoniot jump conditions at shock front
 ⚫ test with observations: C
 compressible flow within shock region

$$\left. \begin{array}{l} \rho_1 v_1 = \rho_2 v_2 \\ P_1 + \rho_1 v_1^2 = P_2 + \rho_2 v_2^2 \\ h_1 + \frac{1}{2} v_1^2 = h_2 + \frac{1}{2} v_2^2 \end{array} \right\} \quad \begin{aligned} \frac{\rho_1}{\rho_2} &= \frac{2\gamma}{\gamma+1} M^2 + \frac{\gamma-1}{\gamma+1} = \frac{1}{C} \\ \frac{P_2}{P_1} &= \frac{2\gamma}{\gamma+1} M^2 - \frac{\gamma-1}{\gamma+1} \end{aligned}$$



❖ Result: additional contribution due to pressure jump

Modified SZ signal (1)

What changes? \diamond

$g(x)$

- shock particle re-acceleration mechanism depending on C
(1st order Fermi acc., stochastic acc. ... Not well known!)
- new emerging electron population
- hybrid electron distribution with high energy power law tail



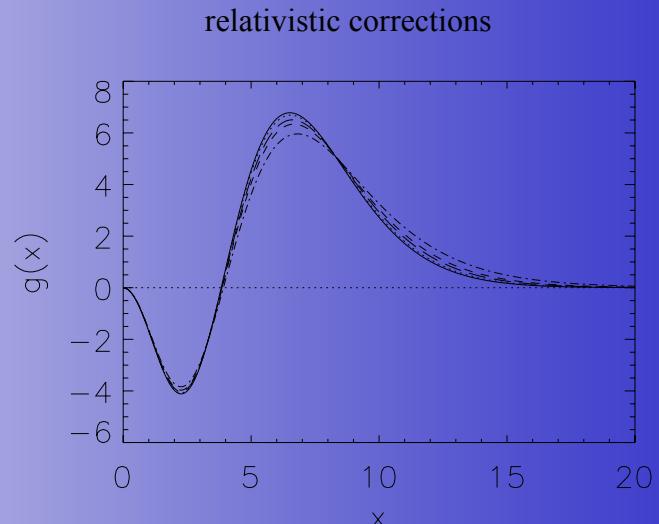
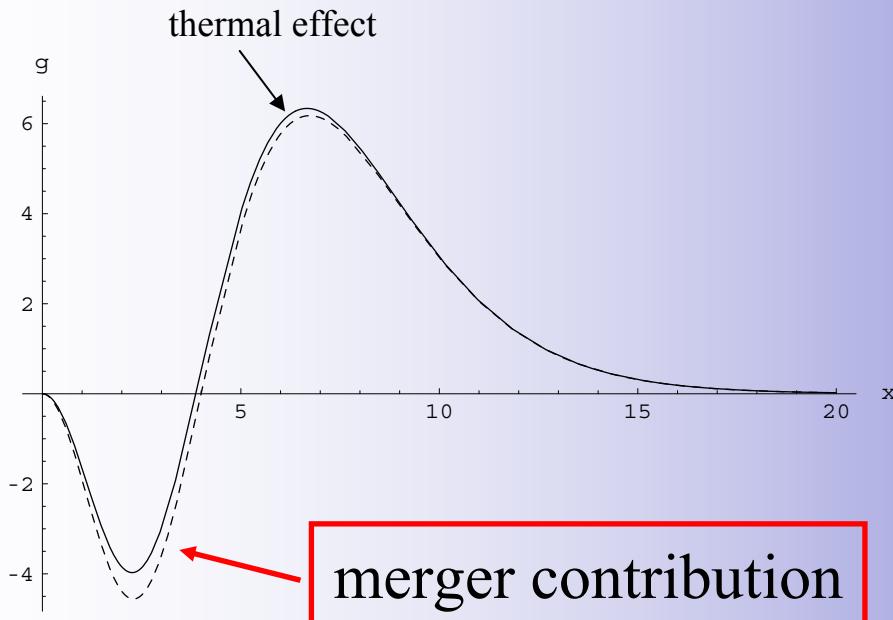
$$\alpha = \frac{C+2}{C-1}$$

\diamond Problem: lower limit p_l for power law!

combination of 2 (independent) electron populations:
 \diamond cross scattering, requires new calculation technique

Modified SZ signal (2)

$$g(x)$$

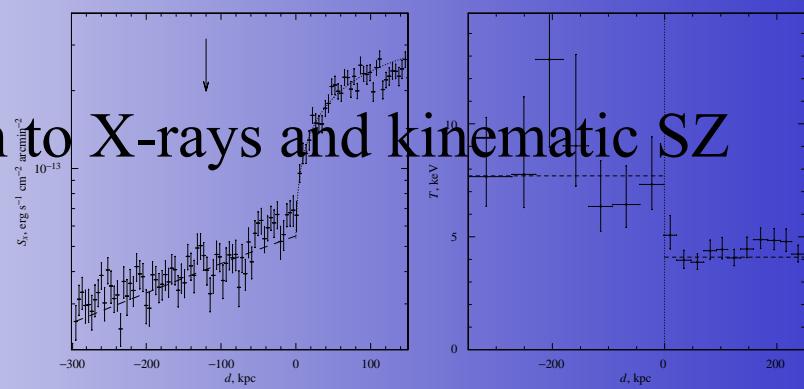
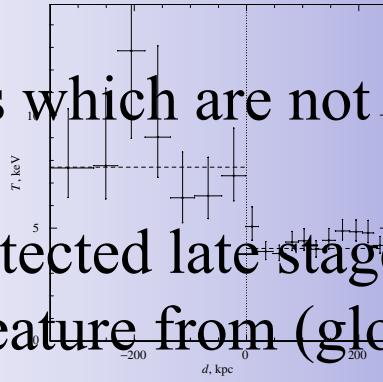
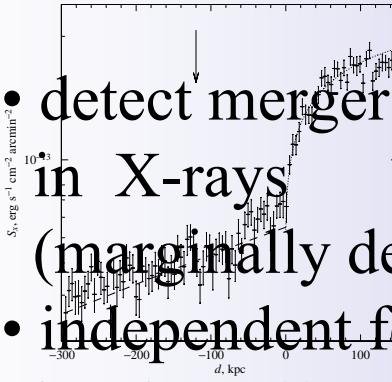


- ❖ Results:
- change in crossover $\sim 10\%$
 - limited to certain regions?

MAPS FOR SUBSONIC MERGERS

Interest:

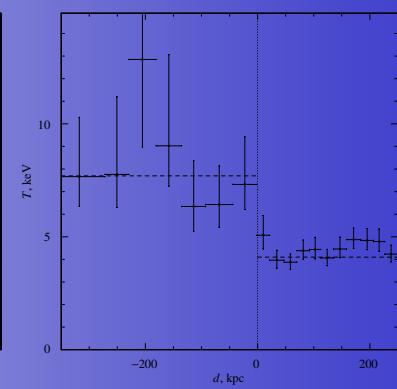
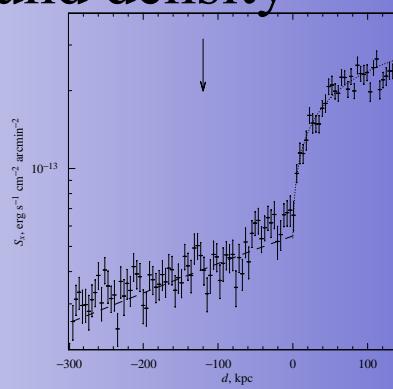
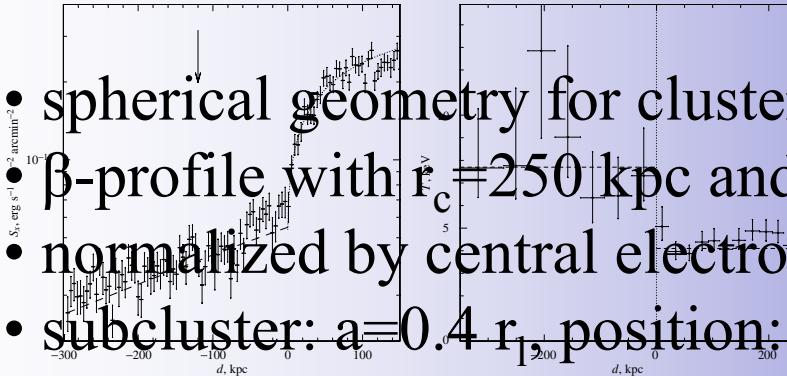
- detect mergers which are not identified by bow shocks in X-rays (marginally detected late stage, cold fronts)
- independent feature from (global) compressibility and heating effects
- look for merger dynamics
- complementary information to X-rays and kinematic SZ



Maps for subsonic mergers

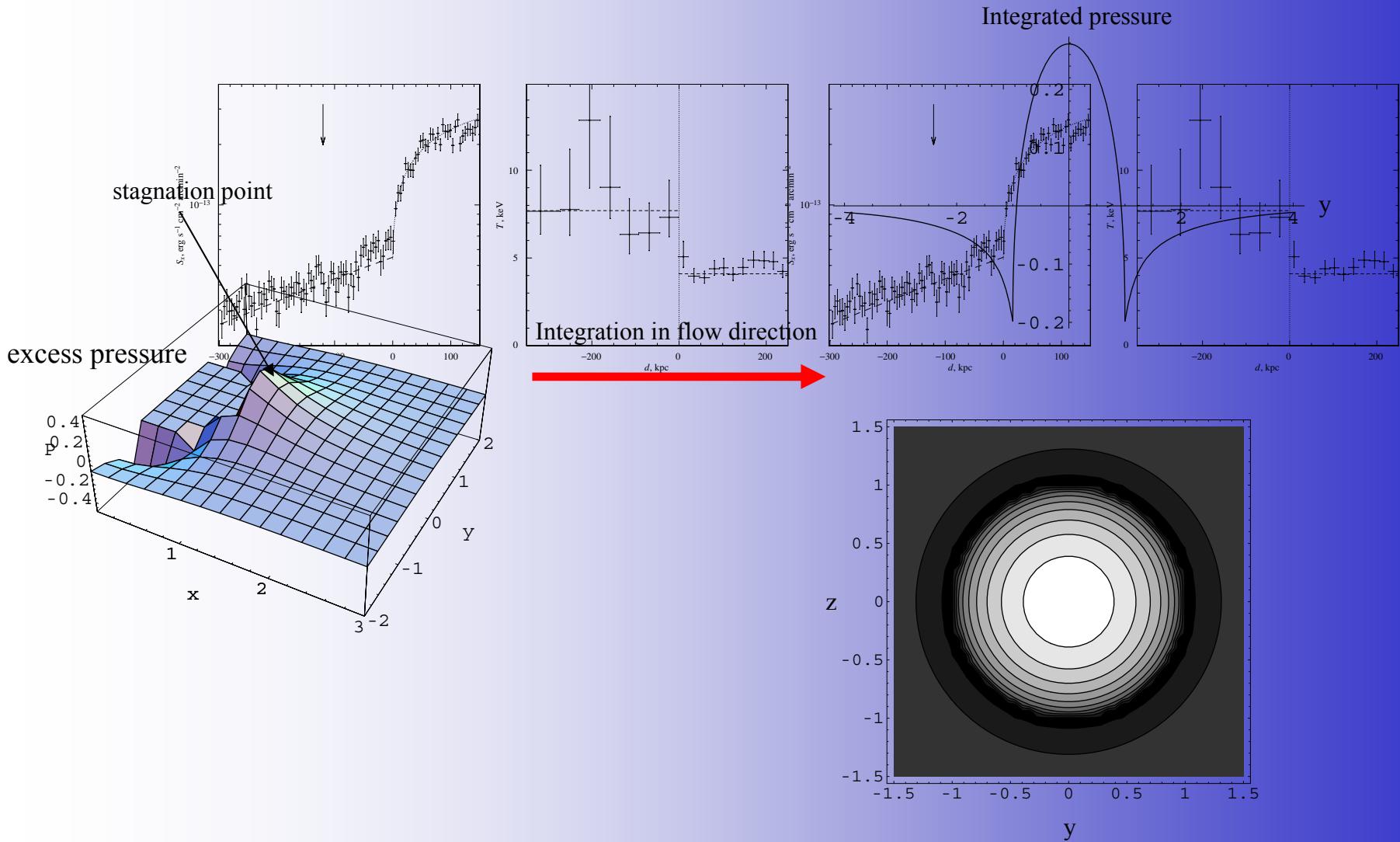
Parameters:

- spherical geometry for cluster and subcluster
- β -profile with $r_c = 250$ kpc and $\beta = 2/3$, $r_l = 1$ Mpc
- normalized by central electron density and temperature
- subcluster: $a = 0.4 r_l$, position: $x = 0.5 r_l$, $y = z = 0$, $M = 0.6$
- subcluster density = background density
- DM structure not changed



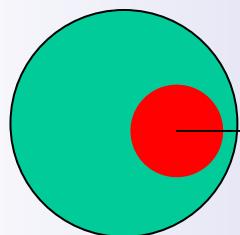
Maps for subsonic mergers

excess pressure feature:

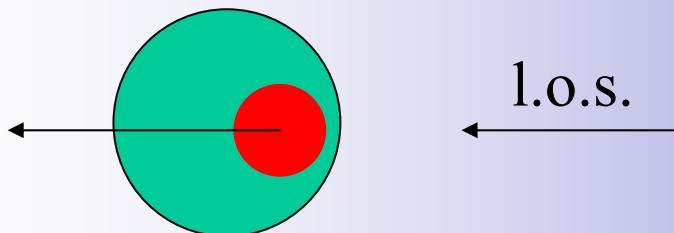


Maps for subsonic mergers

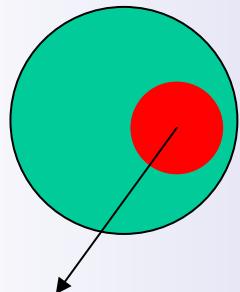
location ⋆ direction of movement



late stage merger, cold front
(slowed down)



infalling substructure

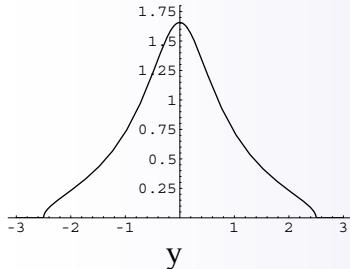


transversally moving
substructure

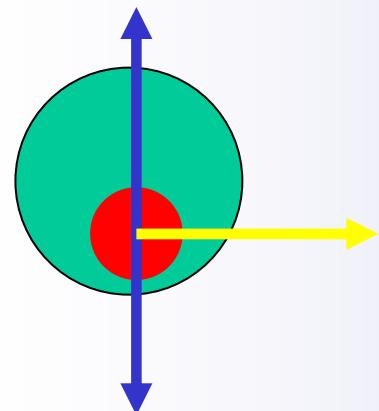
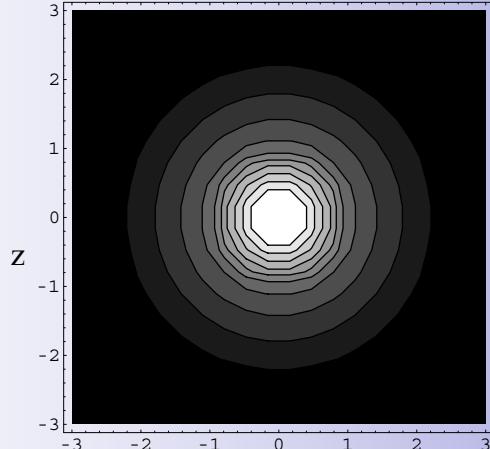
Maps for subsonic mergers

results: excess pressure contribution

integrated pressure

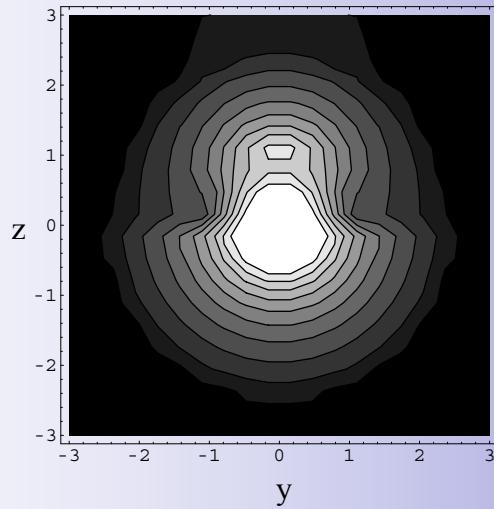


parallel

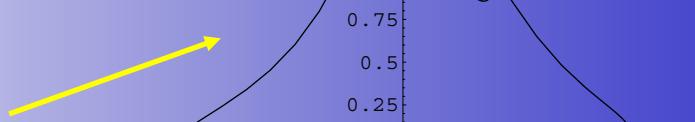
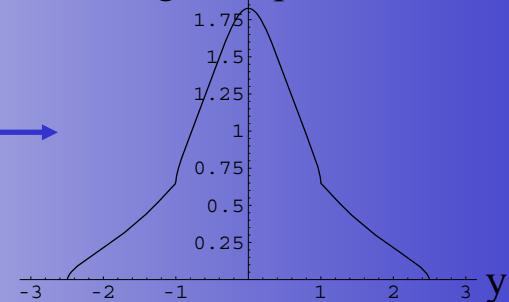


l.o.s.

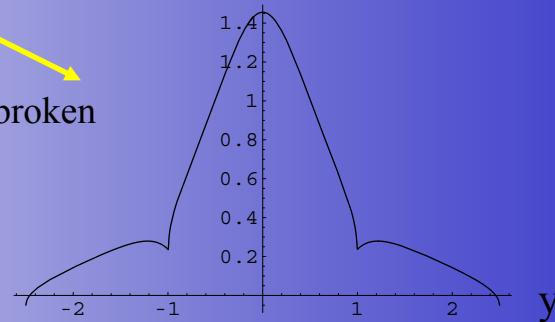
perpendicular



integrated pressure



Symmetry broken



CONCLUSION AND SUMMARY

- model limited to minor or medium-scale merger
 - non-negligible contribution in intensity: $\sim 10\text{-}15\%$
 - transonic: change in spectral function $g(x)$ at crossover frequency $\sim 5\text{-}10\%$
 - (Problem: disentangle from relativistic corrections)
 - complementary/additional information due to excess pressure compared to X-rays (“invisible merger”)
 - different features depending on direction of moving substructure (subsonic mergers), change: $\sim 5\text{-}10\%$
 - complementary to kinematic SZ
 - detect substructures with high resolution SZ mapping (high z)
 - SZ-survey: constrain merger rate and structure formation
- 