

4.5 Post-Newtonian order gravitational radiation

Tanguy Marchand - IAP
marchand@iap.fr

*T. Marchand, L. Blanchet, G. Faye
arxiv:1607.07601 (submitted to QCG)*

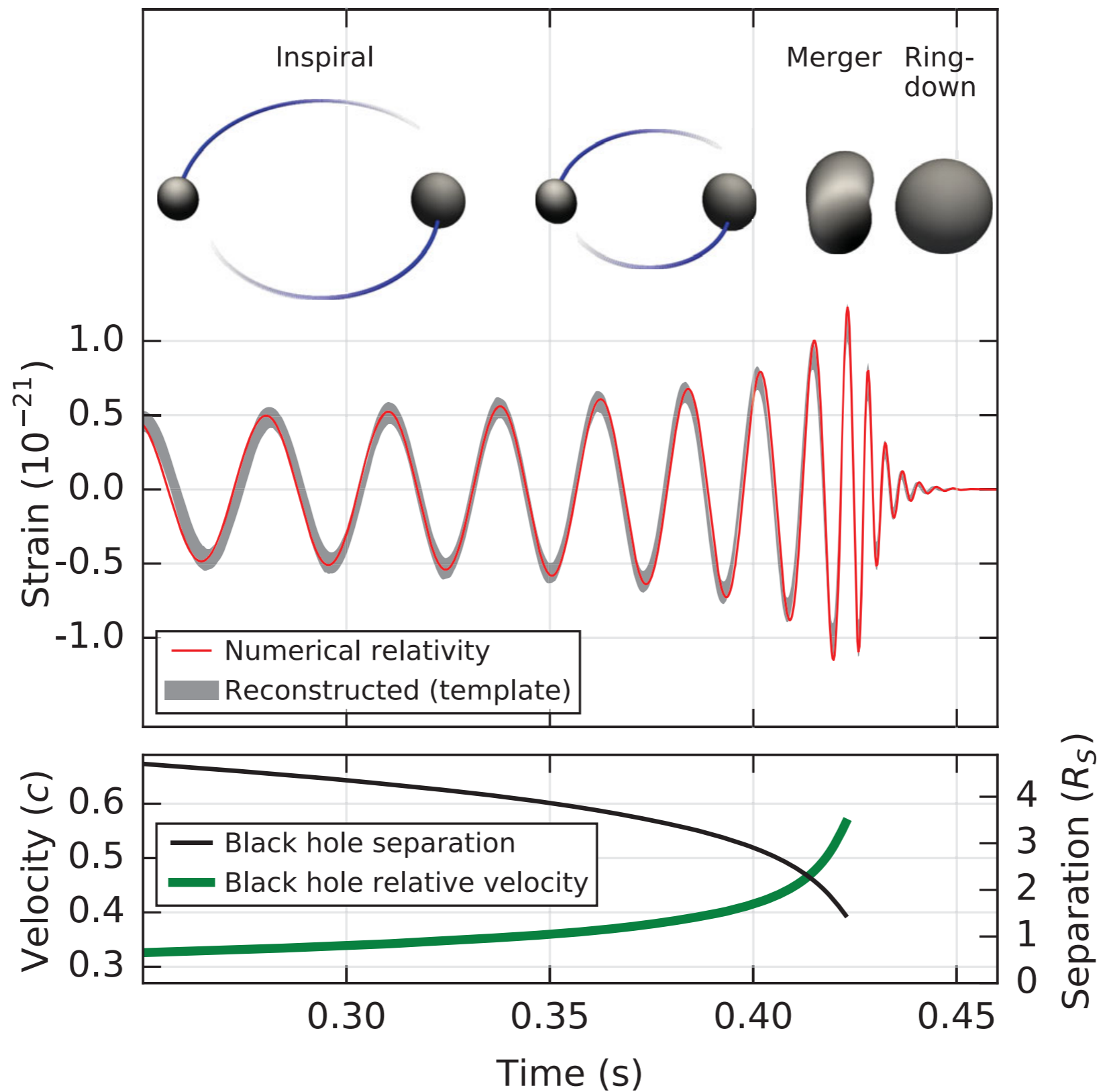
GRAMPA meeting - 30 august 2016

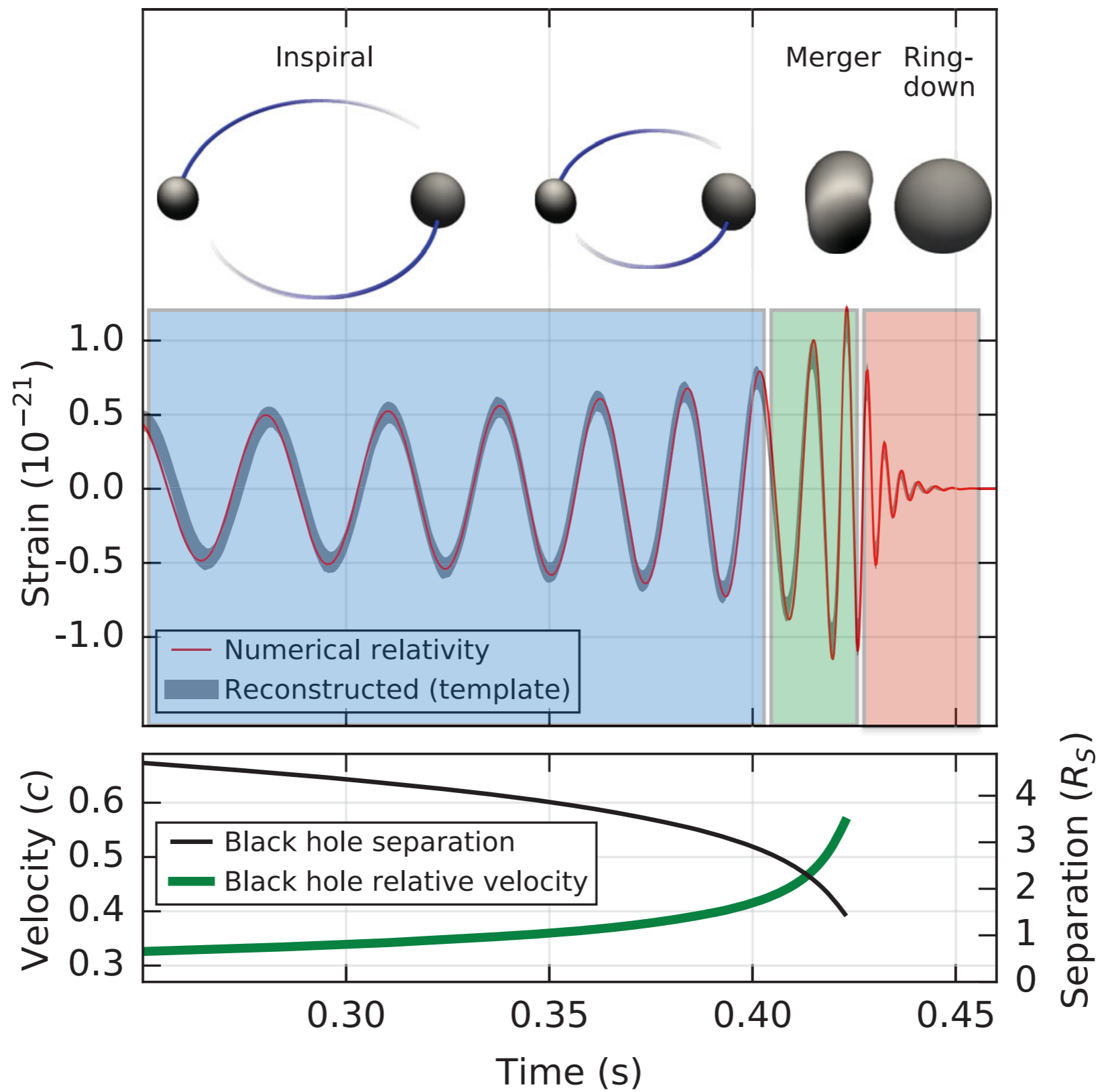
I. Introduction

II. The multipolar-post-Minkowskian algorithm

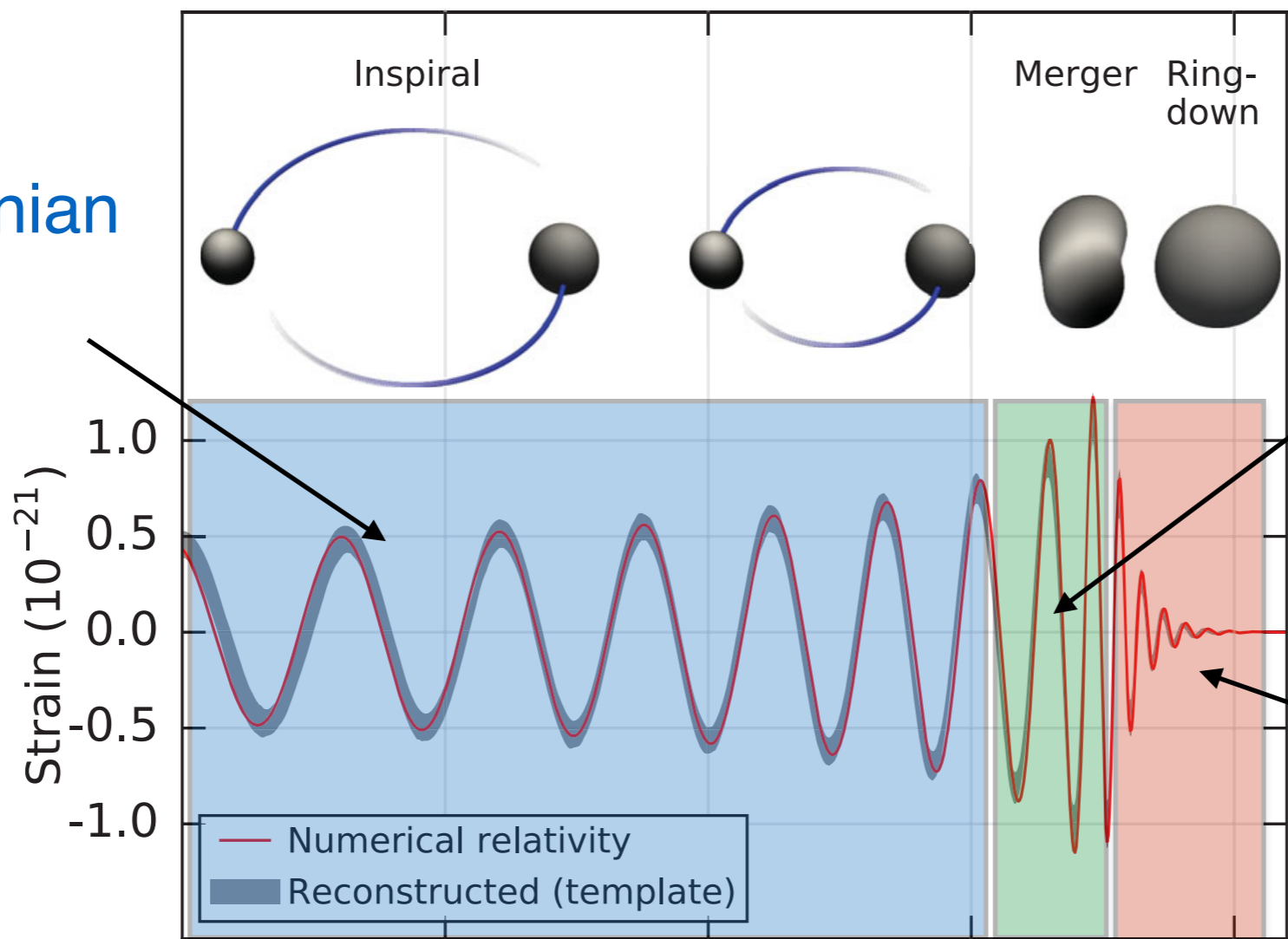
III. The 4.5PN project

I. INTRODUCTION



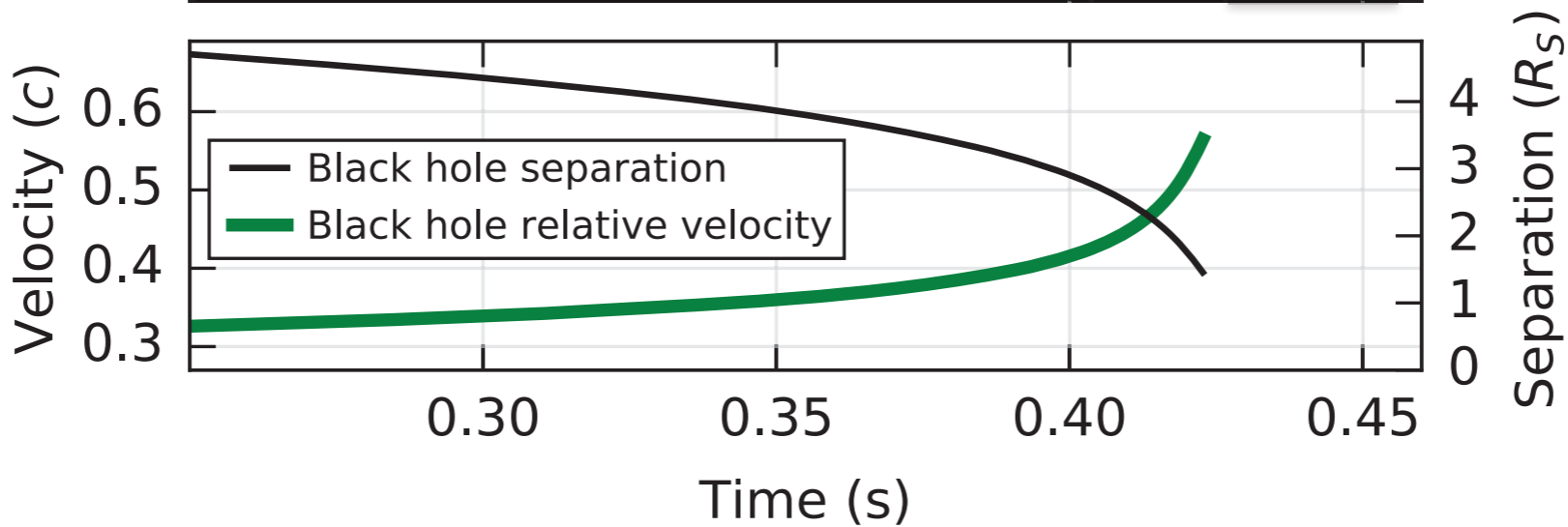


Post-Newtonian theory



Numerical relativity

BH perturbations QNM



PRL 116, 061102 (2016)

Post-Newtonian theory

- ▶ Perturbative expansion of relativistic effects

- ▶ 1 PN $\rightarrow \left(\frac{v}{c}\right)^2$

- ▶ More and more difficulties appear as we go to higher orders

Blanchet-Damour-Iyer formalism

Blanchet-Damour-Iyer formalism

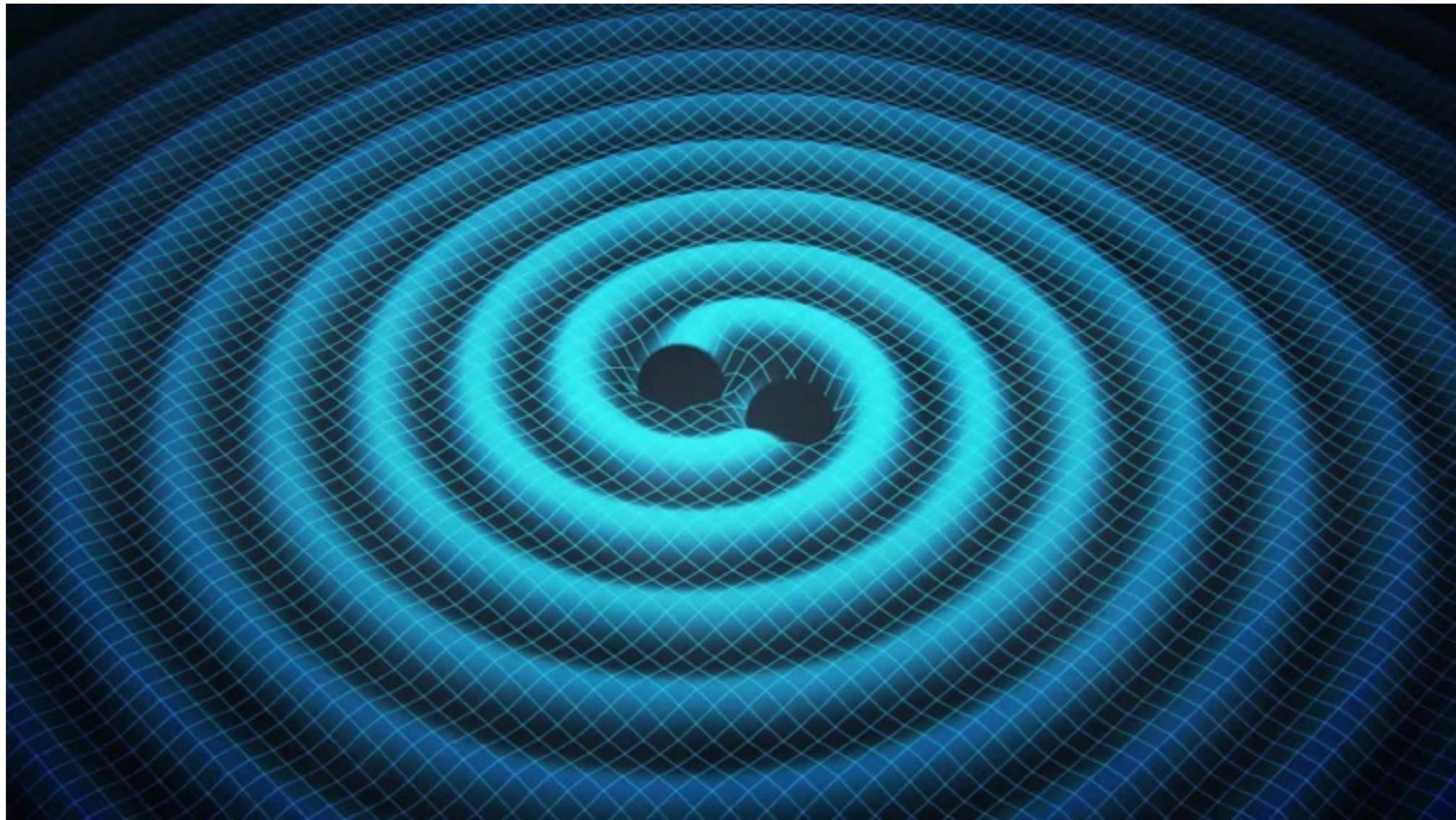


figure: www.virgo-gw.eu

Blanchet-Damour-Iyer formalism

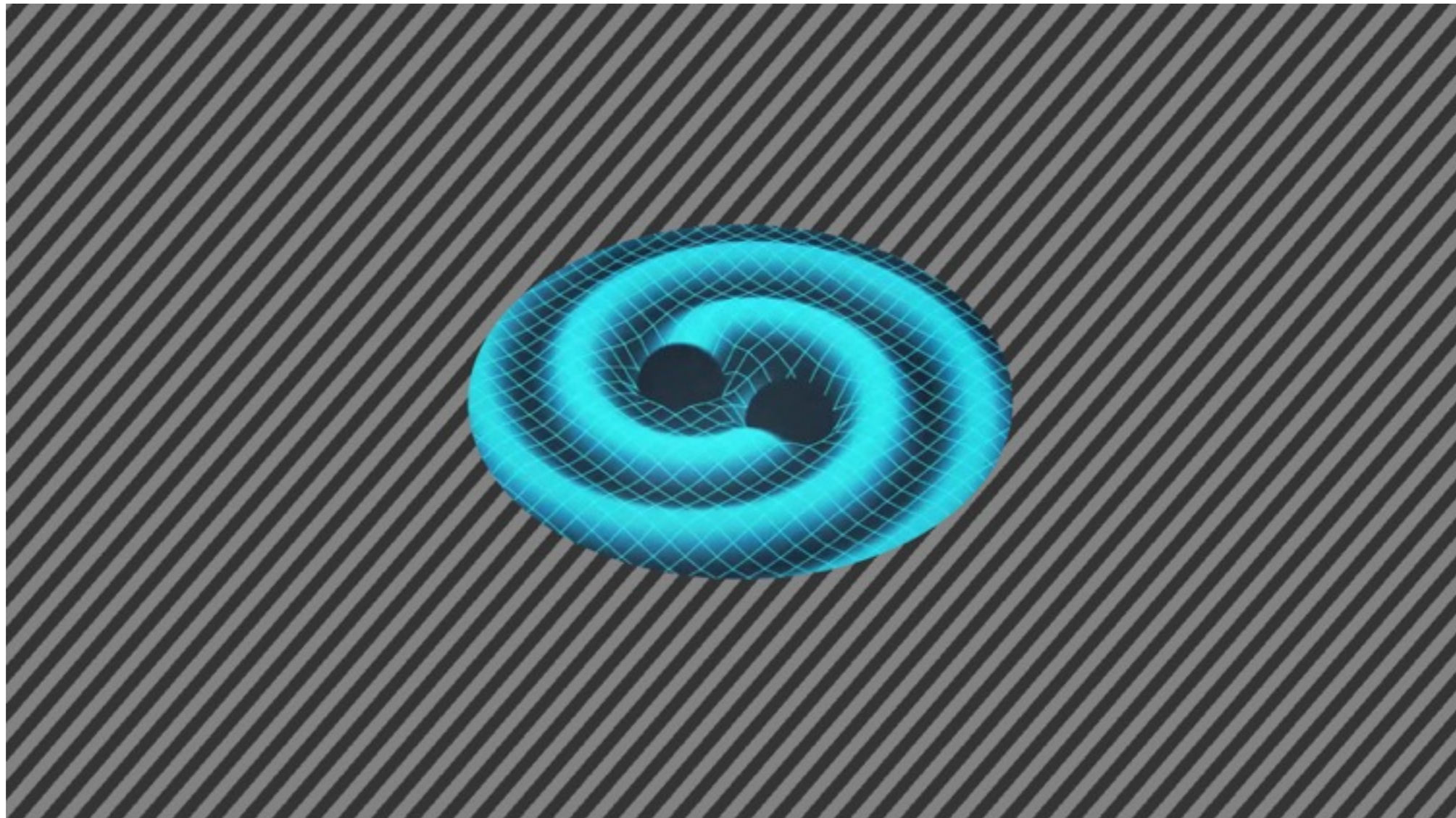


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Blanchet-Damour-Iyer formalism

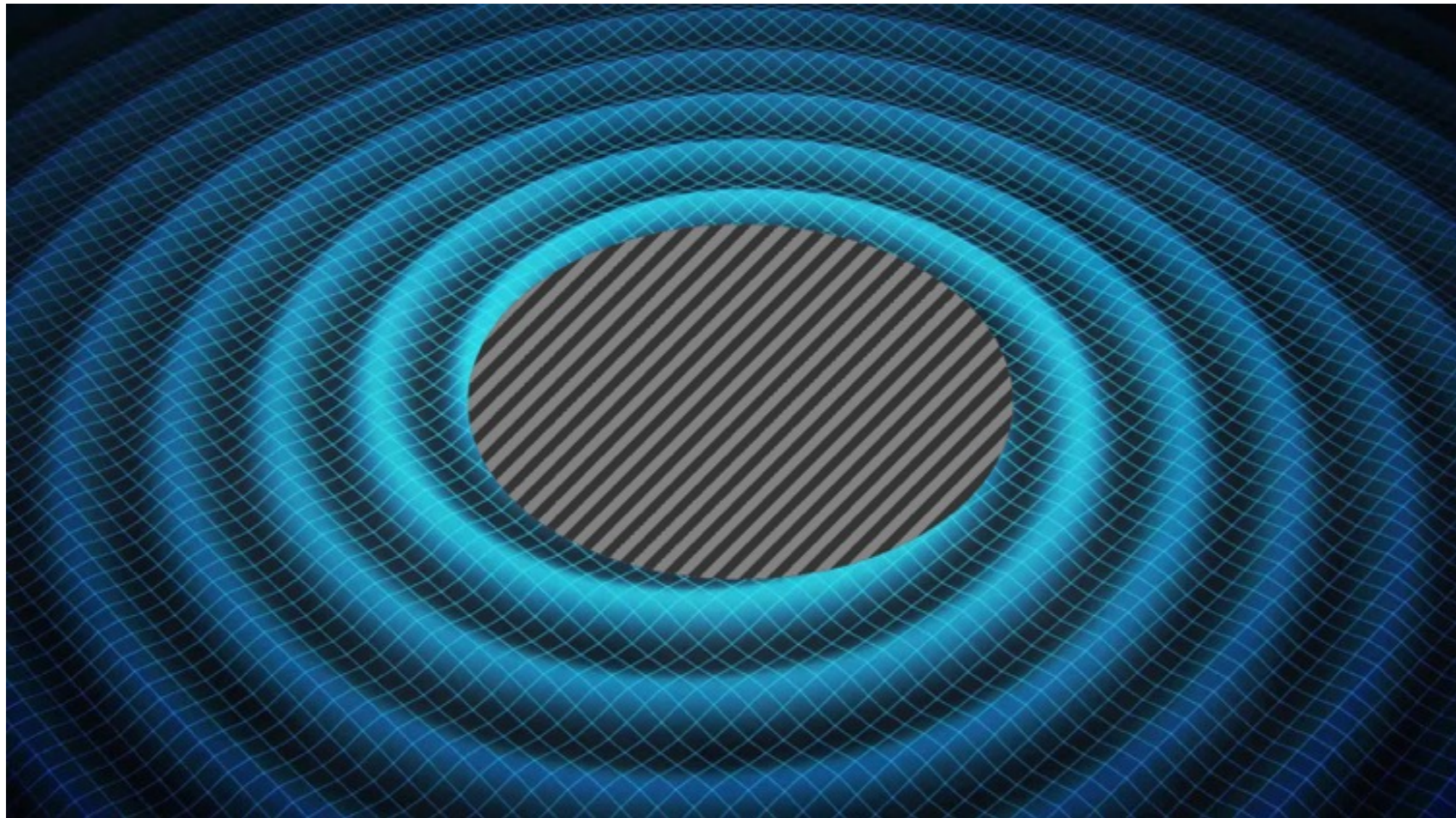


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Blanchet-Damour-Iyer formalism

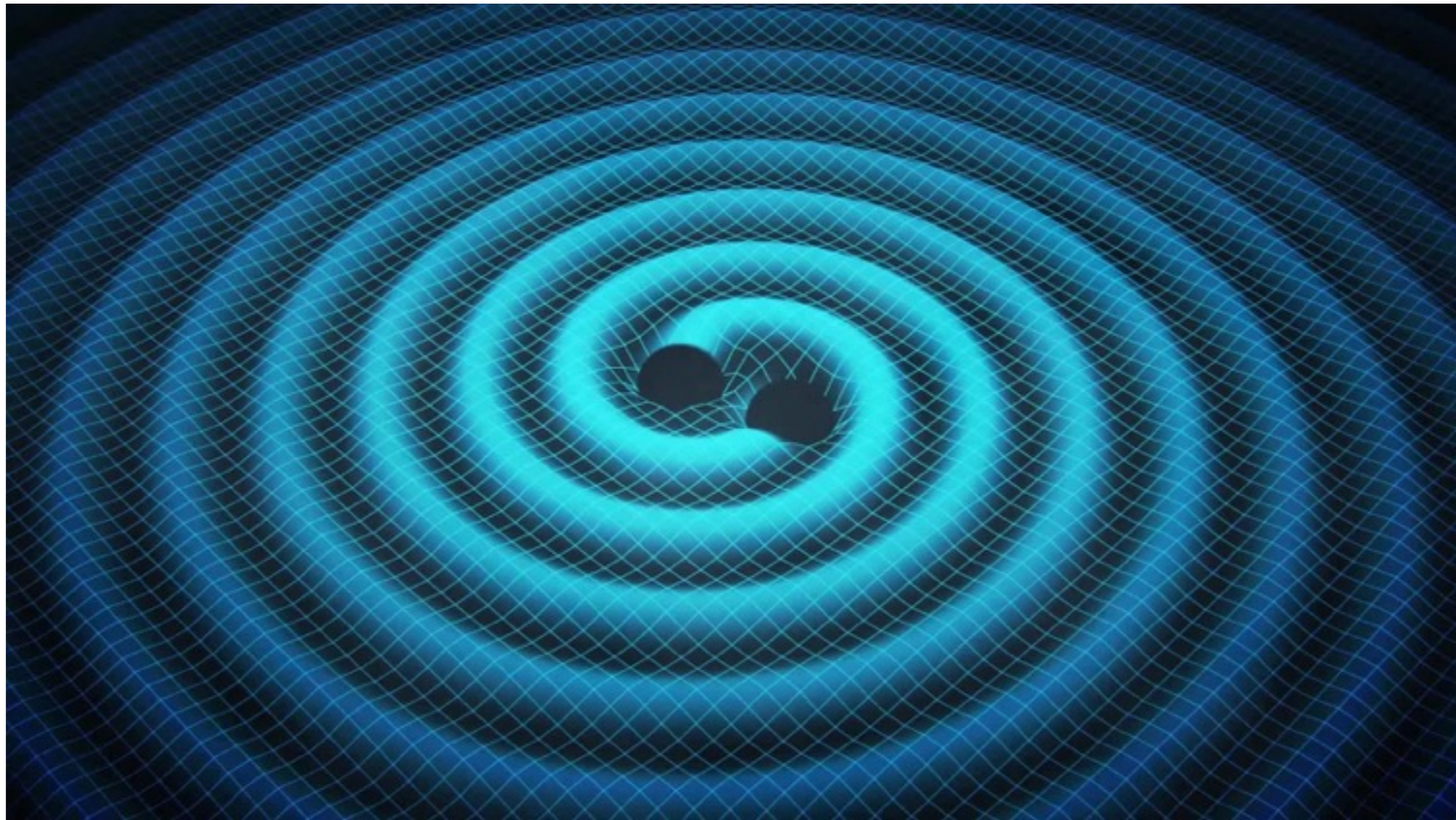


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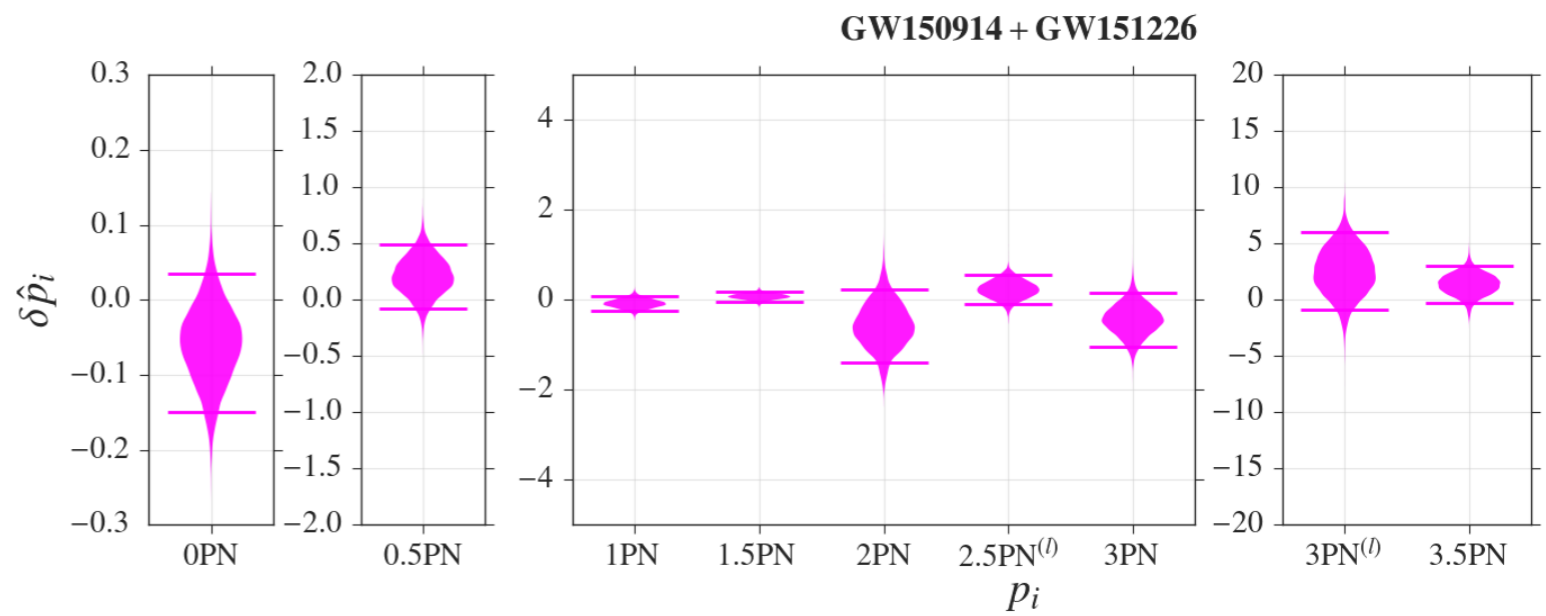
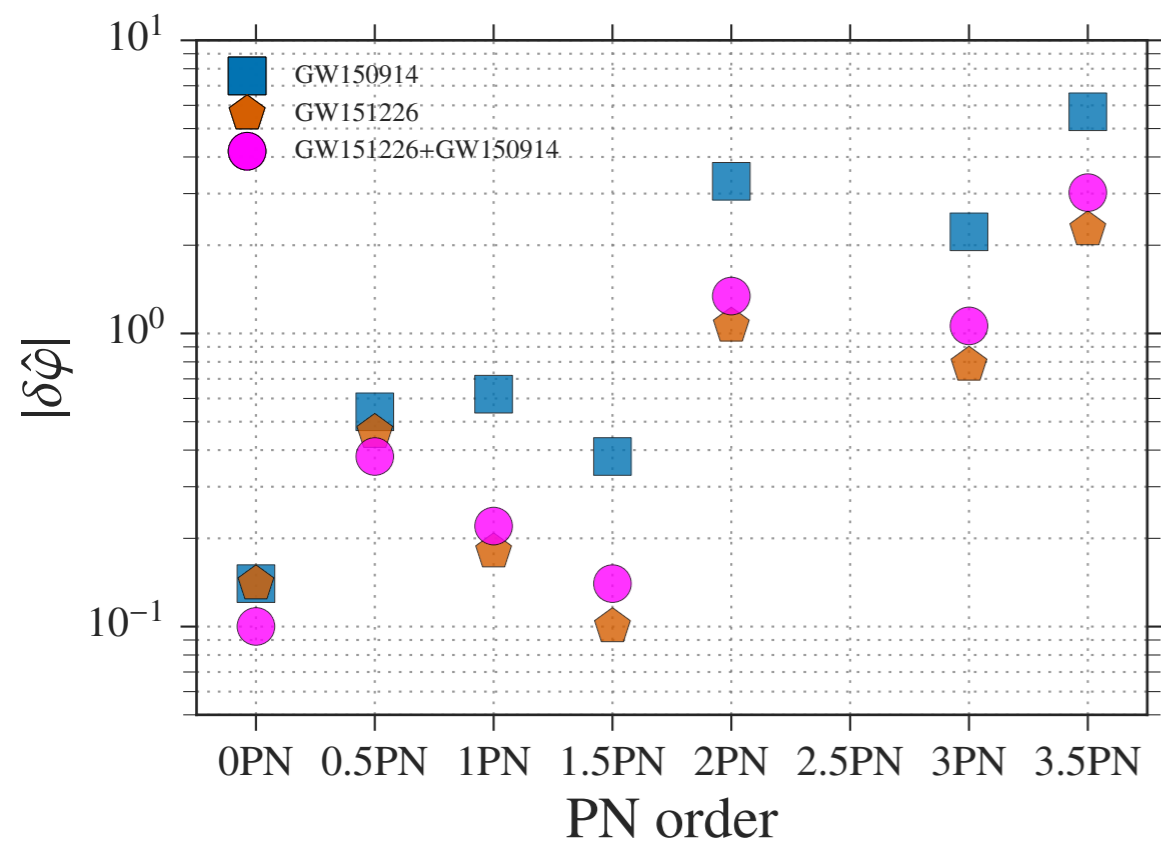
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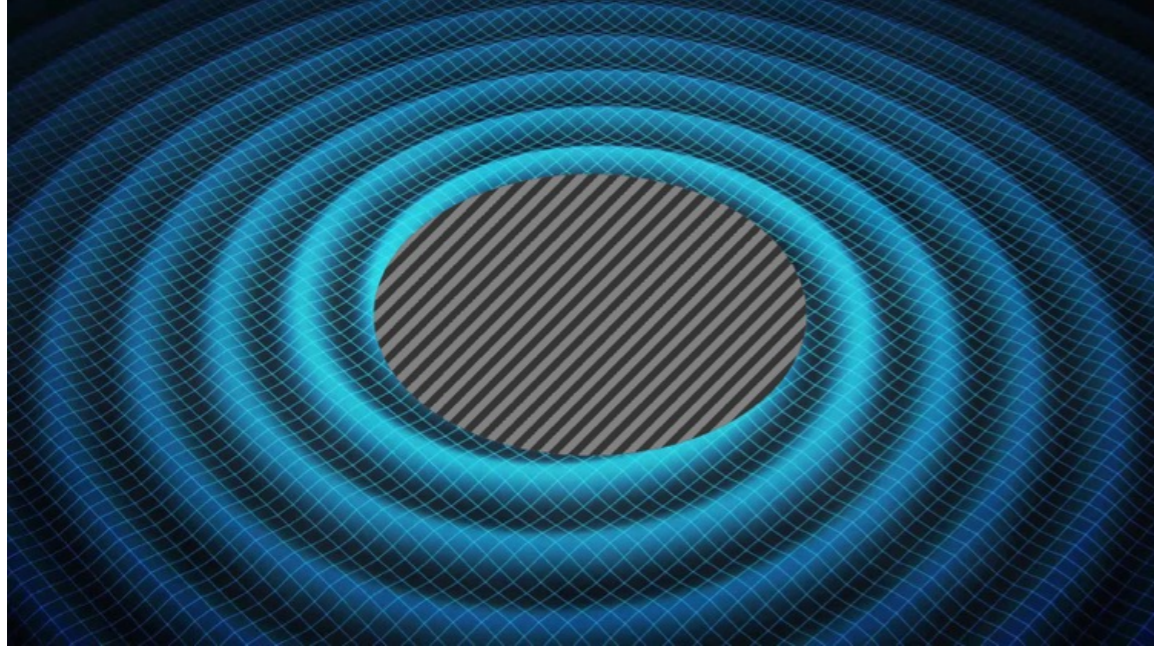
$$\frac{dE_n}{dt} = \mathcal{F}_n \quad \Rightarrow \quad \phi_n = f_n(x)$$



LIGO Scientific and Virgo collaboration arxiv:1606.04856

II. The multipolar post-Minkowskian (MPM) algorithm

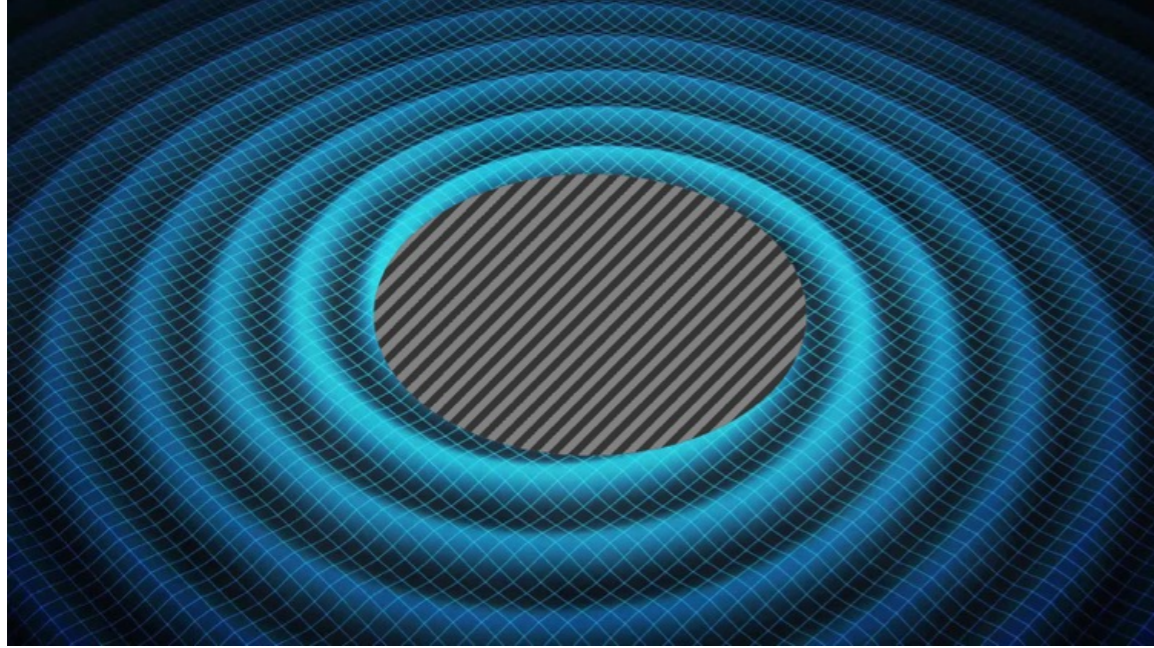
The MPM algorithm



$$G_{\mu\nu}(g_{\alpha\beta}, \partial g_{\alpha\beta}, \partial^2 g_{\alpha\beta}) = 0$$

$$h^{\mu\nu} \equiv \sqrt{-g}g^{\mu\nu} - \eta^{\mu\nu} = \mathcal{G}h^{1\mu\nu} + \mathcal{G}^2h^{2\mu\nu} + \dots$$

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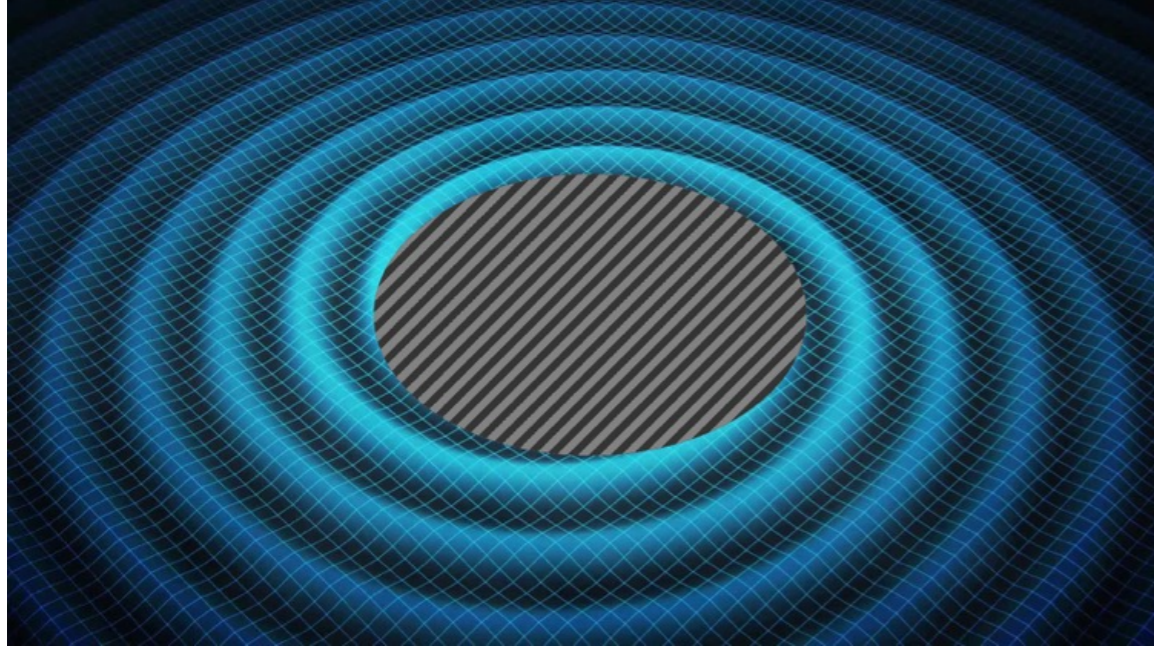


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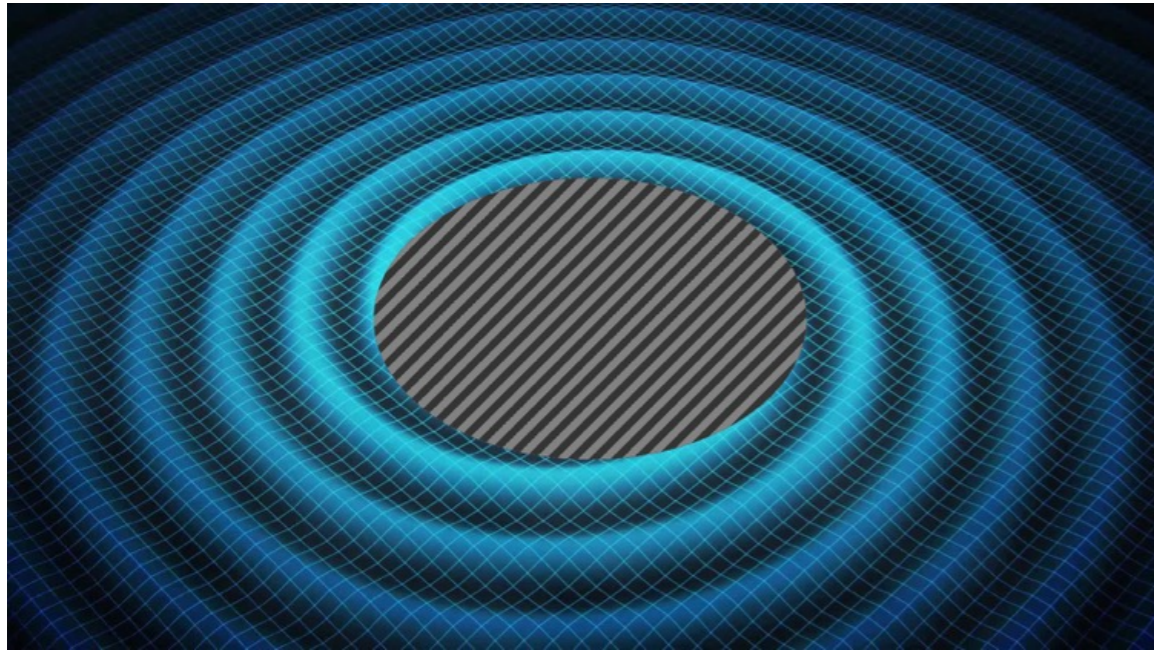
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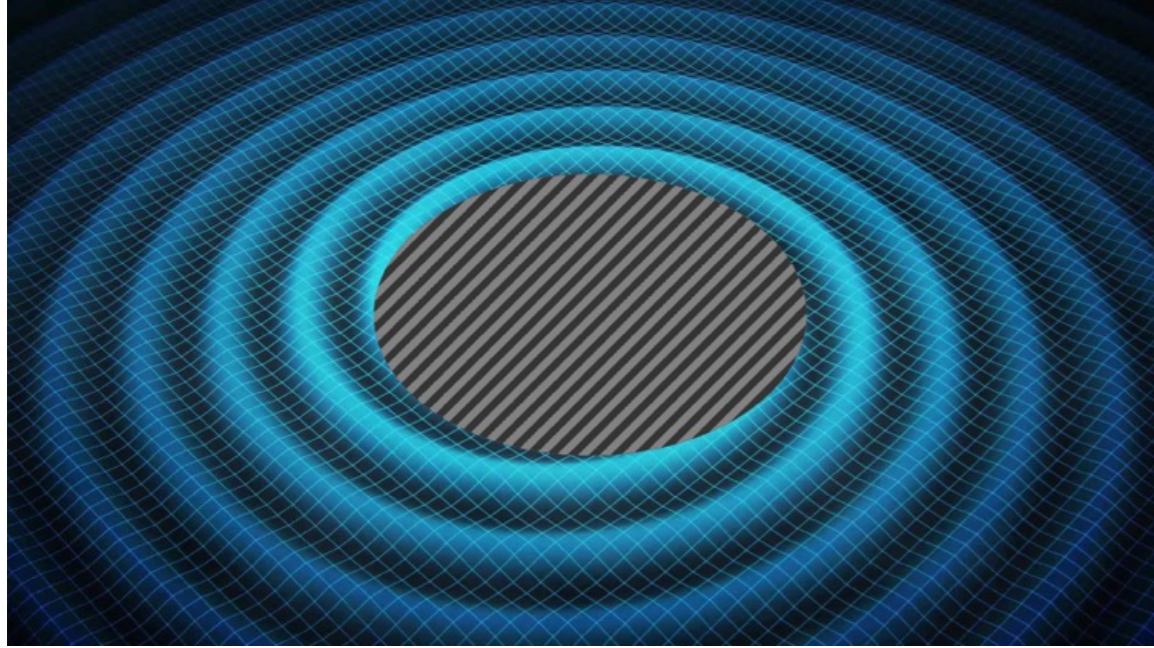
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$$\begin{cases} \square h^i_{\mu\nu} = \Lambda(h^1, \dots, h^{i-1}) \\ \partial^\mu h^i_{\mu\nu} = 0 \end{cases}$$

$$\begin{aligned} h^1_{\mu\nu} &\sim \sum_{l \geq 0} \partial_{i_1, \dots, i_l} \left(\frac{M_{i_1 \dots i_l}(t-r)}{r} \right) + \sum_{l \geq 2} \partial_{i_1, \dots, i_l} \left(\frac{S_{i_1 \dots i_l}(t-r)}{r} \right) \\ &= h^1_M + h^1_{M_{ij}} + h^1_{M_{ijk}} + \dots + h^1_{S_{ij}} + h^1_{S_{ijk}} + \dots \end{aligned}$$

The MPM algorithm

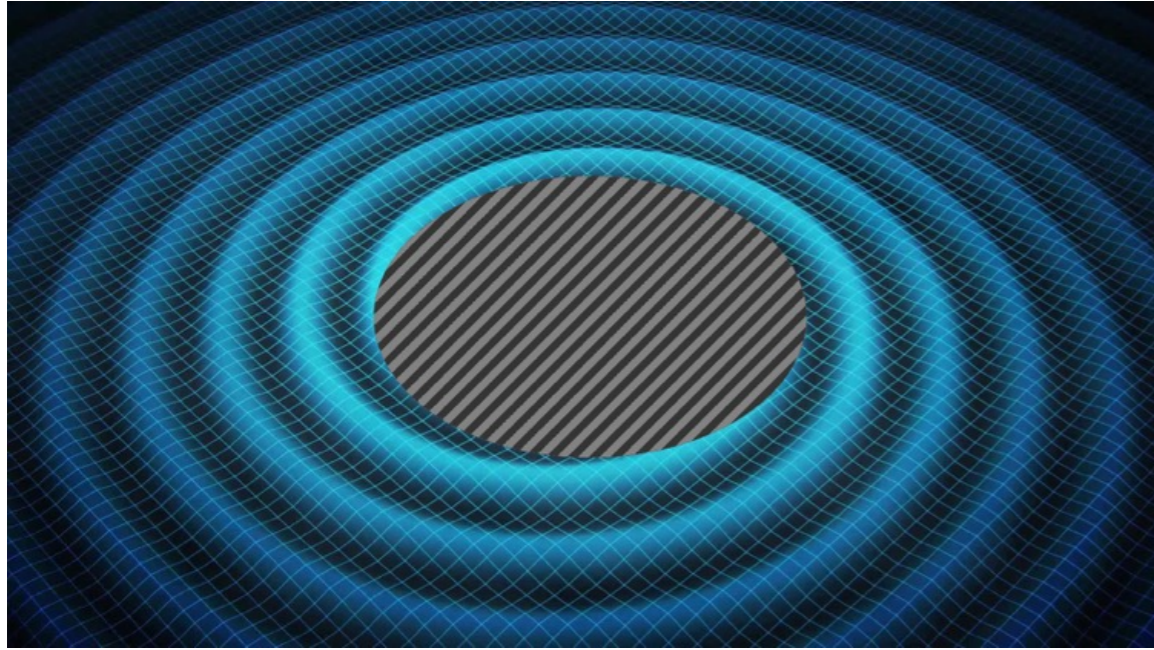


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$$h^2_{\mu\nu} = h^2_{M \times M} + h^2_{M \times M_{ij}} + h^2_{M_{ij} \times M_{ij}} + \dots$$

First issue: regularization

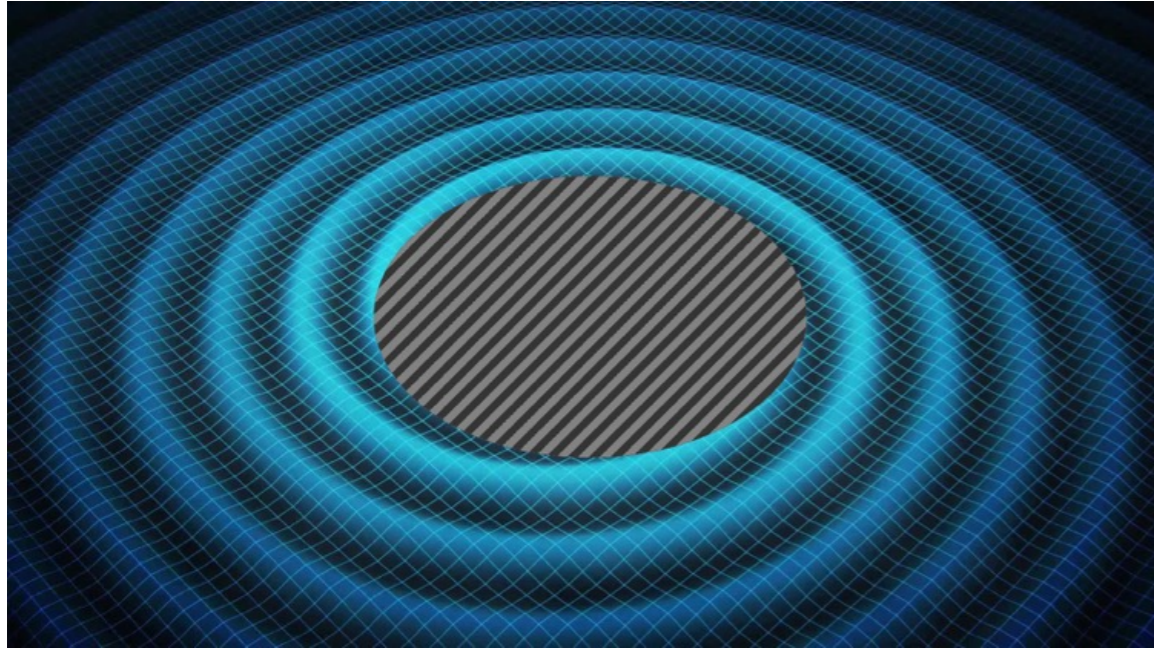


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$$\square^{-1}\Lambda(x, t) = \int d^3x' \frac{\Lambda(x', t - |x - x'|)}{|x - x'|}$$

Issue: $\Lambda \sim_{r \rightarrow 0} \frac{1}{r^k}, k \geq 3$

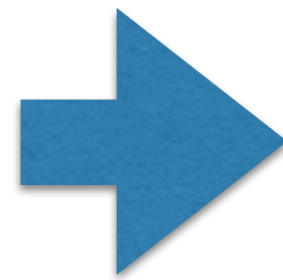
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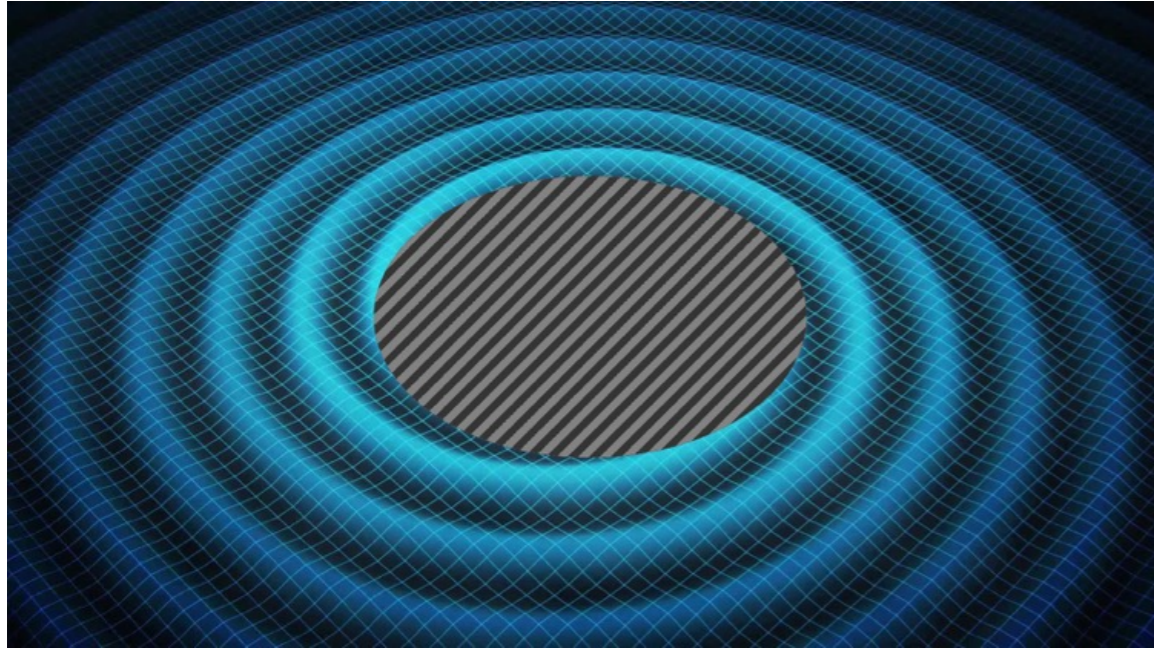
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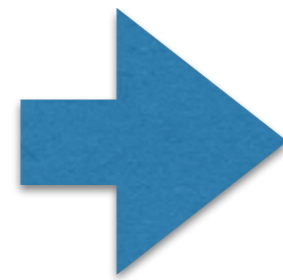
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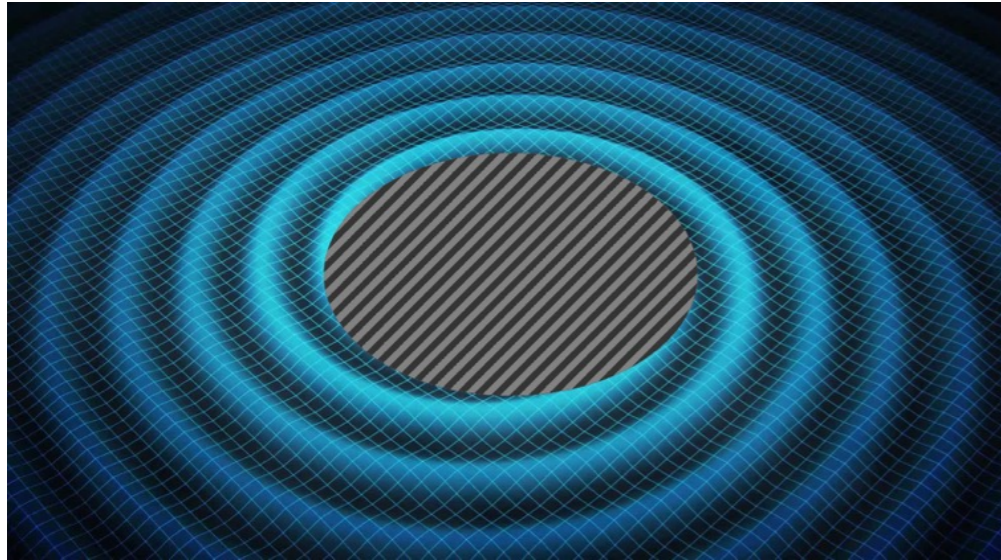
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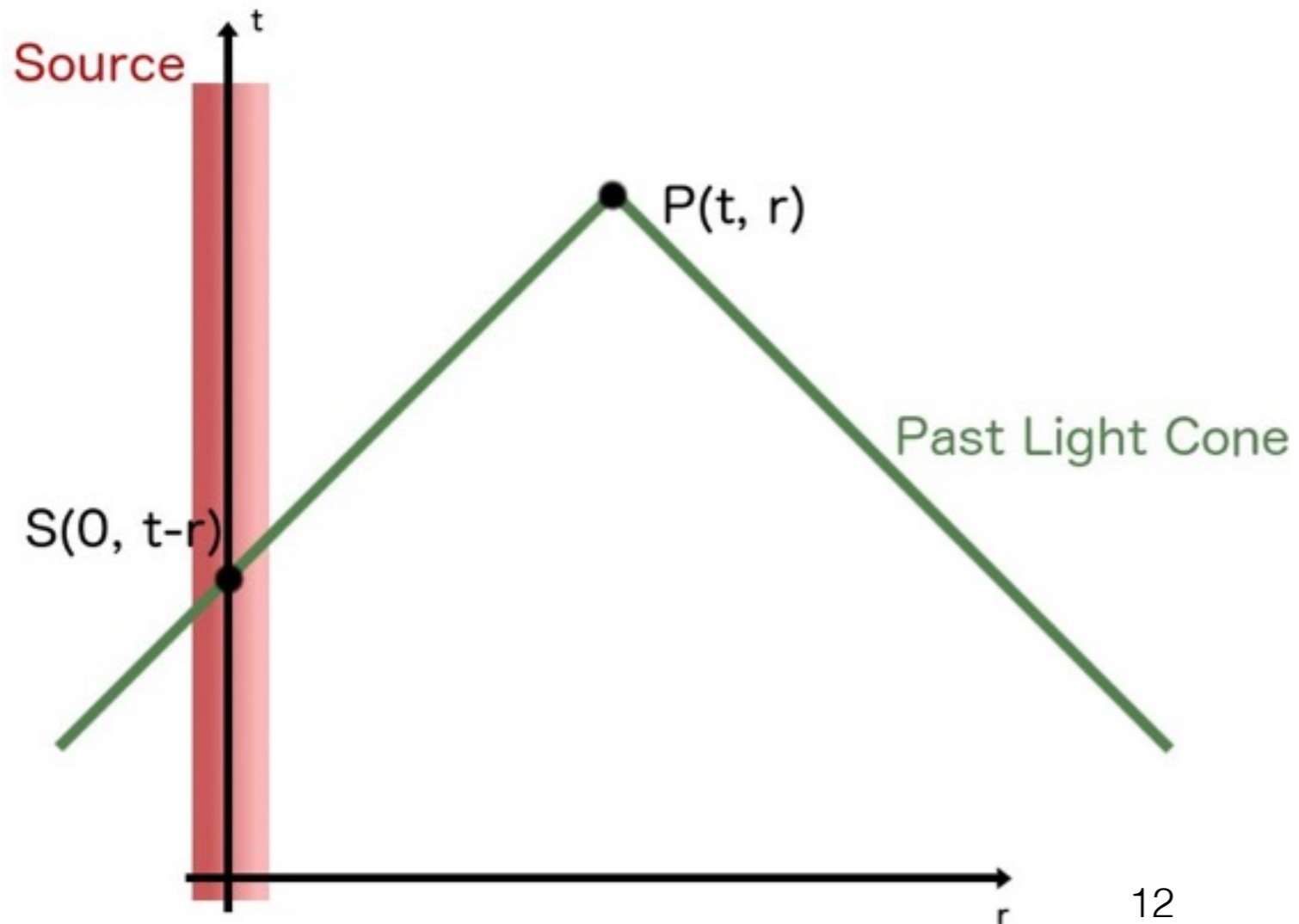
$$\text{FP}_{B=0} \square^{-1} \left[\left(\frac{r}{r_0} \right)^B \Lambda \right]$$

$$\text{FP}_{B=0} \left[\sum_{k \geq -k_0} g_k B^k \right] \equiv g_0$$

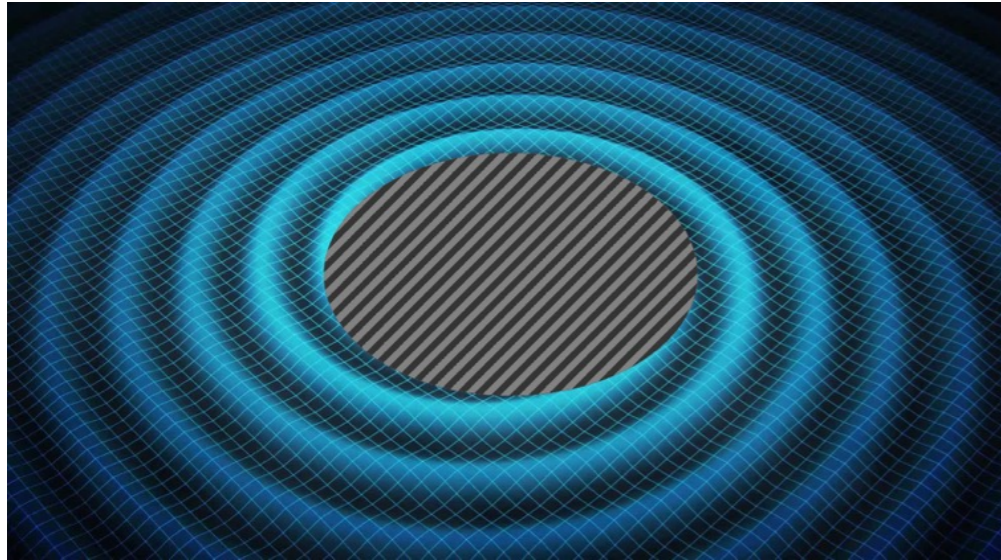
Second issue: tails



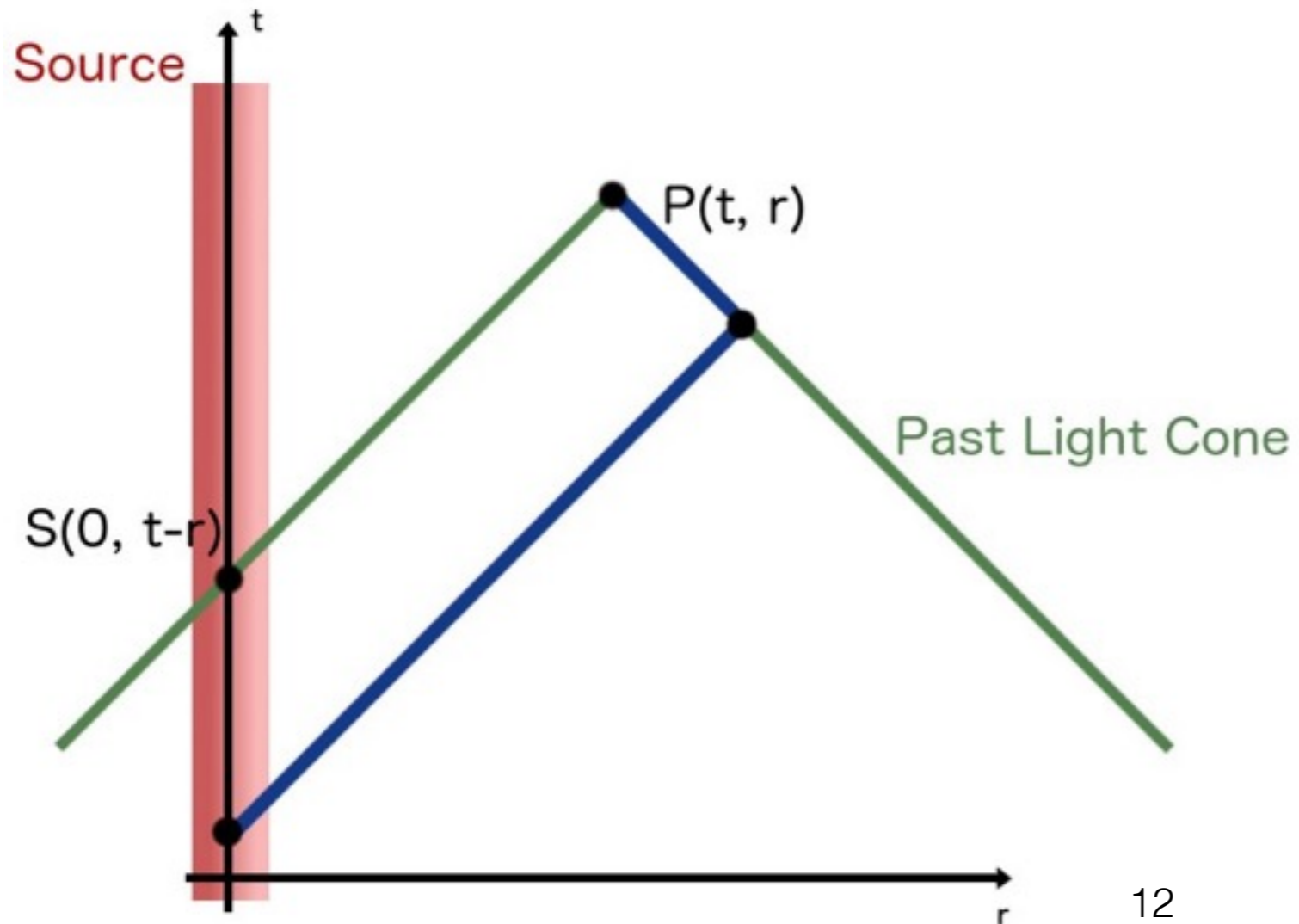
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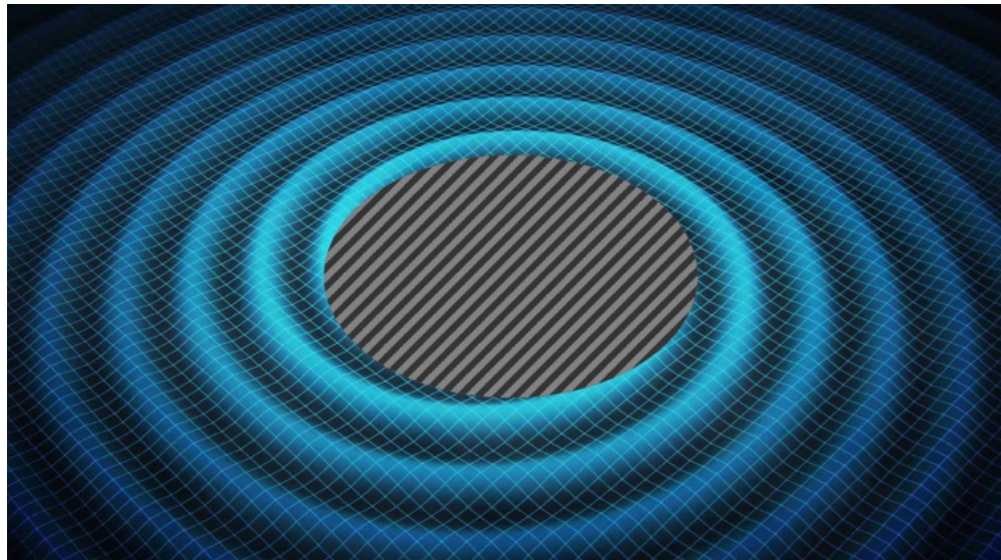
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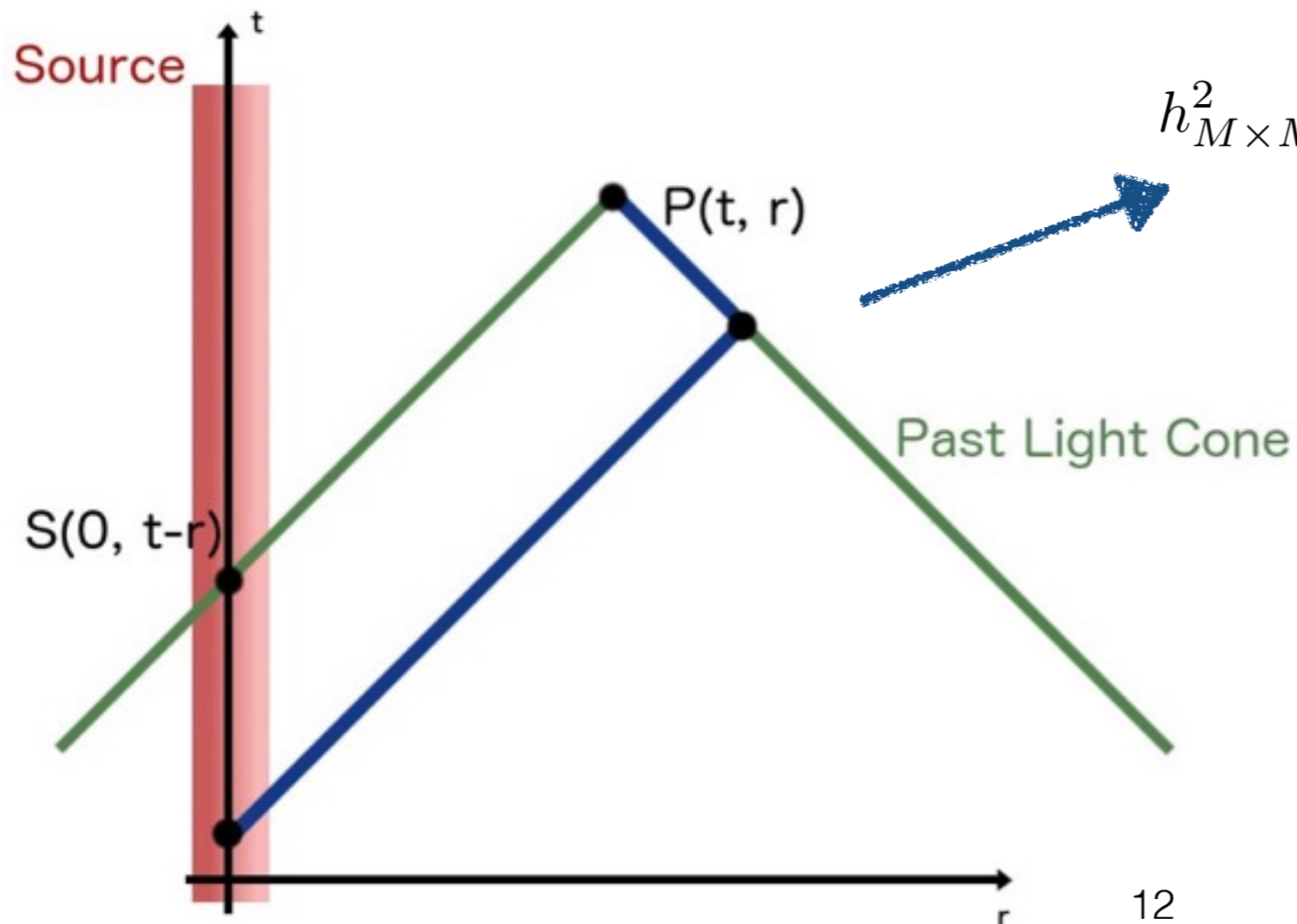
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$$h_{M \times M_{ij}}^2(t, r) \sim M \int_0^\infty d\tau M_{ij}(t - r - \tau) \mathcal{Q}(\tau)$$

III. The 4.5PN project

4.5PN project

Ultimate Goal: compute the flux up to 4.5PN

Done so far: compute all the 4.5PN contributions of the tails:

$$h_{M \times M_{ij}}^2, \quad h_{M \times M \times M_{ij}}^3, \quad h_{M \times M \times M \times M_{ij}}^4$$

arxiv:1607.07601

▶ Required new analytical formulae

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$$\begin{aligned} & \text{FP}_{B=0} \square^{-1} \left[\hat{n}_L \left(\frac{r}{r_0} \right)^B r^{-k} \int_1^\infty dy Q_m(y) F(t - ry) \right] \\ & = -\hat{n}_L \int_1^\infty ds F^{(k-2)}(t - rs) \left(Q_l(s) \int_1^s dy Q_m^{(-k+2)}(y) P_l(y) + P_l(s) \int_s^\infty dy Q_m^{(-k+2)}(y) Q_l(y) \right) \end{aligned}$$

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$$\begin{aligned} & \text{FP}_{B=0} \square^{-1} \left[\hat{n}_L \left(\frac{r}{r_0} \right)^B \ln \left(\frac{r}{r_0} \right) \frac{F(t - r)}{r^2} \right] \\ &= -\frac{\hat{n}_L}{2r} \int_r^\infty ds F(t - s) Q_m \left(\frac{s}{r} \right) \left[\ln \left(\frac{s^2 - r^2}{4r_0^2} \right) + 2H_l \right] \end{aligned}$$

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► Implementing the algorithm into Mathematica

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► Implementing the algorithm into Mathematica

 $h_{M \times M}^2, h_{M \times M \times M}^3, h_{M \times M \times M \times M}^4$

Going to future null infinity

→ $h^2_{M \times M_{ij}}$, $h^3_{M \times M \times M_{ij}}$, $h^4_{M \times M \times M \times M_{ij}}$ but $h \underset{\substack{r \rightarrow \infty \\ t-r=\text{const}}}{\sim} \frac{\ln^k r}{r}$

Going to future null infinity

→ $h^2_{M \times M_{ij}}$, $h^3_{M \times M \times M_{ij}}$, $h^4_{M \times M \times M \times M_{ij}}$ but $h \underset{\substack{r \rightarrow \infty \\ t-r=\text{const}}}{\sim} \frac{\ln^k r}{r}$

→ $X^\mu = x^\mu + \xi^\mu(x) \quad \begin{cases} \xi^0 = -2M \ln\left(\frac{r}{b_0}\right) \\ \xi^i = 0 \end{cases}$

Going to future null infinity

→ $h_{M \times M_{ij}}^2, h_{M \times M \times M_{ij}}^3, h_{M \times M \times M \times M_{ij}}^4$ but $h \underset{\substack{r \rightarrow \infty \\ t-r=\text{const}}}{\sim} \frac{\ln^k r}{r}$

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→
$$U_{ij}(T_R) = M_{ij}^{(2)}(T_R) + \frac{GM}{c^3} \int_0^{+\infty} d\tau M_{ij}^{(4)}(T_R - \tau) \left[2 \ln\left(\frac{c\tau}{2b_0}\right) + \frac{11}{6} \right]$$

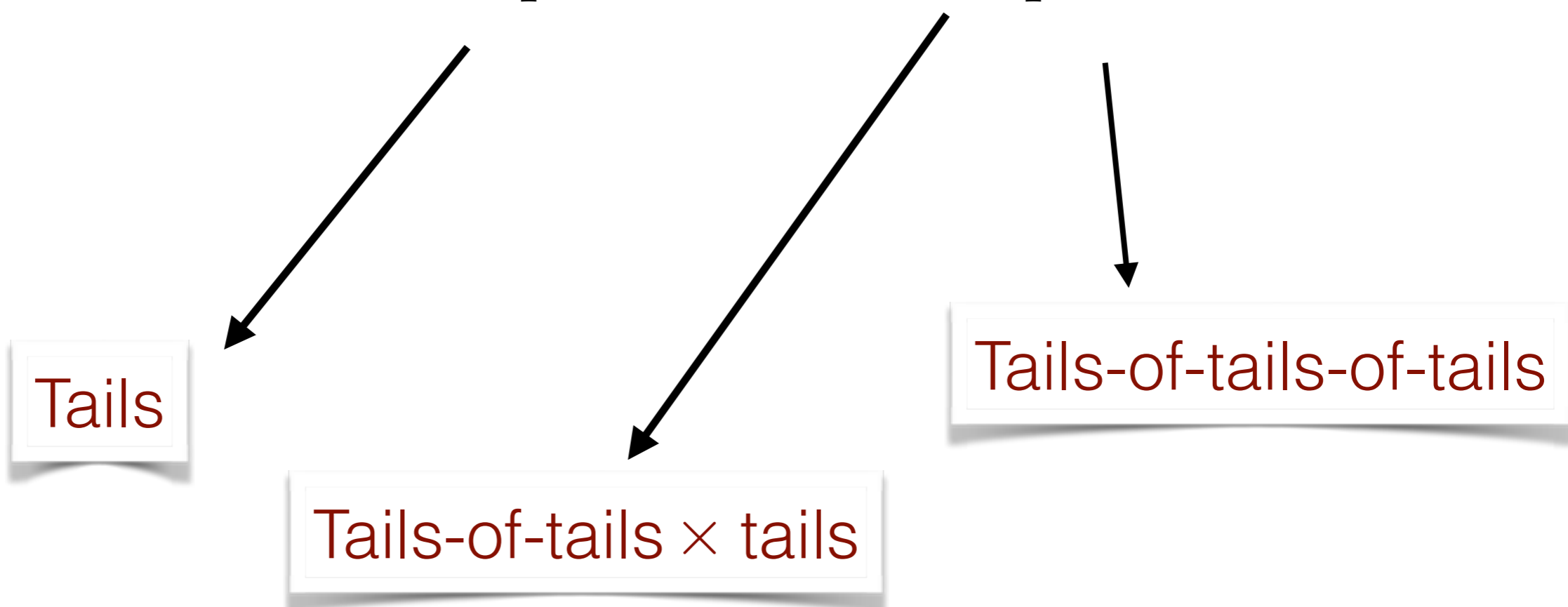
$$+ \frac{G^2 M^2}{c^6} \int_0^{+\infty} d\tau M_{ij}^{(5)}(T_R - \tau) \left[2 \ln^2\left(\frac{c\tau}{2b_0}\right) + \frac{11}{3} \ln\left(\frac{c\tau}{2b_0}\right) - \frac{214}{105} \ln\left(\frac{c\tau}{2r_0}\right) + \frac{124627}{22050} \right]$$

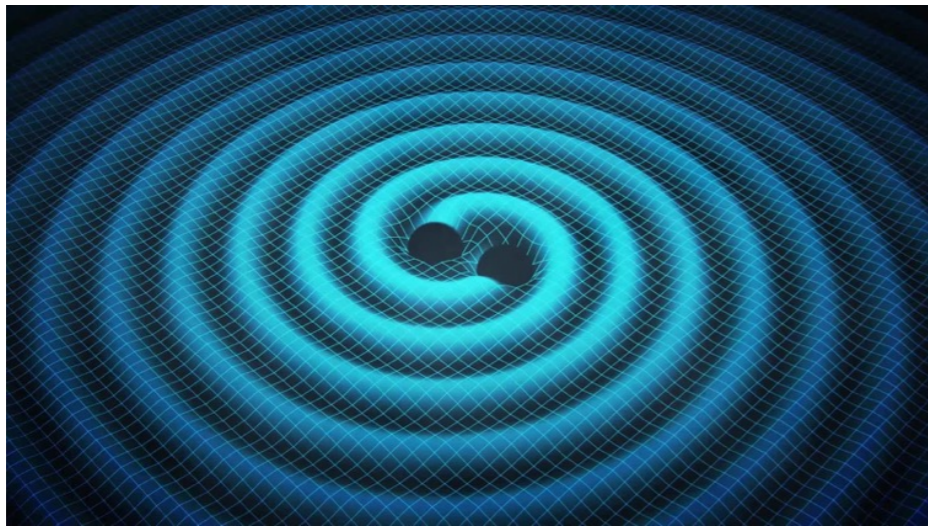
$$+ \frac{G^3 M^3}{c^9} \int_0^{+\infty} d\tau M_{ij}^{(6)}(T_R - \tau) \left[\frac{4}{3} \ln^3\left(\frac{c\tau}{2b_0}\right) + \frac{11}{3} \ln^2\left(\frac{c\tau}{2b_0}\right) + \frac{124627}{11025} \ln\left(\frac{c\tau}{2b_0}\right) - \frac{428}{105} \ln\left(\frac{c\tau}{2b_0}\right) \ln\left(\frac{c\tau}{2r_0}\right) - \frac{1177}{315} \ln\left(\frac{c\tau}{2r_0}\right) + \frac{129268}{33075} + \frac{428}{315} \pi^2 \right] + \mathcal{O}\left(\frac{1}{c^{12}}\right)$$

$$\mathcal{F} = \sum_{l=2}^{\infty} \frac{G}{c^{2l+1}} \left[a_l (U_L^{(1)})^2 + \frac{b_l}{c^2} (V_L^{(1)})^2 \right]$$

At 4.5PN, only non-local terms contribute in the flux

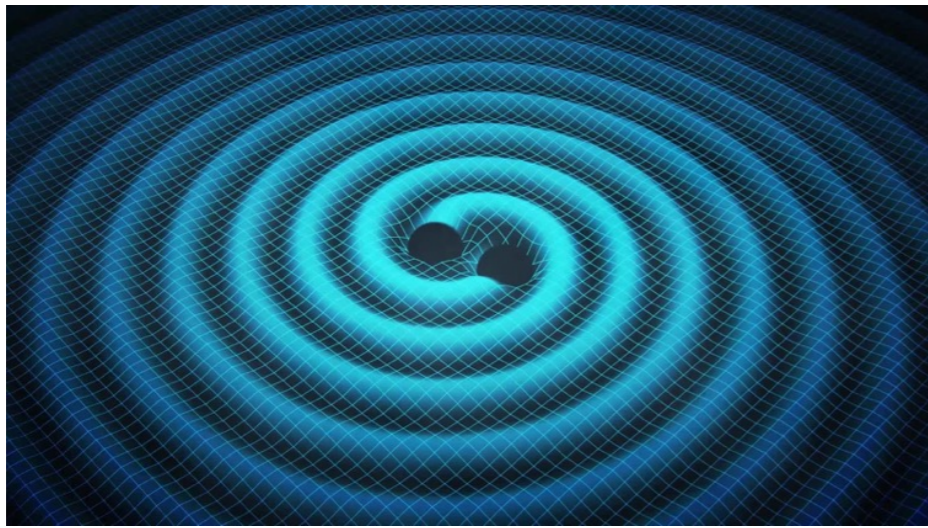
$$\mathcal{F}_{4.5\text{PN}} = \mathcal{F}_{\text{quadratic}} + \mathcal{F}_{\text{quartic}}$$





$$\nu = \frac{m_1 m_2}{(m_1 + m_2)^2}$$

$$\gamma = \frac{Gm}{rc^2}$$



Tails

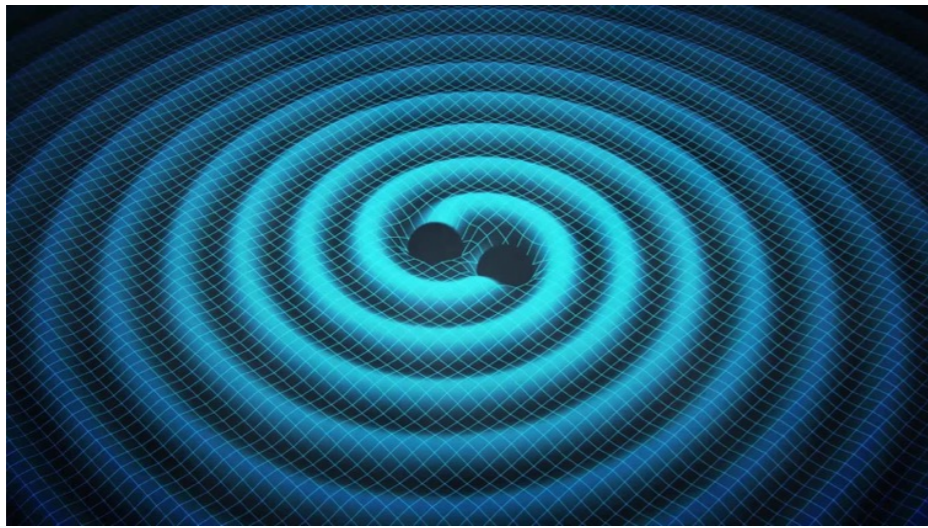
$$\nu = \frac{m_1 m_2}{(m_1 + m_2)^2}$$

$$\gamma = \frac{Gm}{rc^2}$$

$$\mathcal{F}_{\text{quadratic}} = \frac{32c^5}{5G} \nu^2 \gamma^5 \left\{ 4\pi\gamma^{3/2} + \dots \right. \\ \left. + \left(\frac{9997778801}{106444800} - \frac{6848}{105} \ln\left(\frac{r}{r_0}\right) + \left[-\frac{8058312817}{2661120} + \frac{287}{32}\pi^2 + \frac{572}{3} \ln\left(\frac{r}{r'_0}\right) \right] \nu \right. \right. \\ \left. \left. - \frac{12433367}{13824} \nu^2 - \frac{1026257}{266112} \nu^3 \right) \pi\gamma^{9/2} + \mathcal{O}\left(\frac{1}{c^{11}}\right) \right\},$$

Tails-of-tails-of-tails and Tails-of-tails \times tails

$$\mathcal{F}_{\text{quartic}} = \frac{32c^5}{5G} \nu^2 \gamma^5 \left\{ \left(-\frac{467044}{3675} - \frac{3424}{105} \ln(16\gamma) + \frac{6848}{105} \ln\left(\frac{r}{r_0}\right) - \frac{6848}{105} \gamma_E \right) \pi\gamma^{9/2} \right. \\ \left. + \mathcal{O}\left(\frac{1}{c^{11}}\right) \right\}.$$



Tails

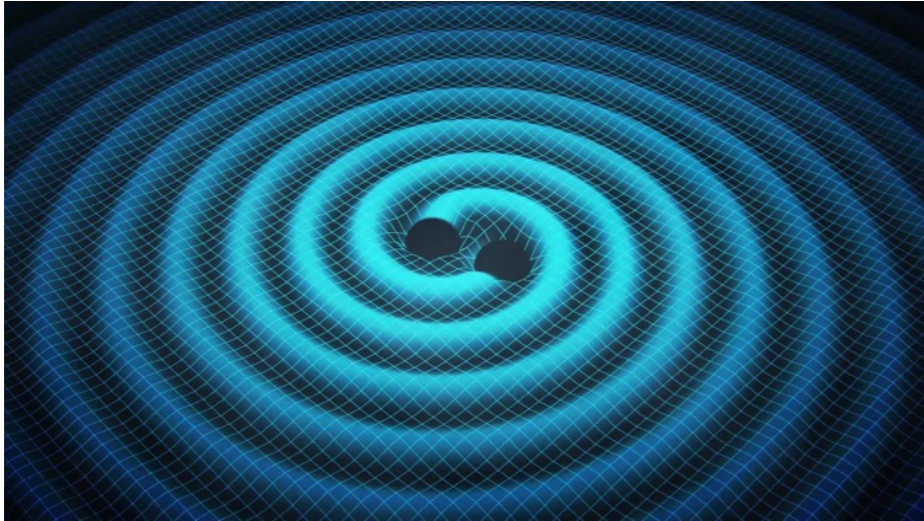
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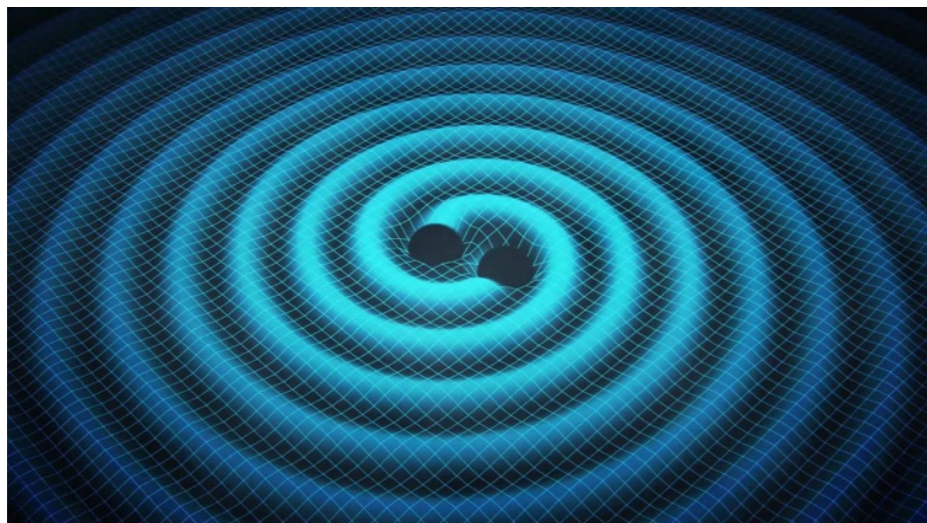
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$$\nu = \frac{m_1 m_2}{(m_1 + m_2)^2}$$

$$x = \left(\frac{Gm\Omega}{c^3} \right)^{2/3} = \mathcal{O} \left(\frac{1}{c^2} \right)$$

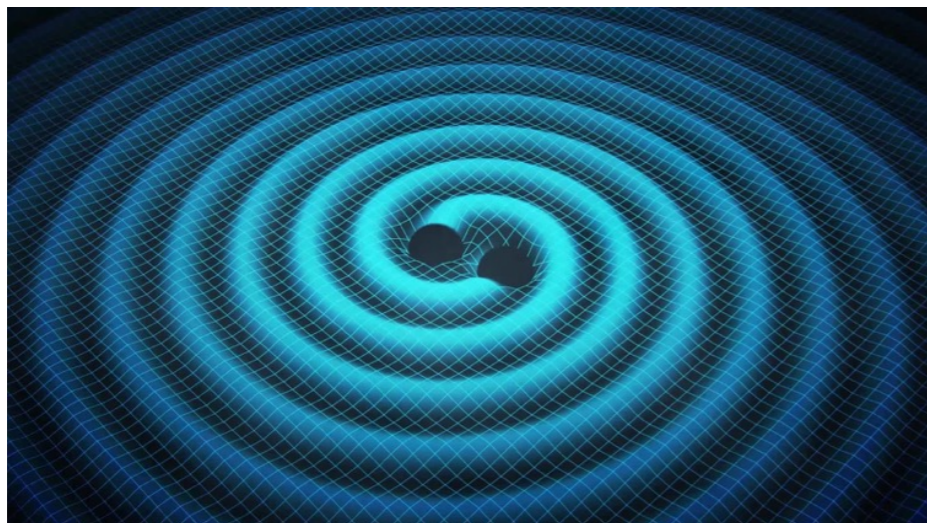


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$$x = \left(\frac{Gm\Omega}{c^3} \right)^{2/3} = \mathcal{O} \left(\frac{1}{c^2} \right)$$

$$\begin{aligned} \mathcal{F}_{\text{total}} = & \frac{32c^5}{5G} \nu^2 x^5 \left\{ 1 + \left(-\frac{1247}{336} - \frac{35}{12} \nu \right) x + 4\pi x^{3/2} + \left(-\frac{44711}{9072} + \frac{9271}{504} \nu + \frac{65}{18} \nu^2 \right) x^2 \right. \\ & + \left(-\frac{8191}{672} - \frac{583}{24} \nu \right) \pi x^{5/2} + \left[\frac{6643739519}{69854400} + \frac{16}{3} \pi^2 - \frac{1712}{105} \gamma_E \right. \\ & \left. - \frac{856}{105} \ln(16x) + \left(-\frac{134543}{7776} + \frac{41}{48} \pi^2 \right) \nu - \frac{94403}{3024} \nu^2 - \frac{775}{324} \nu^3 \right] x^3 \\ & + \left(-\frac{16285}{504} + \frac{214745}{1728} \nu + \frac{193385}{3024} \nu^2 \right) \pi x^{7/2} + (\text{unknown coefficients}) x^4 \\ & + \left(\frac{265978667519}{745113600} - \frac{6848}{105} \gamma_E - \frac{3424}{105} \ln(16x) + \left[\frac{2062241}{22176} + \frac{41}{12} \pi^2 \right] \nu \right. \\ & \left. - \frac{133112905}{290304} \nu^2 - \frac{3719141}{38016} \nu^3 \right) \pi x^{9/2} + \mathcal{O}(x^5) \left. \right\}. \end{aligned}$$

arxiv:1607.07601



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Thank you !